

Particle Detectors

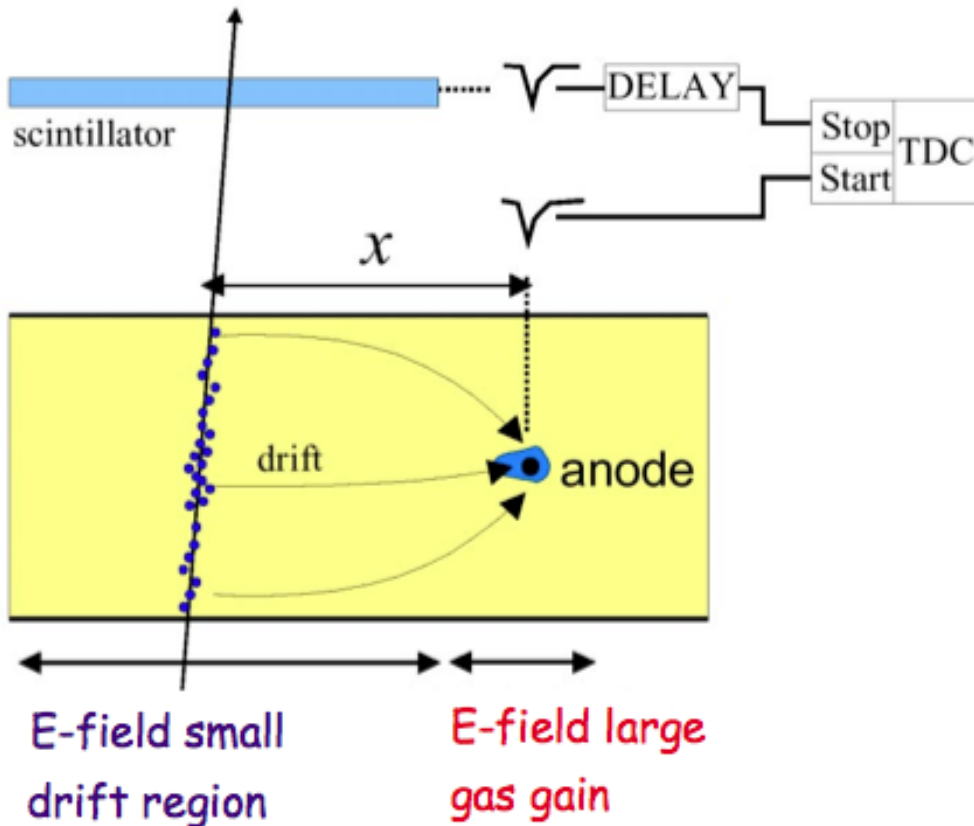
Lecture 17 *16/05/18*

a.a. 2017-2018

Emanuele Fiandrini

Camere a deriva

Possono essere considerate come derivate delle camere proporzionali. In questo caso ricavo la coordinata misurando un tempo.



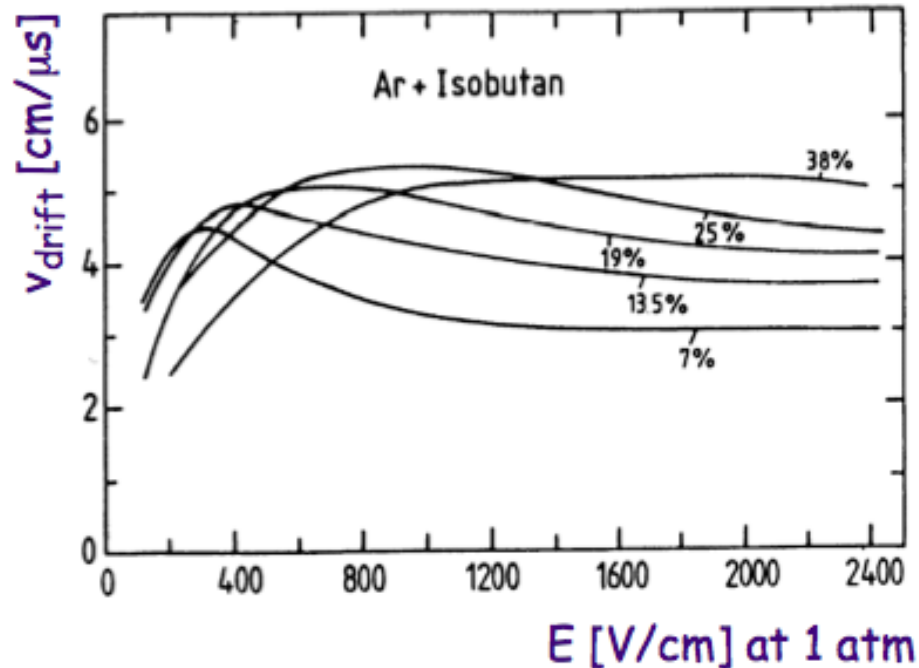
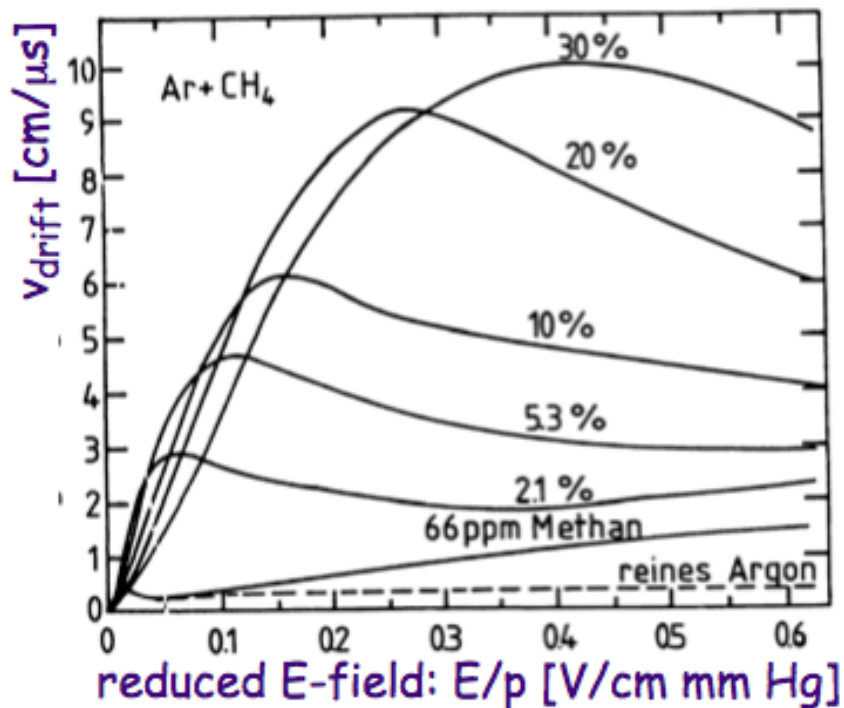
TDC: Time to Digital Converter

Measure arrival time t_1 of electrons at anode wire relative to reference t_0 .

- external definition of time reference t_0 (here by fast scintillator signal)
- x-coordinate given by:

$$x = \int_{t_0}^{t_1} v_D(t) dt$$

- if drift velocity v_D constant over full drift distance: $x = v_D(t_1 - t_0) = v_D \Delta t$
- advantage of drift chambers: much **larger sensitive volume** per read out channel



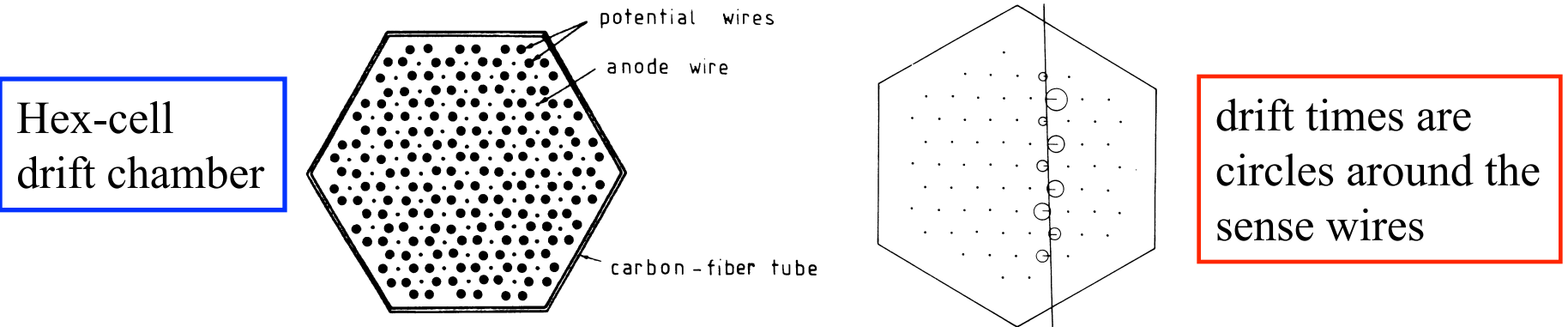
- strong dependence on the choice of the gas mixture
- details of the energy dependence of the ionisation cross section (Ramsauer minimum) result in a characteristic maximum of the E field dependence.
- for stable operation it is useful to operate in the maximum: $\frac{dv_{drift}}{dE} = 0$
- typical drift velocities : $v_{drift} \approx 2-10 \text{ cm}/\mu\text{s} = 20-100 \mu\text{m}/\text{ns}$

Drift Chambers

Drift Chambers are MWPCs where the time it takes for the ions to reach the sense wire is recorded. This time info gives position info:

$$x = \int_{t_0}^{t_s} v(t) dt \quad t_0 = \text{start time}, t_s = \text{stop time} = \text{time electrons reach sense wire}$$

By using the drift time information we can improve our spatial resolution by a factor of 10 over MWPCs (1mm $\rightarrow$ 100 μm).



The trajectory is given by the straight line tangent to all the circles

Camere a deriva

Vantaggi di una camera a deriva rispetto ad una MWPC:

- ❖ più facile da costruire da un punto di vista meccanico (fili più lontani → minori le forze elettrostatiche)
- ❖ meno fili → meno elettronica (anche se più costosa)
- ❖ migliore precisione (non più limitata alla distanza dei fili $\propto (12)^{1/2}$).

I parametri fondamentali sono:

- diffusione (buono se piccola)
- velocità di deriva (ottimo se costante). Tipiche velocità di deriva (con argon-isobutano nelle proporzioni 75%-25% e campi elettrici $E \sim 700-800 \text{ V/cm}$) $\sim 50 \text{ mm}/\mu\text{s}$.

Drift Chamber – Spatial Resolution

Resolution determined by
accuracy of drift time measurement ...

Influenced by:

Diffusion [$\sigma_{\text{Diff.}} \sim \sqrt{x}$]

see above: $\sigma^2 \sim 2Dt = 2Dx/v_D \sim x \dots$

δ -electrons [$\sigma_\delta = \text{const.}$]

independent of drift length; yields constant
term in spatial resolution ...

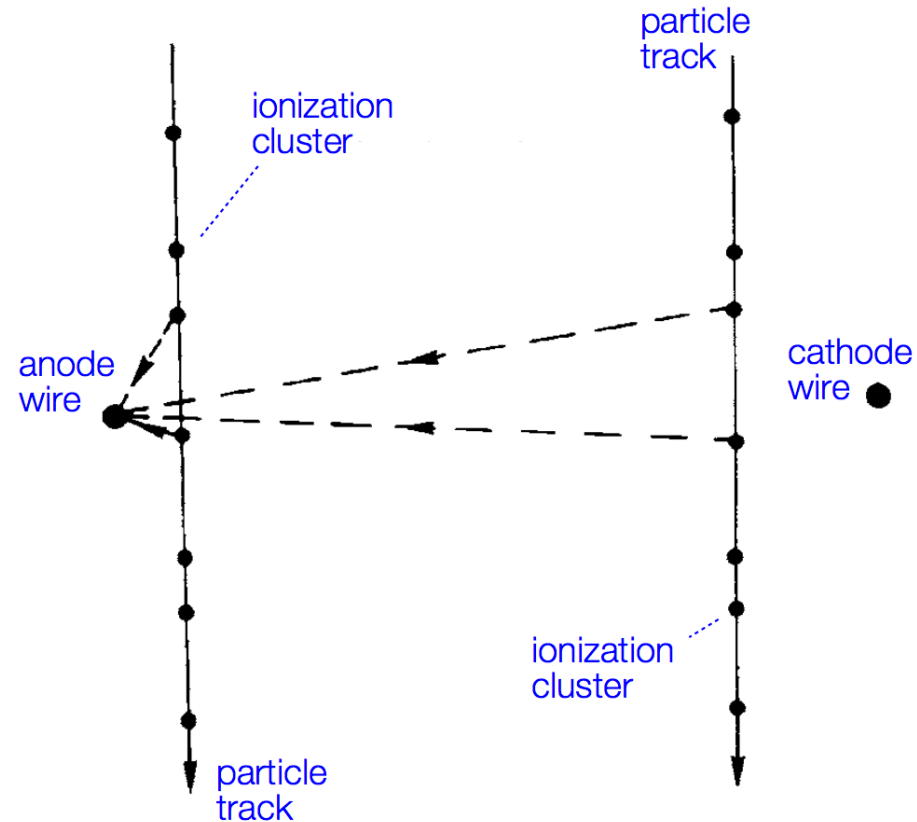
Electronics [$\sigma_{\text{electronics}} = \text{const.}$]

contribution also independent of drift length ...

Primary ionization statistics [$\sigma_{\text{prim}} = 1/x$]

Spatial fluctuations of charge-carrier production result in
large drift-path differences for particle trajectories close to the anode ...

[minor influence for tracks far away from anode]



Primary ionization statistics:

Step 1: Consider a track passing through an anode wire ...

Probability of no ionization within distance d:

$$P_0(d) = e^{-2Nd}$$

with
N: number of ionizations per unit length

Average minimum distance of closest ionization cluster:

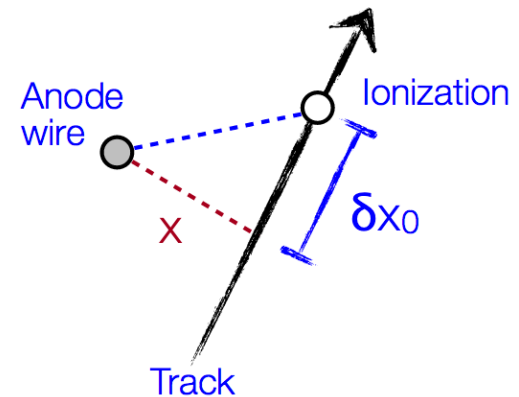
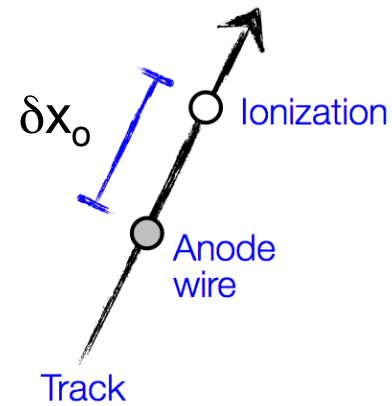
$$\delta x_0 = \langle d_{\min} \rangle = \int_0^{\infty} x e^{-2Nx} 2N dx = \frac{1}{2N}$$

$$\sigma_{\langle d_{\min} \rangle}^2 = \int_0^{\infty} \left(x - \frac{1}{2N}\right)^2 e^{-2Nx} 2N dx = \frac{1}{4N^2}$$

Normalization

Step 2: Track at distance x ...

$$\delta x = \sqrt{x^2 + (\delta x_0)^2} - x = x \left(\sqrt{1 + \left(\frac{\delta x_0}{x}\right)^2} - 1 \right) \approx \frac{x}{2} \left(\frac{\delta x_0}{x}\right)^2 \propto \frac{1}{x}$$



Camere a deriva

Diffusione

In presenza di campo elettrico: $\sigma_x/x \approx \sqrt{\frac{2KT}{eV_0}}$ → limite intrinseco delle camere a deriva.

