# Lecture 14 211119

- Il pdf delle lezioni puo’ essere scaricato da
- http://www.fisgeo.unipg.it/~fiandrin/didattica_fisica/cosmic_rays1920/

# Blast shock waves

A blast shock wave is a shock wave formed by a hot gas bubble expanding supersonically in the ambient medium