Il coefficiente di diffusione diminuisce aumentando il campo elettrico (con campi tipici di 1KV/cm → la risoluzione è ~100μm per una distanza di deriva di 1 cm).

Concludendo:

In presenza di campo elettrico, aumentando il campo diminuisce la diffusione → usare gas “freddi” (CO₂) che hanno elettroni termici anche con alti E → bassa moltiplicazione (male) e lungo tempo di deriva (bene).

Gas “caldi” (Argon) hanno elettroni non termici anche per bassi E → diffusione anisotropa ed in genere $D_T > D_L$

In presenza di un campo magnetico la diffusione lungo B non cambia, mentre nella proiezione ortogonale a B gli elettroni fanno archi di cerchio con raggi v_T/ω → la diffusione diminuisce.

$$\sigma_x^2 = \underbrace{\left(\frac{1}{64N^2} \right) \cdot \frac{1}{x^2}}_{1^{\text{st}} \text{ ionization statistics}} + \underbrace{\frac{2D}{v_d} \cdot x}_{\text{diffusion}} + \underbrace{\sigma_{\text{const}}^2}_{\text{electronics } \delta\text{-electrons}}$$

Possible improvements:

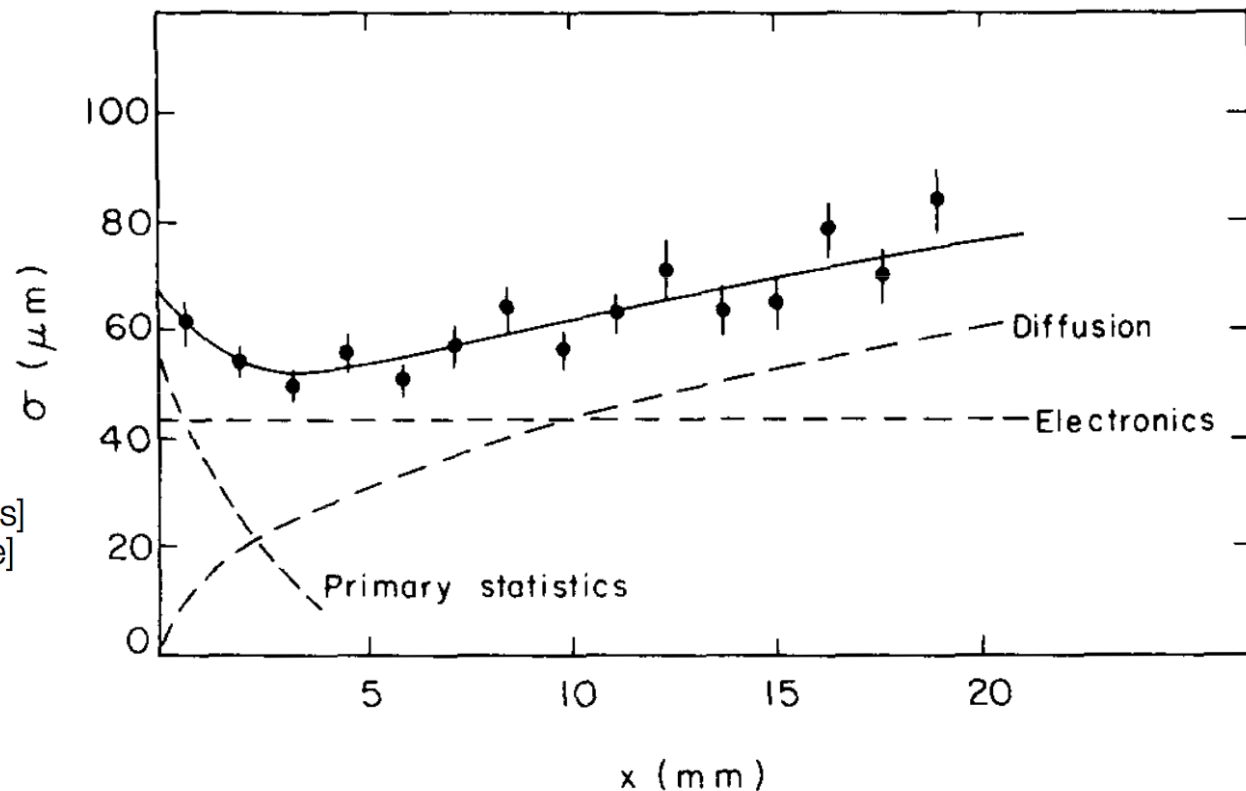
Increase N by increasing pressure ...

Decrease D by increasing pressure ...

$$D \sim \frac{\lambda_0^2}{\tau} \sim \frac{1/n^2}{1/n} \sim \frac{1}{n}$$

[n: particle density in gas]
[increases with pressure]

i.e.: increase pressure ...
[up to 4 atm possible]



La risoluzione non è determinata direttamente dalla spaziatura dei fili → **meno fili, meno elettronica, meno strutture di sostegno rispetto alle MWPC.**

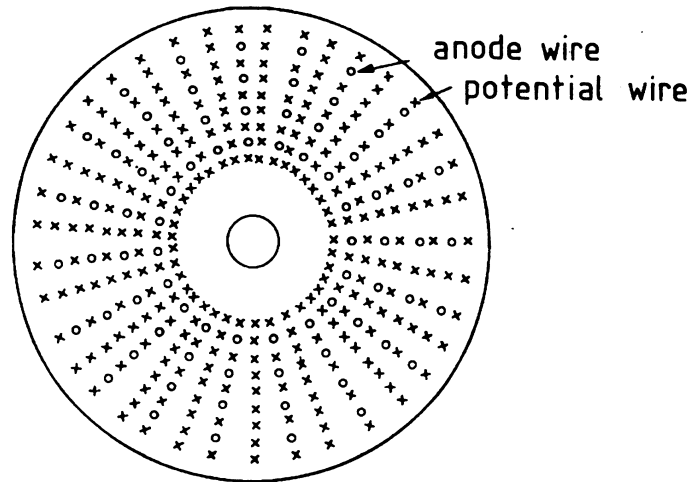
Drift Chambers

Drift chambers come in all sizes, shapes and geometries:

planar $\Rightarrow$ **fixed target** **cylindrical** $\Rightarrow$ **colliding beam**

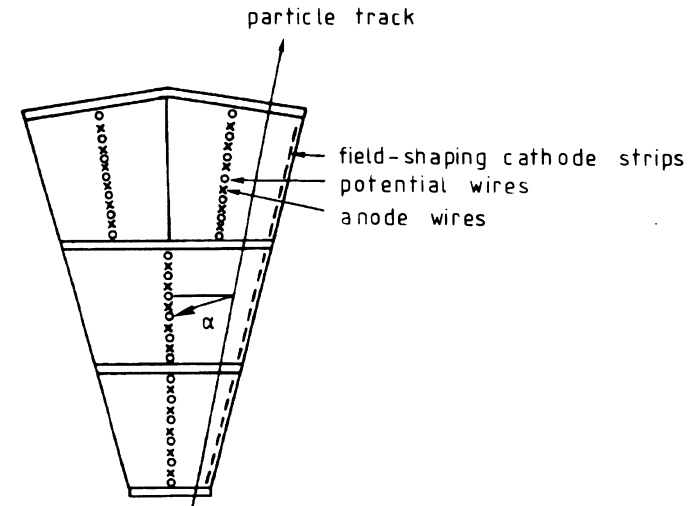
Time info gives a “circle” of constant distance around the sense wire (more complicated in B field)

In almost all cases, wires in different layers are staggered to resolve the left-right ambiguity



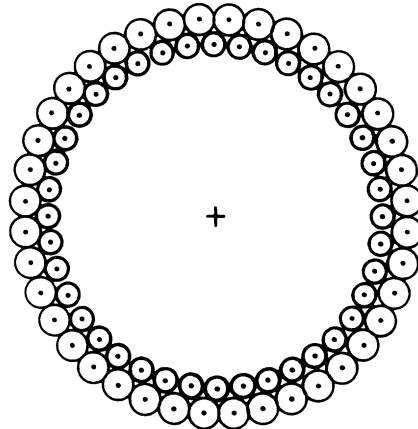
Typical cylindrical DC:

Many wires in same gas volume.
Use small angle stereo for z.
Usually use single hit electronics.
Sense (anode) and field wires.
CLEO, CDF, BELLE, BABAR



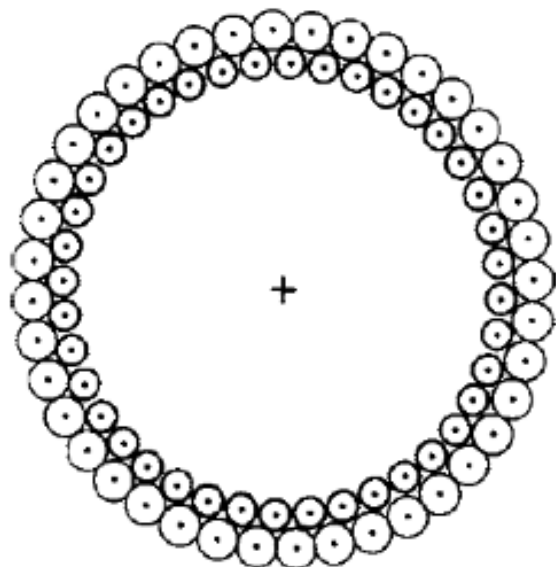
Tube Chamber:

Single sense wire in a cylinder
Can make out of very thin wall tubes.
 $\Rightarrow$ very little material
Small drift cell $\Rightarrow$ single hit electronics
Good cell isolation
 $\Rightarrow$ broken wire only affects one tube
CLEO's PTL detector



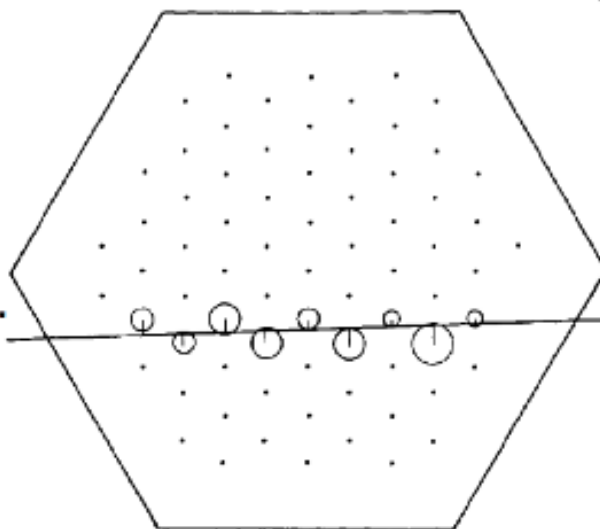
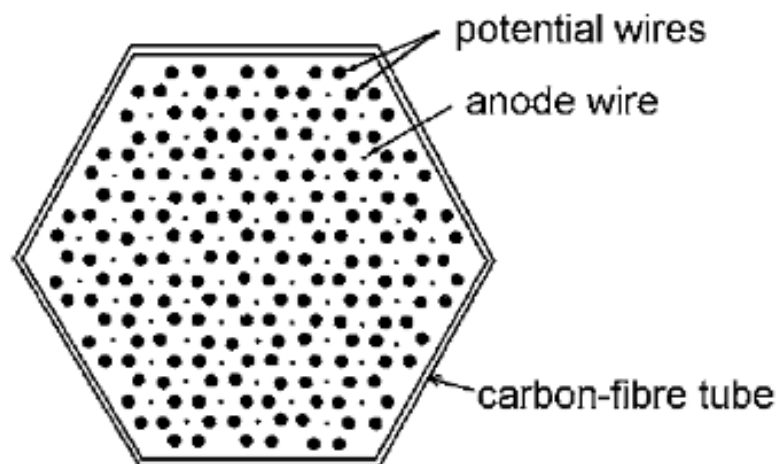
Jet chamber: optimized to resolve two tracks in a “jet”.
Drift direction roughly perpendicular to wire plane.
Single track gives multiple hits on several wires.
Use multi-hit electronics so two tracks on a wire can be resolved.
Lorentz angle must taken into account
 $\Rightarrow$ wires are “slanted”

Cylindrical configuration of thin-wall straw-tube chambers or multi-wire drift module with 70 drift cells in one carbon-fibre container of 30mm diameter



← advantage:
one broken wire =
one channel missing

one broken wire →
can kill a module



Reconstruction of a particle passage through a multi-wire drift module. The circles indicate the measured drift times of the fired anode wires. The particle track is a tangent to all drift circles

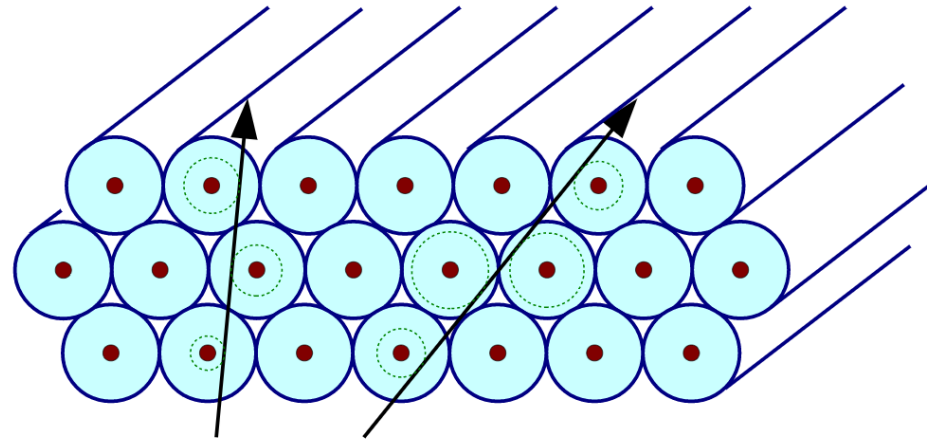
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The straw tubes have diameters of 5-10mm and are frequently operated at overpressure. These detectors allow for spatial resolutions of $30\mu\text{m}$

Straw Tubes

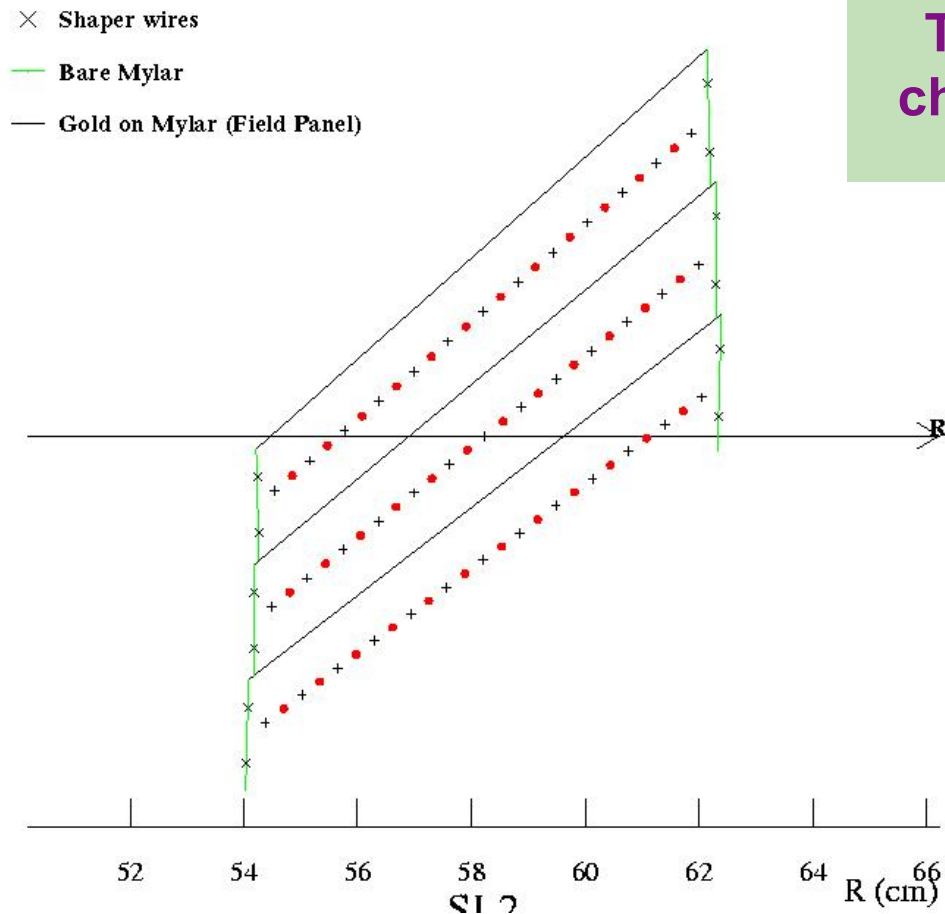
Densely packed array of individual cylindrical drift tubes

- well-defined electric field
 - no low-field regions
- reduced risk if a wire breaks
 - only one cell affected, not the whole detector
- assembly intrinsically self-supporting
 - no need for massive support structure
 - (typical wire tension 50 g $\Rightarrow$ force of 1 ton on end plates of drift chamber with 20 000 wires)
- production: "shoot" wire through each straw using pressurized air (or use vacuum suction)
- employed e.g. in ATLAS, LHCb



Large Area Drift Chambers

The “open cell” drift chamber uses field and sense wires: **field wires** create shape of electric field, **sense wires** detect time of arrival of pulse.

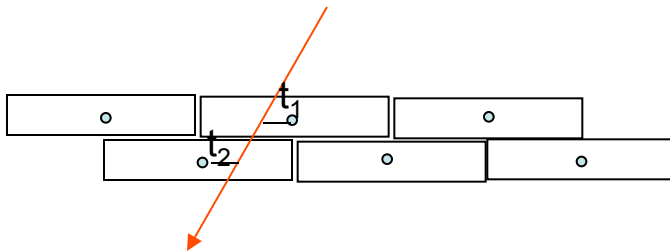


This is the design of the CDF drift chamber - the cells are tilted to take into account $E \times B$ drift!



Camere a deriva

Ambiguità destra sinistra: la misura di tempo non può discriminare fra destra e sinistra →
staggering



$t_1 + t_2 =$ tempo per percorrere $\frac{1}{2}$ cella

Se la cella è di ± 5 cm ($50 \text{ mm}/\mu\text{s}$) →

$t_1 + t_2 = 1 \mu\text{s}$.

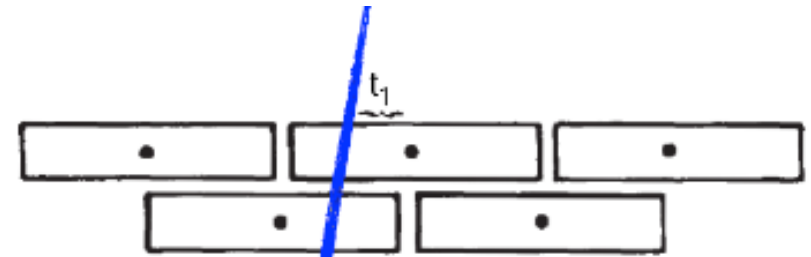
Difficulty:

Time measurement cannot distinguish whether particle has passed right or left from a wire ...

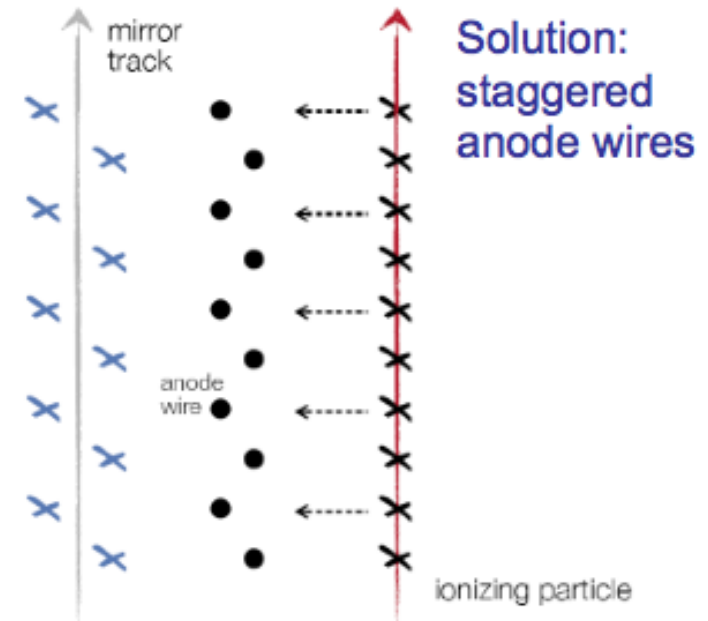
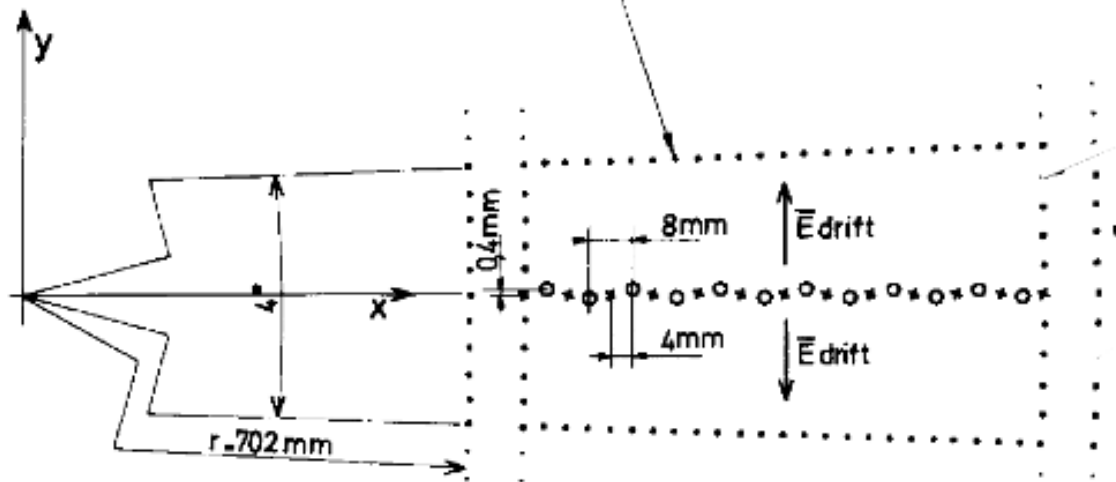
"Left-Right Ambiguity"

Solution: "Staggered wires"

Use multiple (two) layers displaced relative to each other ...



Staggered wires

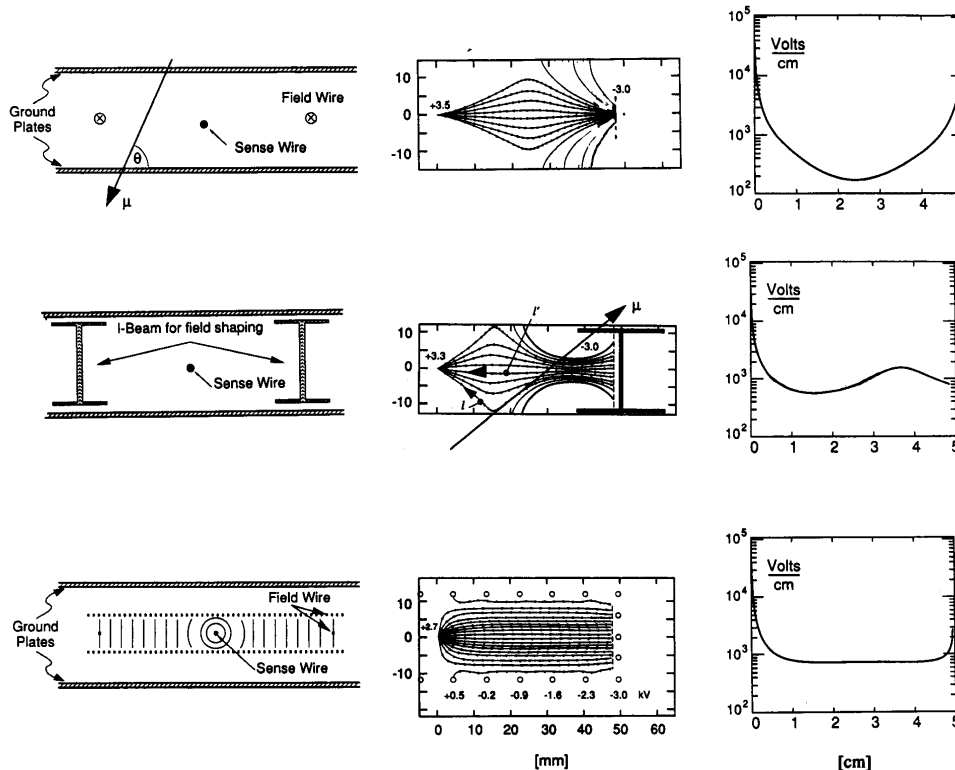


Solution:
staggered
anode wires

Camere a deriva

Camere a deriva piane.

- Ottimizzare la geometria in modo da avere E costante.
- Scegliere un gas con una velocità di deriva che dipenda poco da E → **relazione spazio-tempo lineare.**



$E(x)$ non è costante → v_D non è costante.

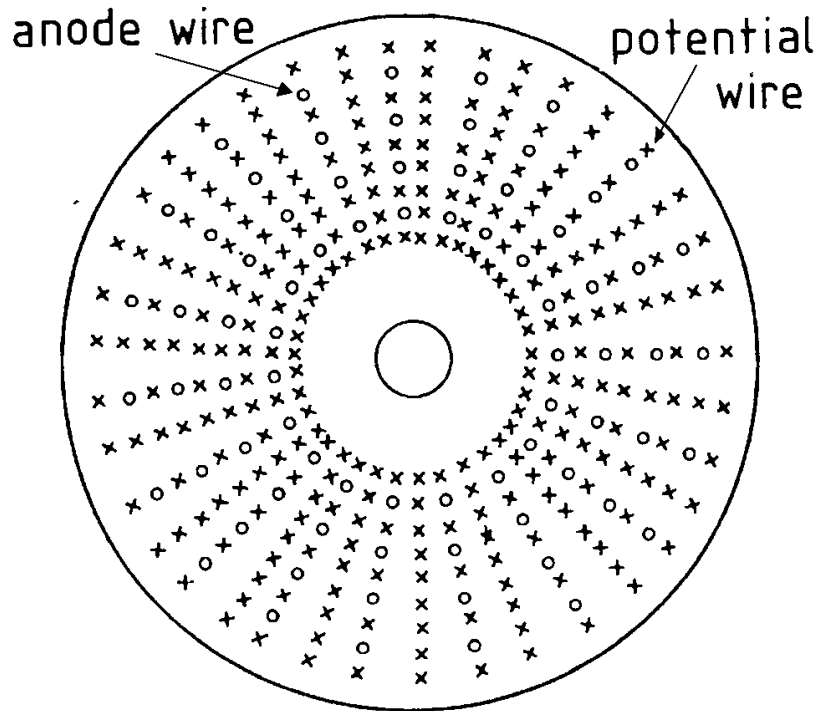
Un po' meglio della configurazione di sopra

E costante $s=s_0+bt$

Camere a deriva

Camere a deriva cilindriche.

In un esperimento ad anelli di collisione conviene usare camere a deriva cilindriche, in quanto ermetiche e facile coprire un grande angolo solido.

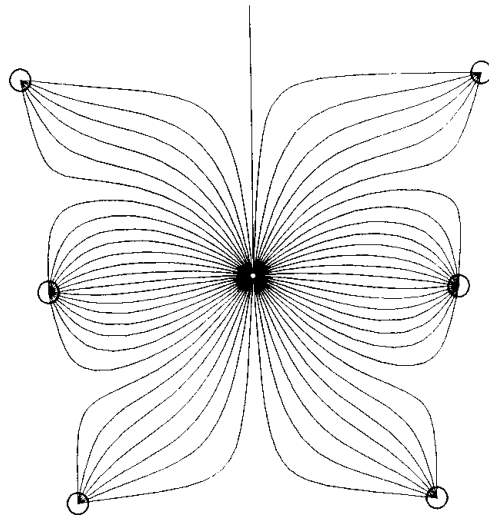
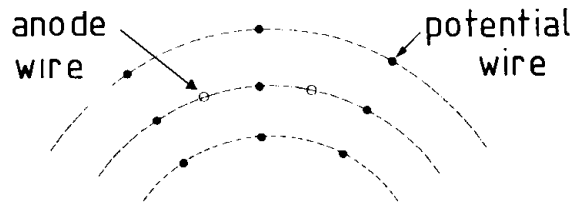


Sezione trasversale di una camera cilindrica.

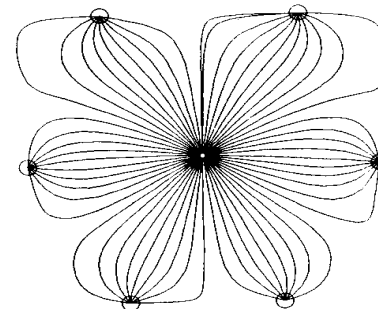
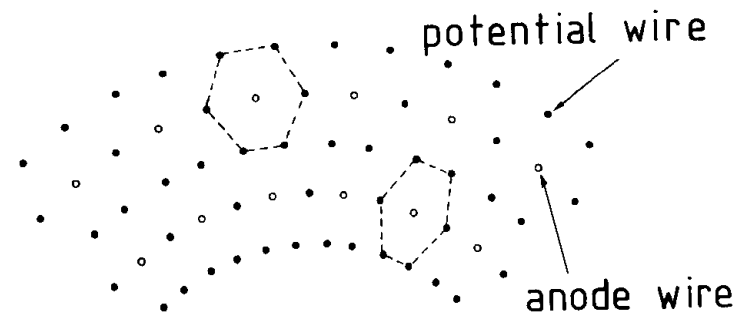
Strati di fili (**anodi, sense**) sono separati da fili di potenziale (**catodi**)

Camere a deriva

Camere a deriva cilindriche.



Configurazione più semplice che fornisce un campo non proprio al meglio



La qualità del campo può essere migliorata con una configurazione a cella chiusa, come indicato .

Camere a deriva

La misura della coordinata z (quella // al filo anodico) è normalmente misurata tramite

- ❖ divisione di carica (filo resistivo)
- ❖ misura di tempo (linee di ritardo)

Le camere a deriva cilindriche tipiche hanno 10÷15 piani di anodi → non sufficienti per identificare le particelle con misure di dE/dx

→ **camere a jet** : ottimizzate per avere il massimo numero di misure nella direzione radiale.

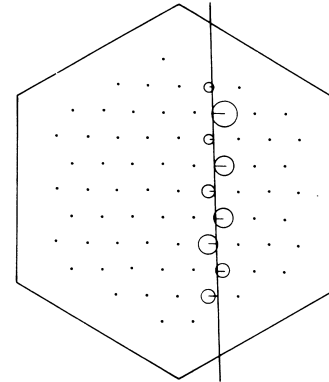
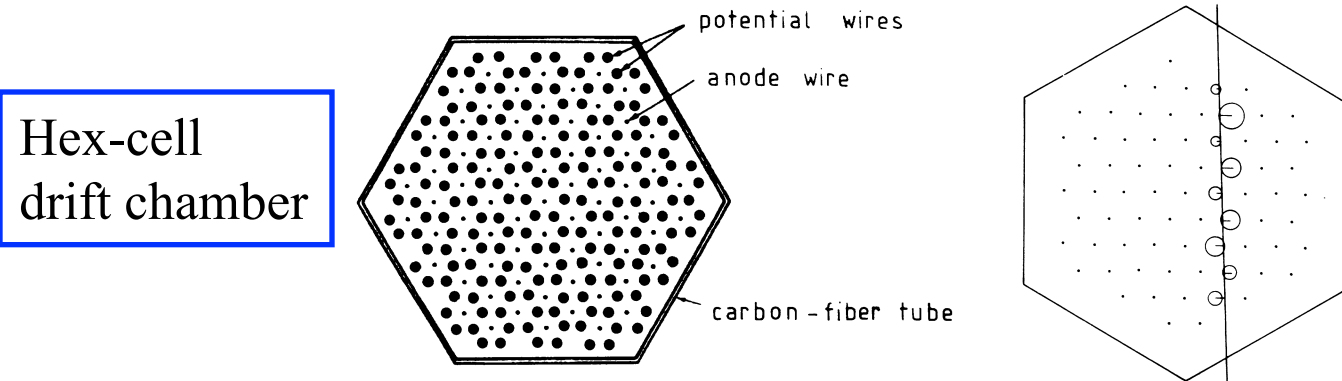
→ **camere a jet** : ottimizzato il campo per avere una velocità di deriva costante

Drift Chambers

Drift Chambers are MWPCs where the time it takes for the ions to reach the sense wire is recorded. This time info gives position info:

$$x = \int_{t_0}^{t_s} v(t) dt \quad t_0 = \text{start time}, t_s = \text{stop time} = \text{time electrons reach sense wire}$$

By using the drift time information we can improve our spatial resolution by a factor of 10 over MWPCs (1mm $\rightarrow$ 100 μm).



drift times are
circles around the
sense wires

Camere a deriva

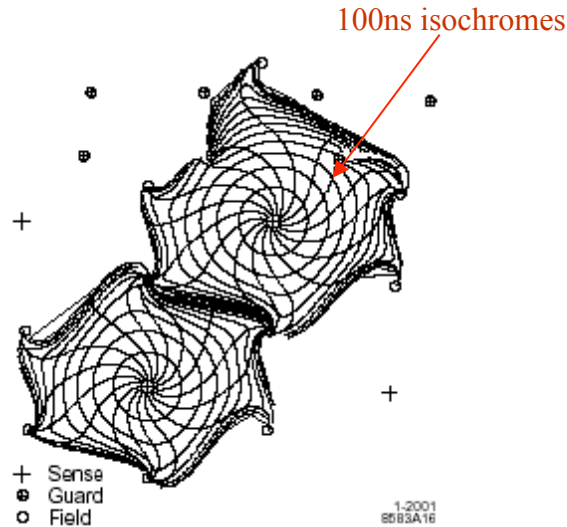
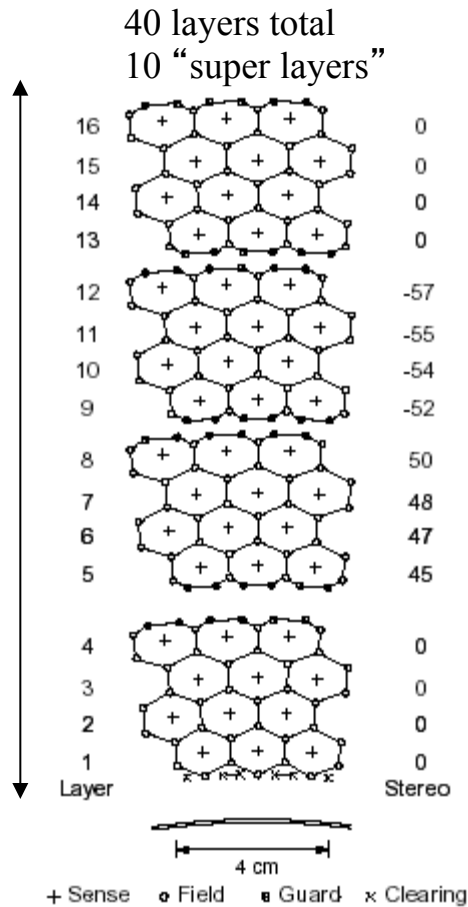
Vantaggi di una camera a deriva rispetto ad una MWPC:

- ❖ più facile da costruire da un punto di vista meccanico (fili più lontani → minori le forze elettrostatiche)
- ❖ meno fili → meno elettronica (anche se più costosa)
- ❖ migliore precisione (non più limitata alla distanza dei fili / $(12)^{1/2}$).

I parametri fondamentali sono:

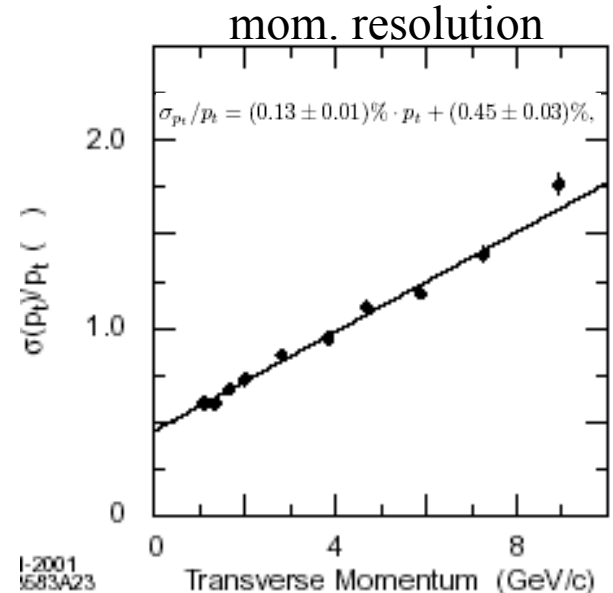
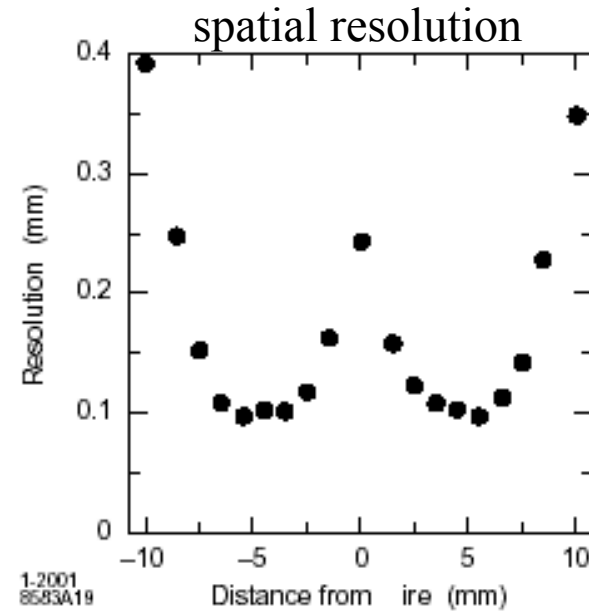
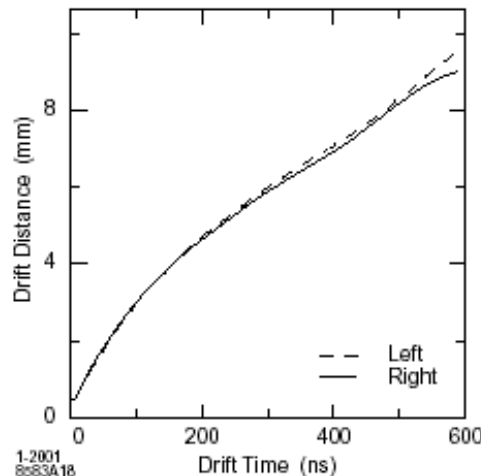
- diffusione (buono se piccola)
- velocità di deriva (ottimo se costante). Tipiche velocità di deriva (con argon-isobutano nelle proporzioni 75%-25% e campi elettrici $E \sim 700-800 \text{ V/cm}$) $\sim 50 \text{ mm}/\mu\text{s}$.

A Real Life Drift Chamber-BaBar



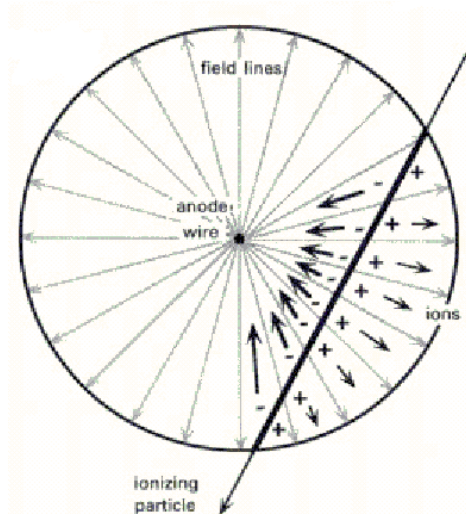
In $B=1.5T$ the ions do not drift straight to the sense wire (anode)

Time to distance relationship complicated!



7104 sense wires (20 μ m diameter)
30gm tension in each wire, sag $\sim$ 200 μ m
In order to measure "z" (along wire)
some wires are "slanted" at a slight angle
AR/Isobutane gas (80/20%)
HV $\sim$ 1950V

Example: ATLAS Muon Drift Tubes



<u>Parameter</u>	<u>Design Value</u>
Gas Mixture	Ar / N ₂ / CH ₄
	91%/ 4%/ 5%
Gas Pressure	3bar absolute
Track ionisation	330/cm
Gas gain	2×10^4
Wire potential	3270V
Electric Field at the wire	$205 \times 10^3 \text{V/cm}$
Electric Field at the wall	340V/cm
Maximum Drift time	500ns
Average drift velocity	30 $\mu\text{m/ns}$
Resolution	80 μm

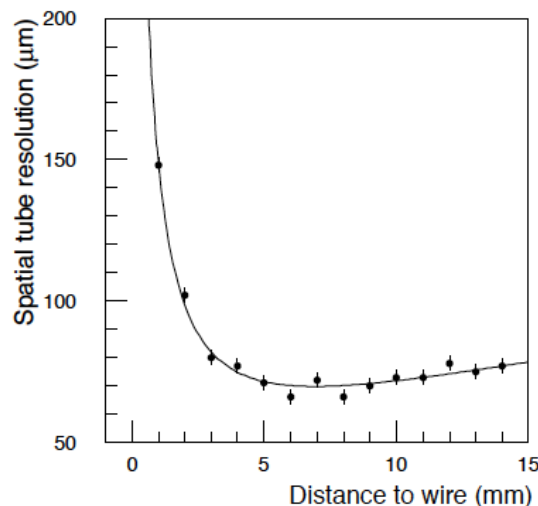


Figure 6-18 Spatial resolution of a single drift tube as a function of the track distance from the wire.

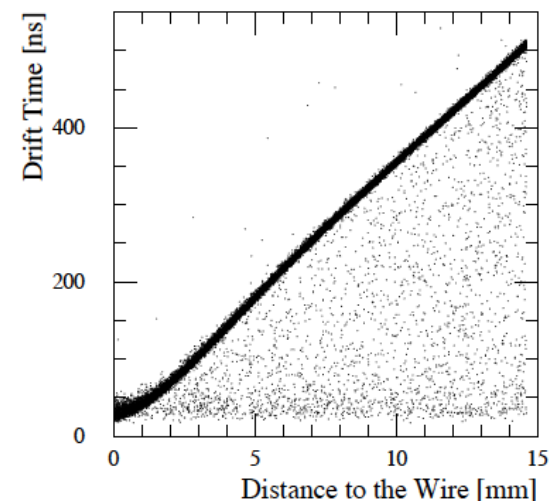
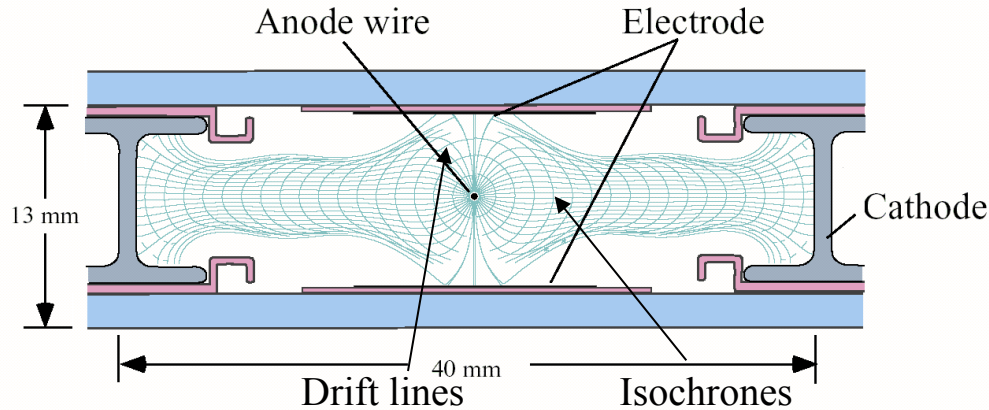


Figure 6-19 r-t relationship for the gas mixture Ar (91)/N₂ (4)/CH₄ (5). The background below the main band is due to Delta rays.

More Complex Geometry: CMS Muon DTs

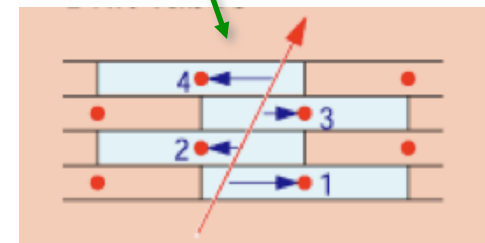
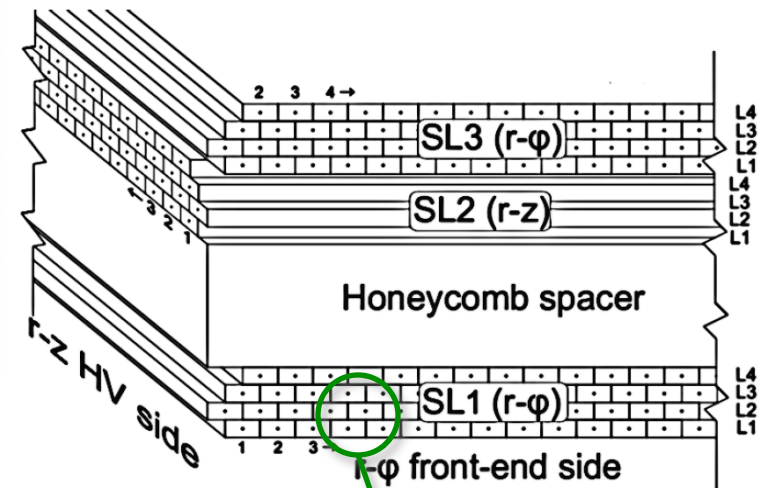
- Additional field shaping electrodes ensure a linear space-time relationship:



- Operating parameters:

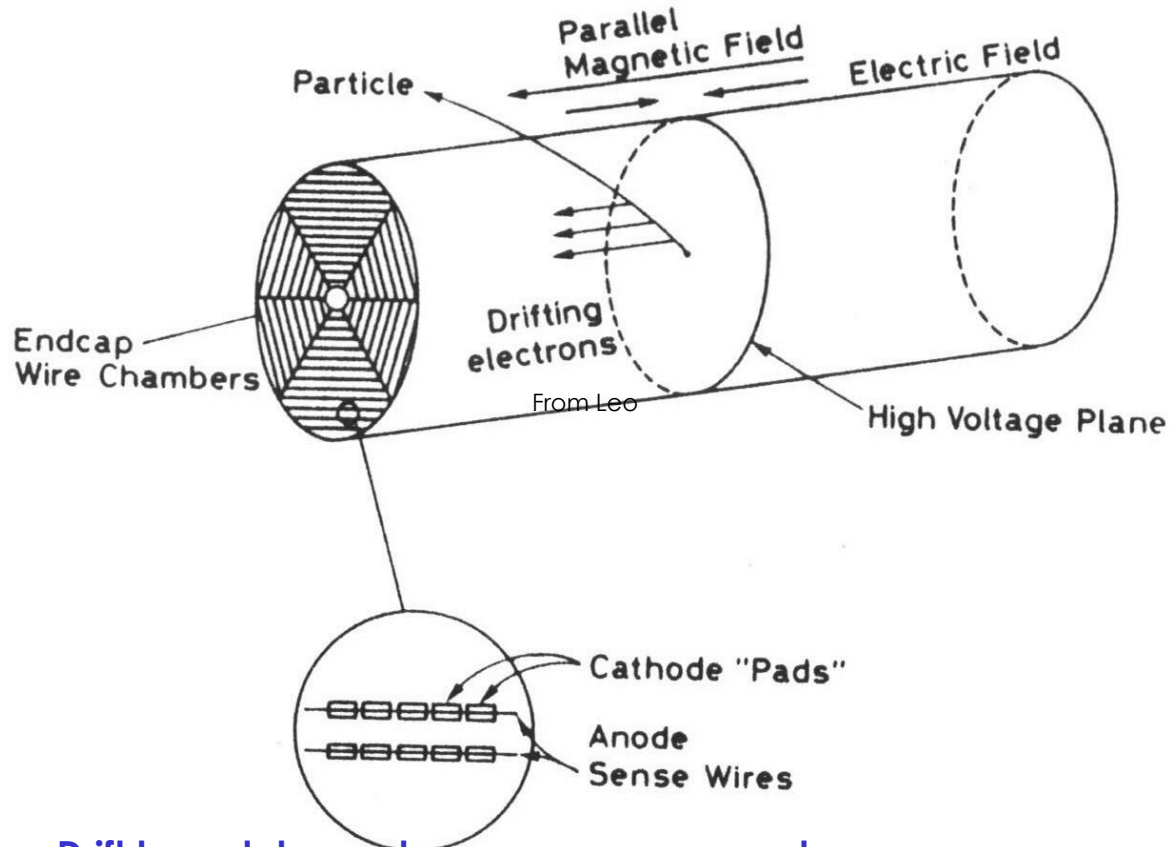
Gas Mixture	Ar/CO ₂ (85%/15%)
Wire voltage	+3600
Electrode strip voltage	+1800
Cathode strip voltage	-1200
Gain	9×10^4

- Alternating layers oriented perpendicular to each other give measurement in 2 dimensions



Time Projection Chamber

- The most sophisticated gas position detector is the Time Projection Chamber or TPC: a 3D tracking detector capable of providing information on many points of a particle track along with information on the specific energy loss, dE/dx , of the particle.
- The TPC makes use of ideas from both the MWPC and drift chamber.
- The detector is essentially a large gas-filled cylinder with a thin high voltage electrode at the center. (At high energy colliders, the diameter and length of the cylinder can be as large as two meters).
- When voltage is applied, a uniform electric field directed along the axis is created. A parallel magnetic field is also applied.
- The ends of the cylinder are covered by sector arrays of proportional anode wires arranged as shown. Parallel to each wire is a cathode strip cut up into rectangular segments. These segments are also known as cathode pads.



Drift to endplane where x,y are measured

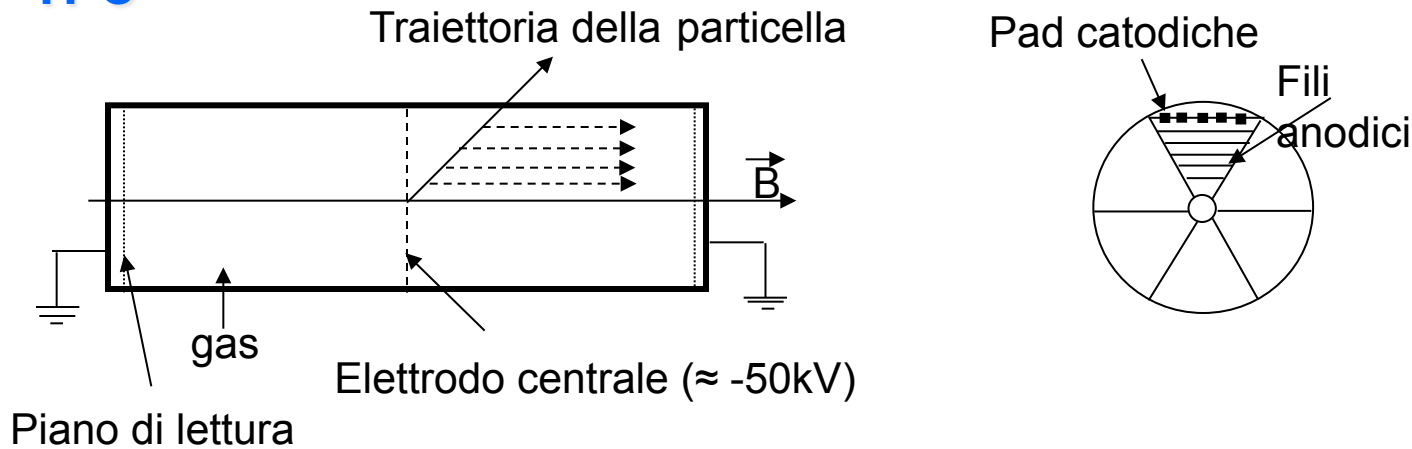
Drift-time provides z

Analogue readout provide dE/dx

Magnetic field provide p (and reduce transverse diffusion during drift)

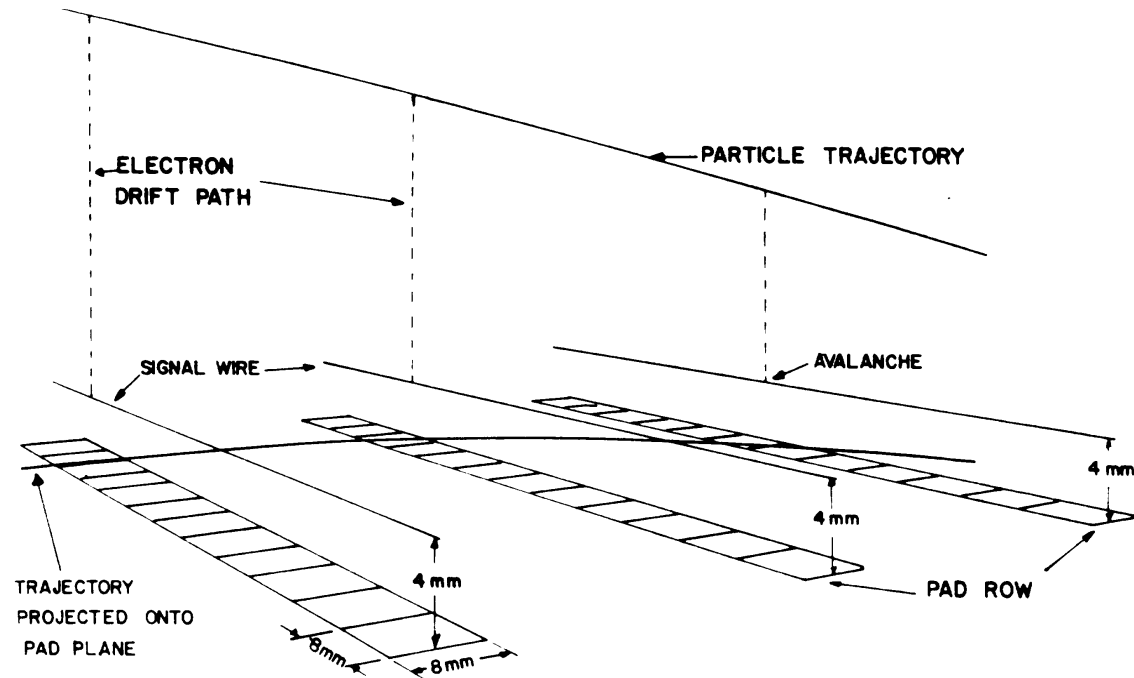
TPC

TPC

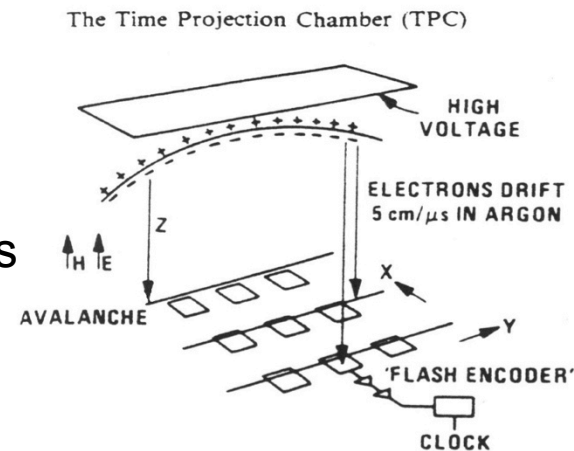


- ❖ La camera è divisa in 2 metà tramite un elettrodo centrale
- ❖ Gli elettroni di ionizzazione primaria si muovono nel campo elettrico verso le placche finali della camera (normalmente delle MWPC). Campo magnetico $\parallel$ al campo elettrico. La diffusione ortogonale al campo è soppressa dal campo B .
- ❖ Il tempo di arrivo degli elettroni sulle placche finali fornisce la coordinata lungo l'asse del cilindro (z). La moltiplicazione degli elettroni avviene vicino agli anodi. x e y si ottengono dagli anodi e dal catodo della MWPC suddiviso normalmente in pad.

Time Projection Chamber

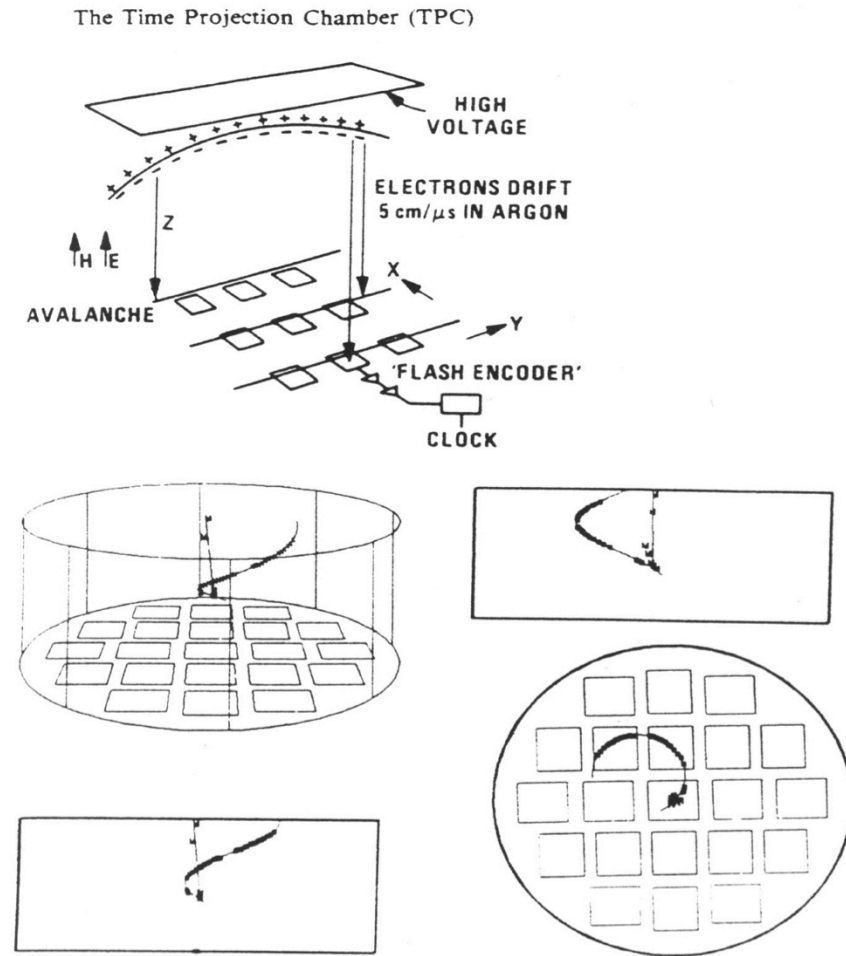


- At a colliders, the detector is positioned so that its center is at the interaction point. The TPC thus subtends a solid angle close to 4π . Particles from IP pass through the cylinder volume producing free e- which drift to the endcaps where they are detected by the anode wires as in a MWPC.
- This yields the position of a space point projected onto the endcap plane. One coordinate is given by the position of the firing anode wire while the 2nd is obtained from the signals induced on the row of cathode pads along the anode wire. Using the center-of-gravity method, this locates the position of the avalanche along the firing anode wire.
- The third coordinate, along the cylinder axis, is given by the drift time of the ionization e-. Since all ionization electrons created in the sensitive volume of the TPC will drift towards the endcap, each anode wire over which the particle trajectory crosses will sample that portion of the track. This yields many space points for each track allowing a full reconstruction of the particle trajectory.



TPC

- Because of the relatively long drift distance, diffusion, particularly in the lateral direction, becomes a problem.
- B field confines the e^- to helical trajectories about the drift direction. This reduces diffusion by as much as a factor of 10.
- In order to avoid deviating the trajectories of the drifting e^- , the B and E fields must be in perfect alignment and uniform over the volume of the drift zone down to about one part in 10^4 .



TPC

La **TPC** permette di determinare un punto nello spazio (x, y, z ovvero r, ϕ, z).

Il segnale analogico sull'anodo fornisce dE/dx .

$E/B \rightarrow$ angolo di Lorentz = 0 e la velocità di deriva è quindi parallela sia al campo elettrico che magnetico.

Il campo magnetico sopprime la diffusione $\perp$ al campo (Gli elettroni spiralizzano attorno a B.) Per $E \sim 50 \text{KV/m}$ e $B \sim 1.5 \text{ T} \rightarrow$ raggi di Larmor $\sim 1 \mu\text{m}$

Richieste:

- i. Per misurare bene la coordinata z bisogna conoscere perfettamente la $v_D \rightarrow$ calibrazione tramite laser e correzioni per la pressione e temperatura.
- ii. La deriva avviene su lunghe distanze $\rightarrow$ gas molto puro e sempre monitorato.
- iii. Poco materiale (solo gas) $\rightarrow$ minimizzo lo scattering multiplo e la conversione dei fotoni.

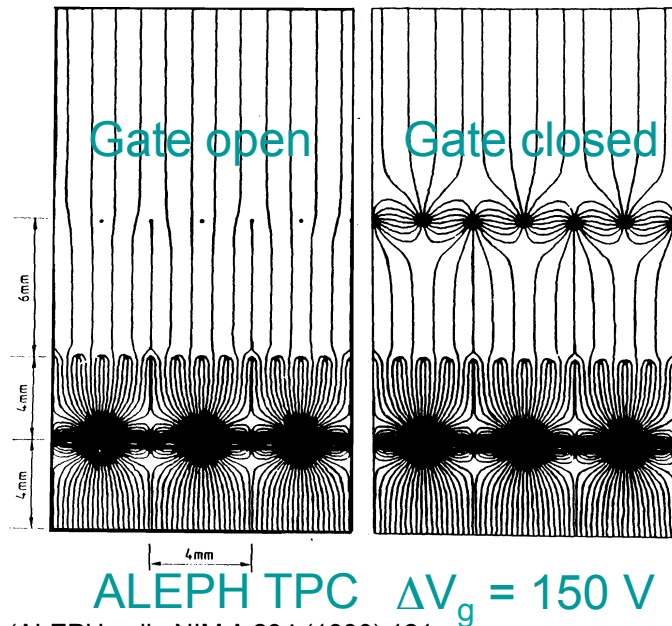
Esempi:

PEP-4 TPC $p=8.5$ atmosfere, $\text{Ar}=80\%$, $\text{CH}_4=20\%$ $V_{\text{centr}}=-55\text{kV}$ $B=1.325\text{T}$ lunga 2m e con raggio 1m.

Aleph TPC lunga 4.4 m e diametro 3.6 m, risoluzione $\sigma_{r\phi}=173\mu\text{m}$, $\sigma_z=740 \mu\text{m}$ per leptoni isolati.

TPC

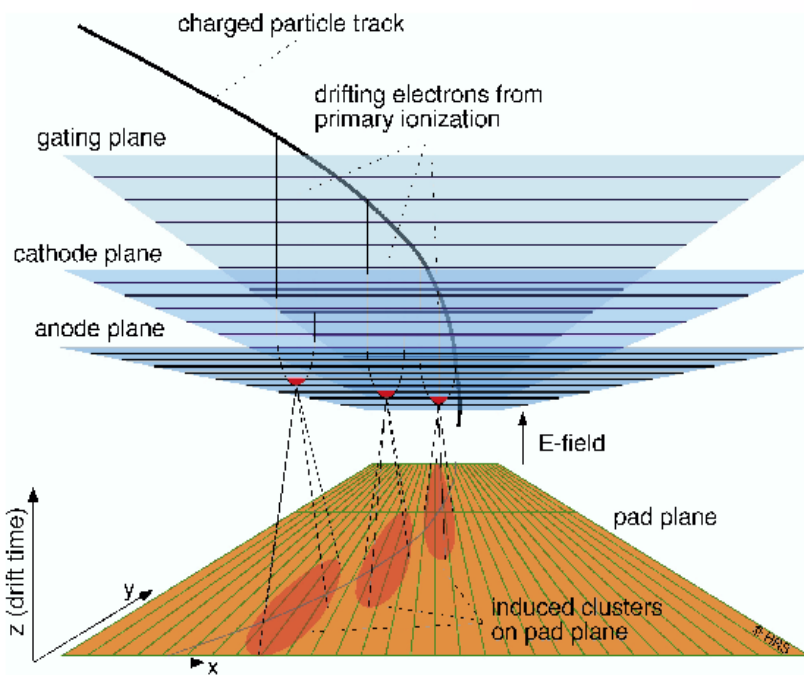
Molti ioni positivi creati nella zona di moltiplicazione vicino agli anodi della MWPC che possono andare fino all'elettrodo centrale → carica spaziale che deteriora il campo → si introduce una griglia (gate)



Il gate è normalmente chiuso, viene aperto solo per un breve tempo quando un trigger esterno segnala un evento interessante → passano gli elettroni. Viene chiuso di nuovo per impedire agli ioni di tornare verso l'elettrodo centrale.

(ALEPH coll., NIM A 294 (1990) 121,
W. Atwood et. Al, NIM A 306 (1991) 446)

TPC



- A problem which arises during operation is the accumulation of a space charge in the drift volume due to positive ions from avalanches drifting back towards the central cathode. These ions are sufficiently numerous that a distortion of the electric field in the drifting volume occurs.
- This is prevented by placing a grid at ground potential just before the anode wires. Positive ions are then captured at this grid rather than drifting back into the sensitive volume. The grid also serves to separate the drift region from the avalanche zone and allows an independent control of each.

TPC

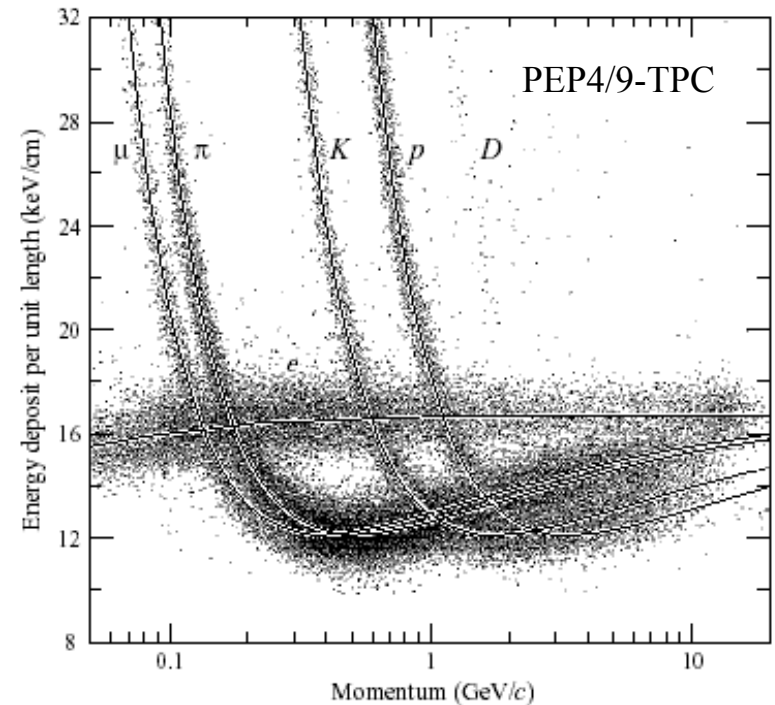
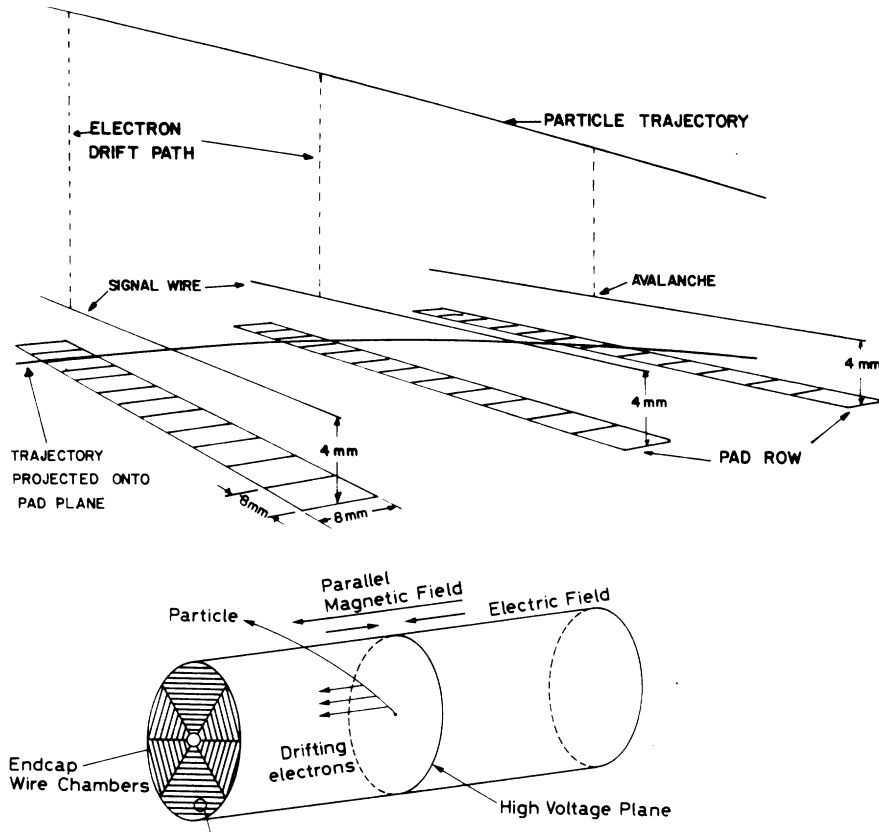
- Since the Q collected at the endcaps is $\propto$ the energy loss of the particle, the signal amplitudes from the anode also provide information on the dE/dx of the particle. If p is known from the curvature of its trajectory in the B field, for example, then this information can be used to identify the particle. In order for this method to work, however, sufficient resolution in the dE/dx measurement must be obtained. This is much more difficult to realize as many factors must be considered, e.g., electron loss due to attachment, wire gain variations in position and time, calibration of the wires, saturation effects, choice of gas and operating pressure, etc., all of which require careful thought!
- Because of the very large amount of data produced for each event, an important consideration is the readout and data acquisition system for a TPC. An approach which has been used is to use flash ADCs directly coupled to the sense wires. These ADCs are sufficiently fast such that several wires can be multiplexed into one ADC.

Time Projection Chamber

TPC measures all 3 space coordinates

$\sigma_x = \sigma_y \sim 0.1\text{-}0.2\text{ mm}$ (drift time), $\sigma_z \sim 0.2\text{-}1\text{ mm}$ (readout pad size) Used at LEP, RHIC

Many hits per track (>100) $\Rightarrow$ excellent dE/dx measurement



Drawbacks:

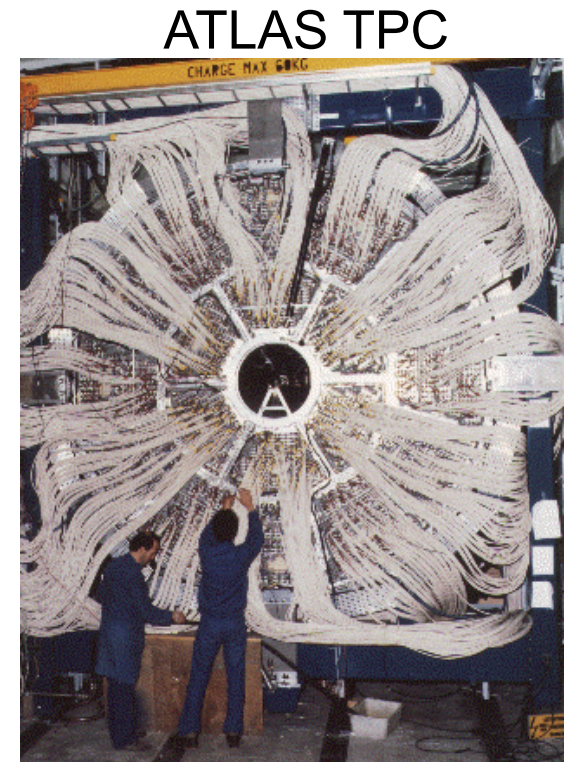
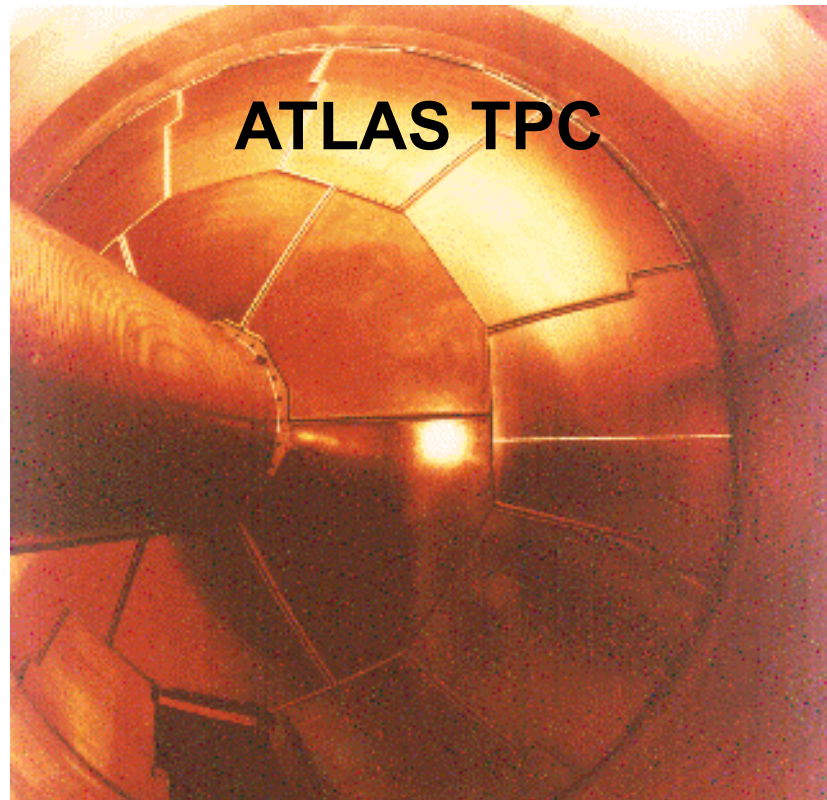
Very complicated electric field shaping: $\mathbf{E} \parallel \mathbf{B}$ to reduce effects of diffusion

Long drift times $\Rightarrow$ complicated gas system

Lots of electronic channels $\Rightarrow$ complicated electronics

Time Projection Chambers

- Note that the electric and magnetic fields are parallel and must be very homogeneous to permit accurate reconstruction. Laser “tracks” are used for calibration and alignment but extracting good calibration constants is tricky.
- Diffusion of the drifting electrons would normally smear out the measured track but the magnetic field limits this by causing the electrons to spiral in the drift direction



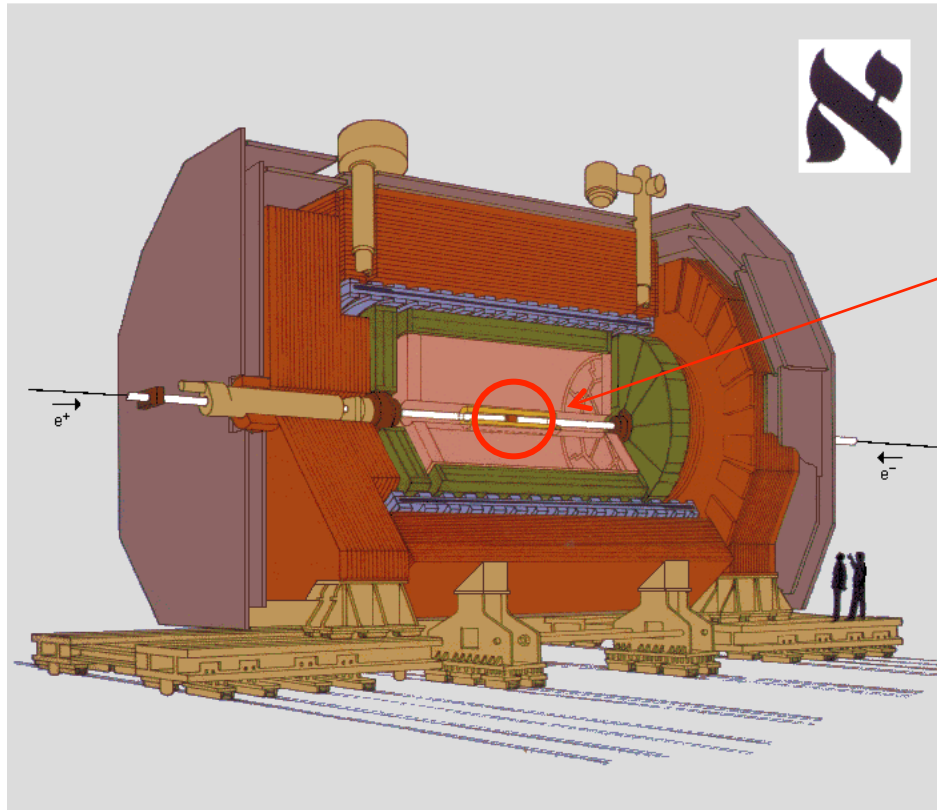
Rivelatori a stato solido

In fisica subnucleare si sono iniziati ad usare come rivelatori di posizione da meta' degli anni '80 circa.

Con camere a deriva a gas è possibile misurare posizioni con la precisione di $\sim 100 \mu\text{m}$ nel caso di traccia isolata (peggiore nel caso di tracce molto vicine) nel piano perpendicolare ai fili, mentre si hanno mm o peggio lungo la direzione del filo.

Con i rivelatori a stato solido si si possono raggiungere granularità (cioe' la segmentazione degli anodi e catodi di raccolta di carica) di pochi μm → possibile raggiungere precisioni di pochi μm richieste per la determinazione di vertici secondari e/o parametri d'impatto o per misure precise di impulso.

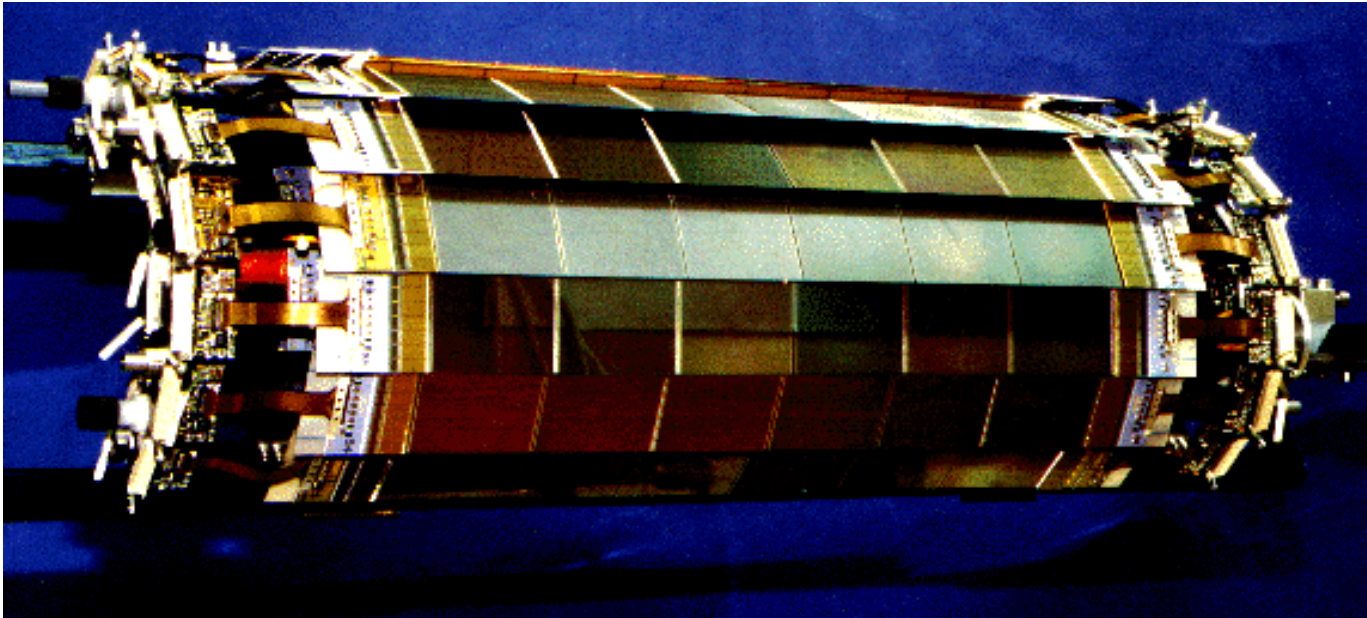
ALEPH



Gli esperimenti a LEP (Aleph, Delphi, Opal ed L3) furono i primi a essere equipaggiati con rivelatori a microstriscie di silicio.

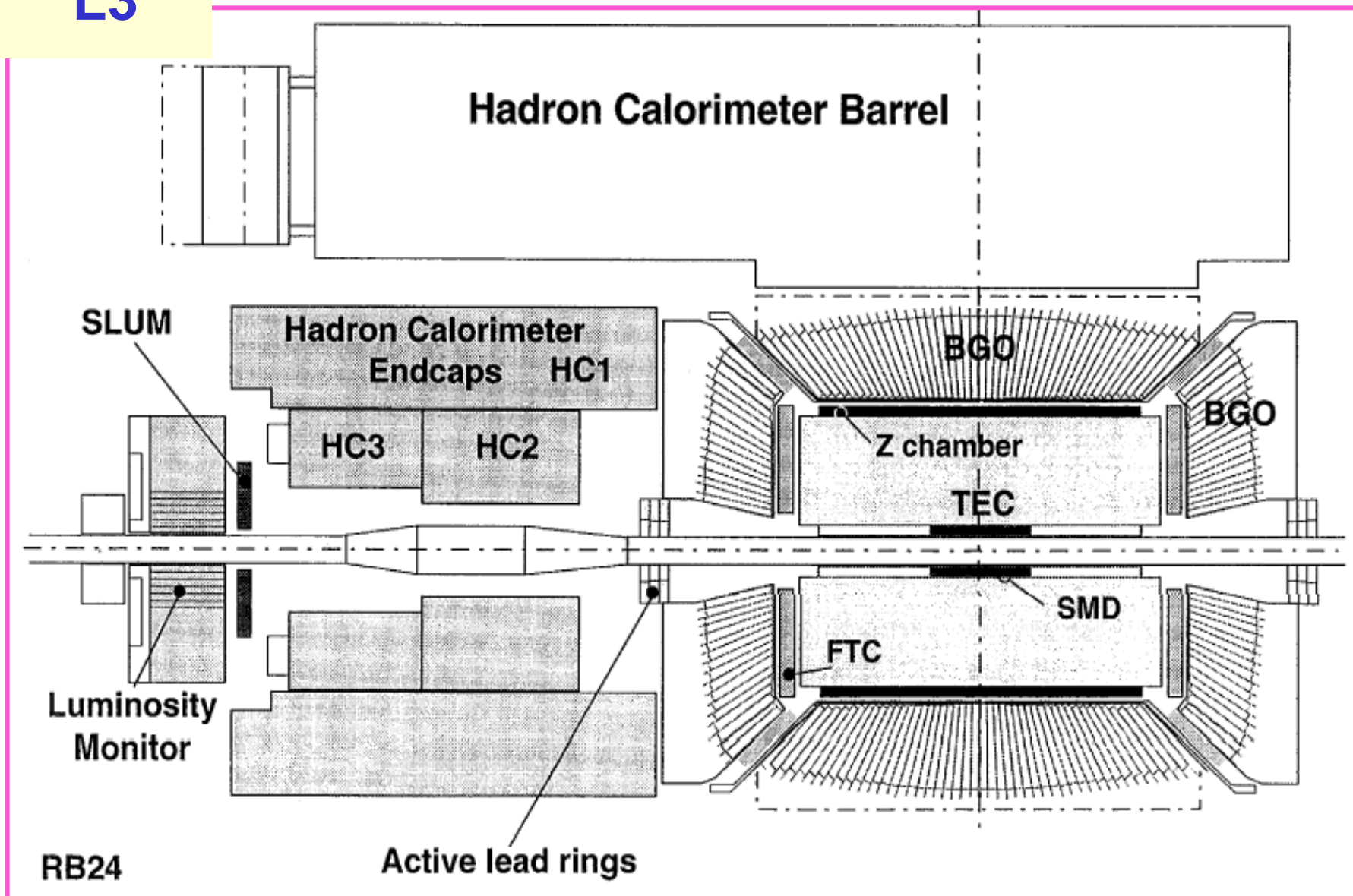
The ALEPH Detector

ALPEH – VDET (the upgrade)



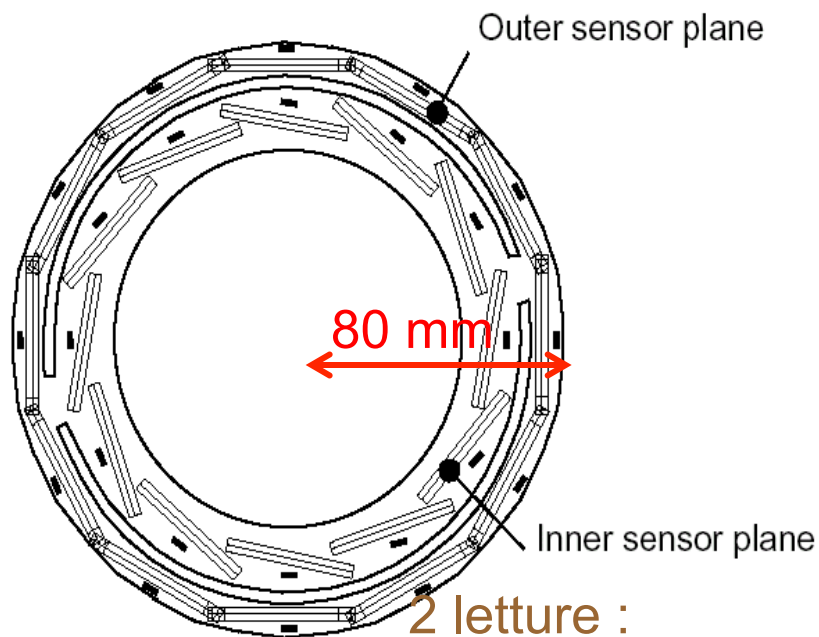
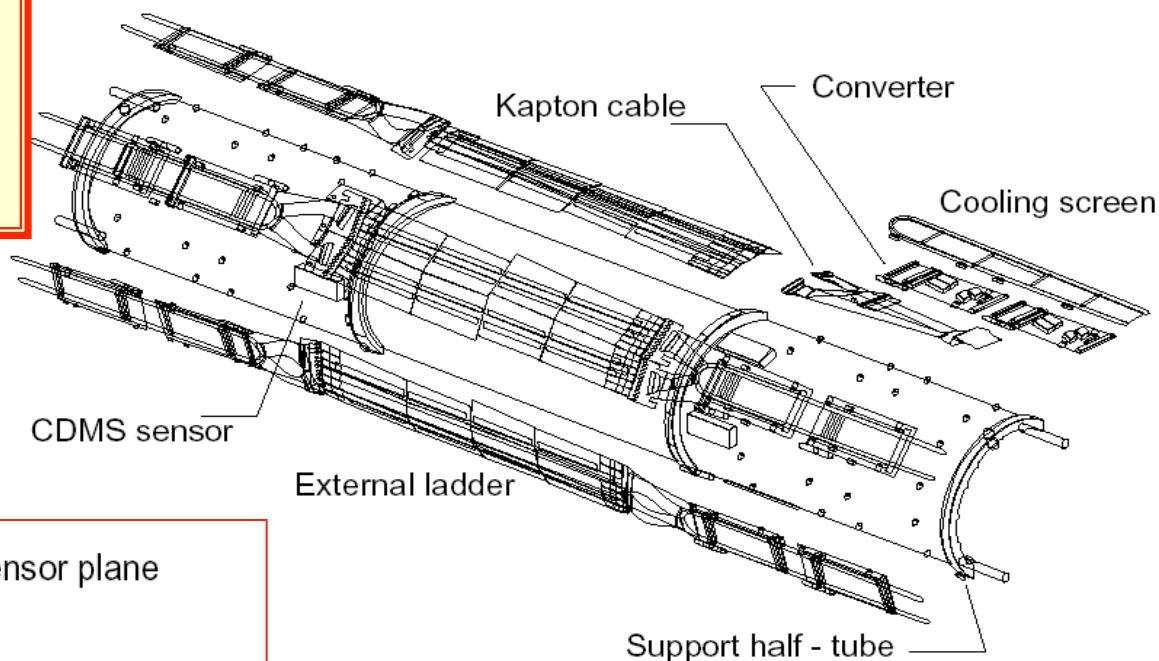
- 2 silicon layers, 40cm long, inner radius 6.3cm, outer radius 11cm
- 300 μm Silicon wafers giving thickness of only $0.015X_0$
- S/N $r\Phi = 28:1$; $z = 17:1$
- $\sigma_{r\phi} = 12\mu\text{m}$; $\sigma_z = 14\mu\text{m}$

L3



SMD : silicon μ -vertex detector - 1

96 tavolette di silicio
70 mm $\times$ 40 mm $\times$ 300 μ m



2 letture :

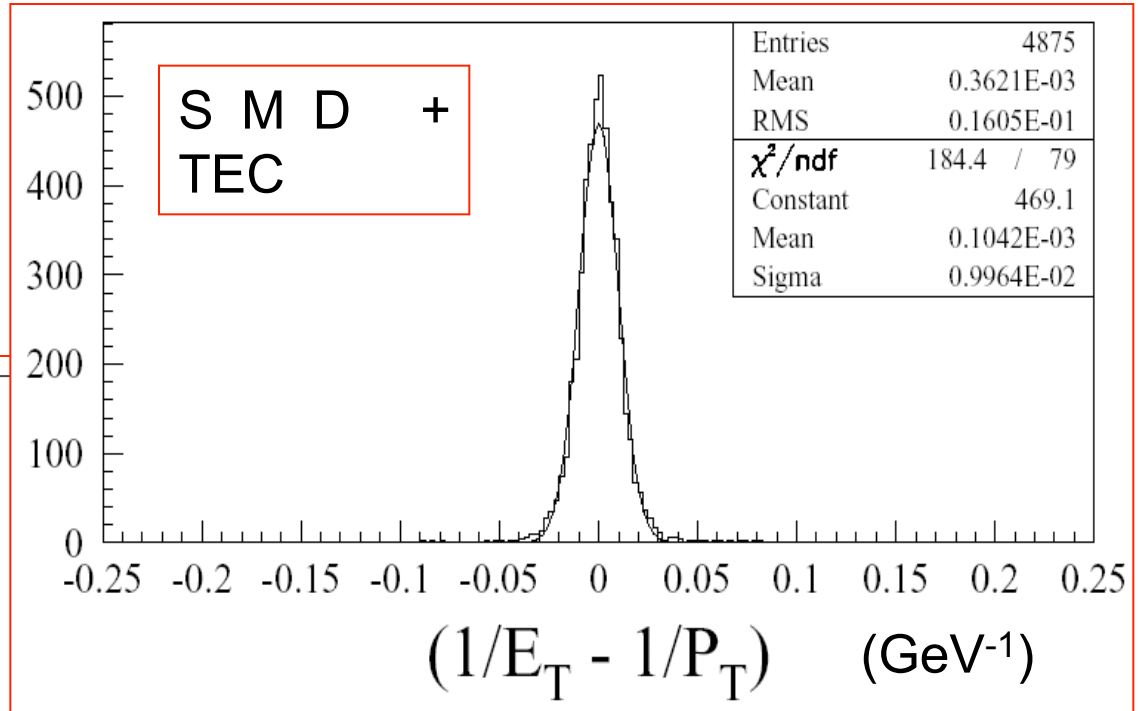
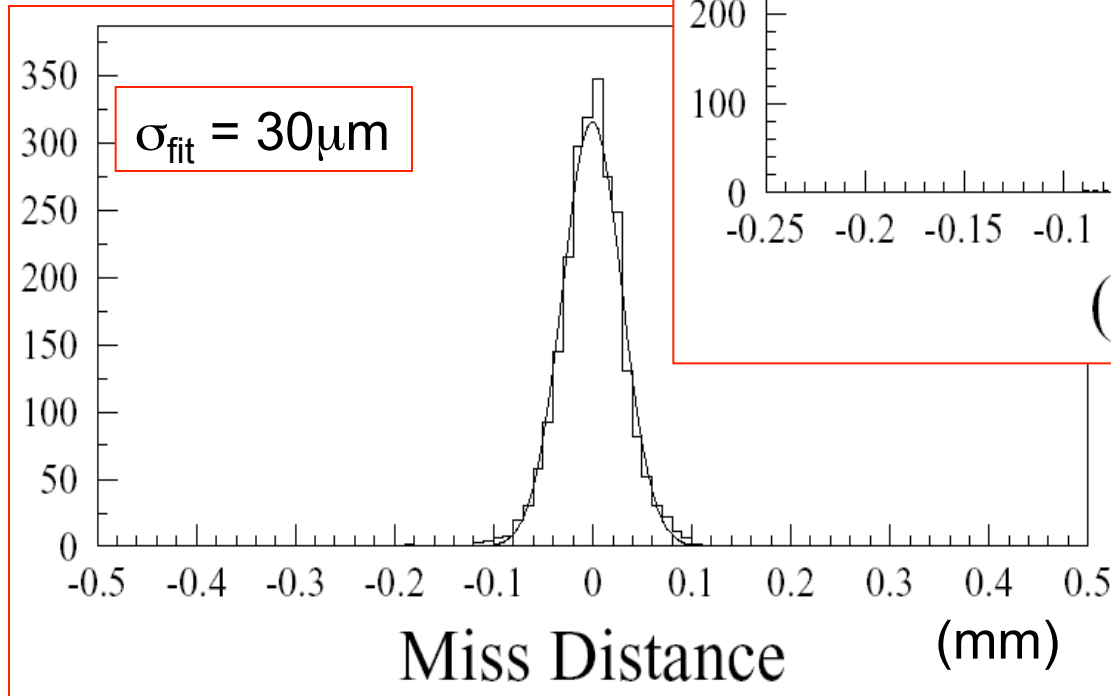
- 50 μ m in $r\phi$;
- 150÷200 μ m in z

Il primo rivelatore costruito a Pg, insieme a molti altri colleghi (bilei, santocchia, bertucci, ambrosi, servoli, pauluzzi, mantovani, alpat,...).

Tutta la tradizione "silicio" di perugia inizia da li'

Potete vederlo esposto nel pianerottolo al II piano del dipartimento

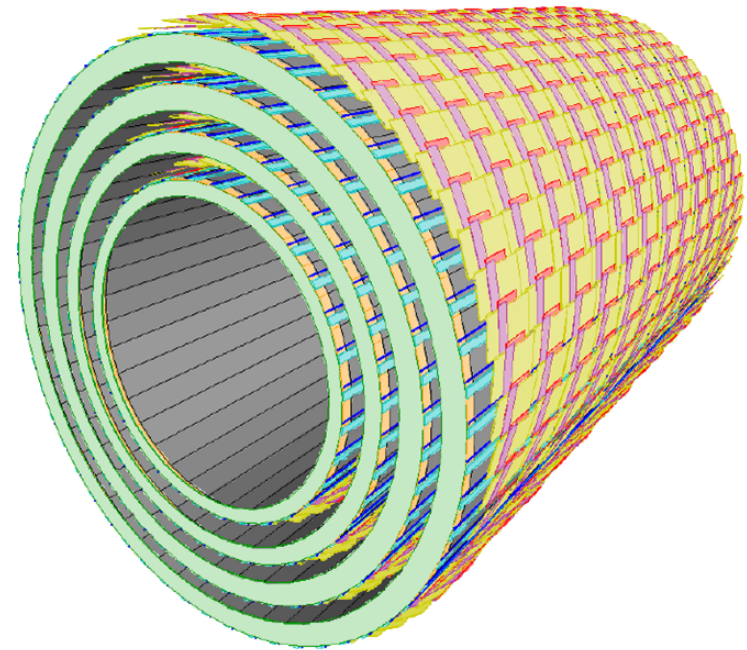
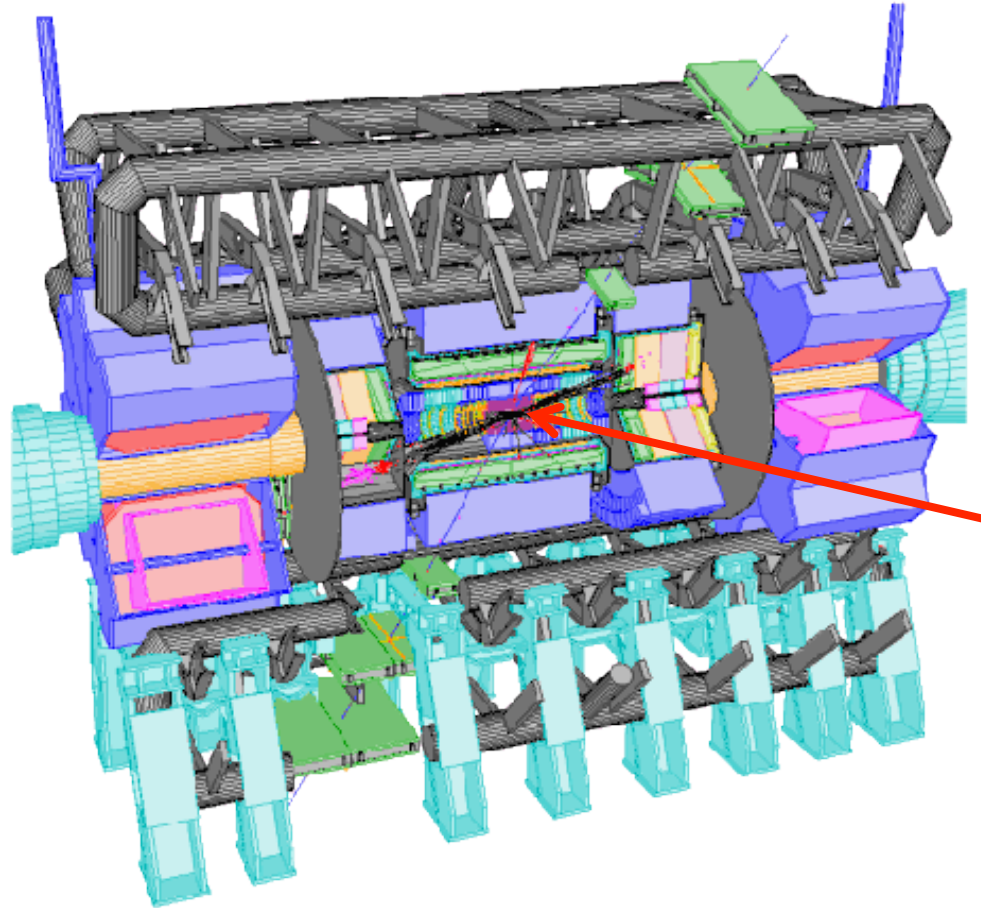
SMD : silicon μ -vertex detector - 2



per la rivelazione dei vertici secondari, vedi LHC.

ATLAS

■ A monster !



- 2112 Barrel modules mounted on 4 carbon fibre concentric Barrels, 12 in each row
- 1976 End-cap modules mounted on 9 disks at each end of the barrel region

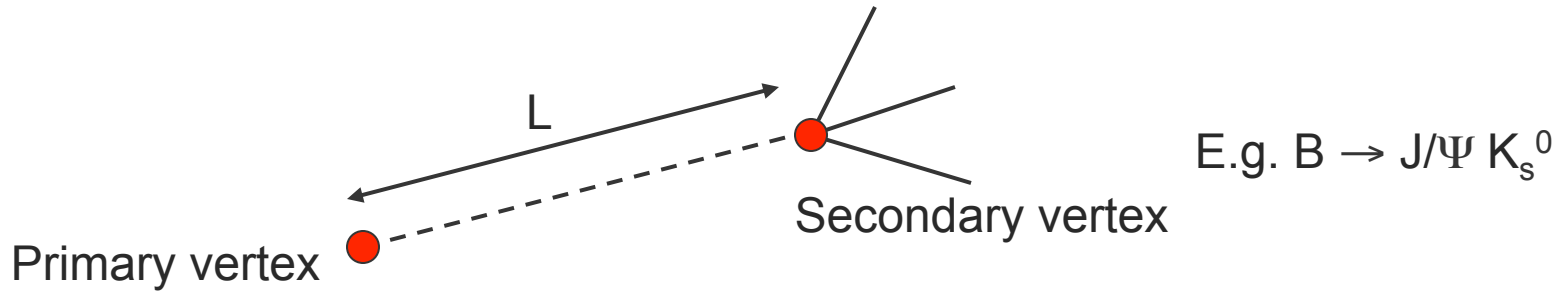
What is measured

- Measure space points of the passage of particles
- Deduce
 - Vertex location
 - Decay lengths
 - Impact parameters

Signature of Heavy Flavours

Stable particles $\tau > 10^{-6} \text{ s}$		$c\tau$
n		2.66km
μ		658m
Very long lived particles $\tau > 10^{-10} \text{ s}$		
$\pi, K^{\pm}, K_L^0$	2.6×10^{-8}	7.8m
$K_S^0, E^{\pm}, \Delta^0$	2.6×10^{-10}	7.9cm
Long lived particles $\tau > 10^{-13} \text{ s}$		
$\tau^{\pm}$	0.3×10^{-12}	91 μm
B_d^0, B_s^0, Δ_b	1.2×10^{-12}	350 μm
Short lived particles		
π^0, η^0	8.4×10^{-17}	0.025 μm
ρ, ω	4×10^{-23}	10 ⁻⁹ $\mu\text{m}!!$

Decay lengths



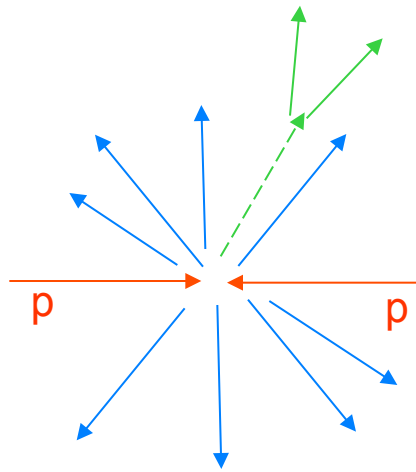
$$L = p\tau/m = E\tau/mc$$

- By measuring the decay length, L , **and** the momentum, p , or energy E **and** the mass the lifetime τ of the particle can be determined
- Need spatial accuracy on both production and decay point

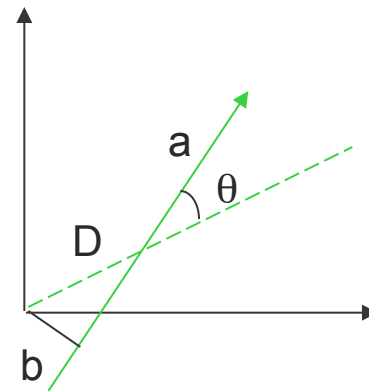
Rivelatori a stato solido

Parametro d'impatto (a un collider): distanza di min avvicinamento della estrapolazione di una traccia al vertice primario dell'interazione

A titolo di esempio consideriamo una particella neutra a breve vita media ed alta massa, prodotta ad un collider (e.g. un **D0**, **vita media $\sim 0.5-1 \times 10^{-12}$ e $c\tau = 315 \mu\text{m}$**) che decade in 2 particelle cariche (es. un K^- e un π^+), che per semplicità considereremo di massa trascurabile rispetto alla massa del D (**$m=1.85 \text{ GeV}$**). Essendo la particella in questione neutra l'unico modo di rivelarla è tramite il **parametro d'impatto**, ovvero lo spostamento del vertice secondario (di decadimento) dal vertice primario (di produzione).



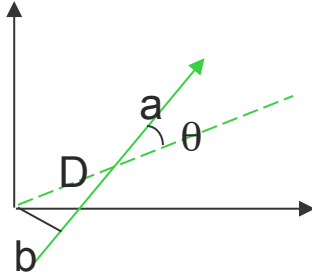
$$D \rightarrow a + x$$



D è la particella che decade in **a** che ha un parametro d'impatto **b** + un' altra (x) non indicata.

Rivelatori a stato solido

Il parametro d'impatto b è proporzionale al percorso di D e all'angolo θ . Ma:



$$D = \gamma_D (c\tau_D)$$

$$b \cong D\vartheta$$

$$\vartheta = \frac{p_T^a}{p^a} \Rightarrow b \sim \frac{p_T^a}{p^a} \gamma_D (c\tau_D)$$

Se il decadimento è a 2 corpi $D \rightarrow a + x$ e posso trascurare le masse di a e x (rispetto alla massa del D) $\rightarrow$

$$p_T^a \sim M_D/2$$

$$p^a \sim p_D/2$$

$$\vartheta \sim \frac{M_D}{p_D} \sim \frac{1}{\gamma_D}$$

ponendo $\beta=1$ e $p/M=\beta\gamma$

$\rightarrow$ il parametro d'impatto è indipendente dall'impulso $b \sim (c\tau)_D$.

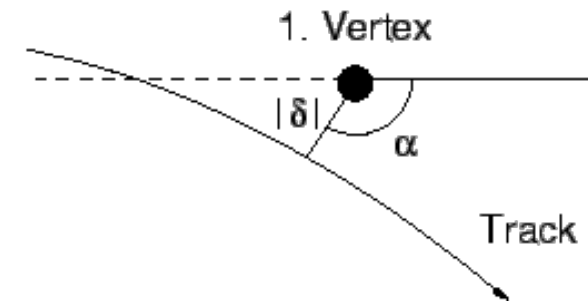
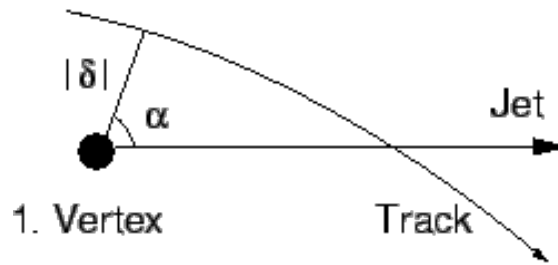
$db/b = d\tau/\tau \rightarrow$ se voglio un errore del 10% su τ , devo misurare b al 10%

Se la particella ha una vita media di $\sim 10^{-12}$ s $\rightarrow$ devo determinare distanze dell'ordine di ~ 300 μm $\rightarrow$ mi servono risoluzioni sul punto di max avvicinamento al vertice primario di ~ 30 μm $\rightarrow$ stato solido va bene

$$\alpha < \pi/2 \rightarrow \delta = +|\delta|$$

$$\alpha > \pi/2 \rightarrow \delta = -|\delta|$$

The impact parameter of a track is defined as the transverse distance of closest approach (DCA) of the track to the primary vertex point with a sign.



The signed impact parameter is defined by reference to the direction of the reconstructed particle or the jet to which the track is associated.

The signed impact parameter of a track is defined as positive if the angle between the jet direction and the line joining the primary vertex to the point of DCA $< 90^\circ$ and it is defined as negative otherwise.

The sign allows to statistically disentangle detector resolution effects and effects from the decay lifetime of the heavy hadron.

Tracks from the decays of long lived particles will have positive impact parameters, if resolution effects are neglected. Tracks produced at the primary vertex result in a symmetric distribution around 0, i.e. negative impact parameters predominantly result from detector resolution effects.

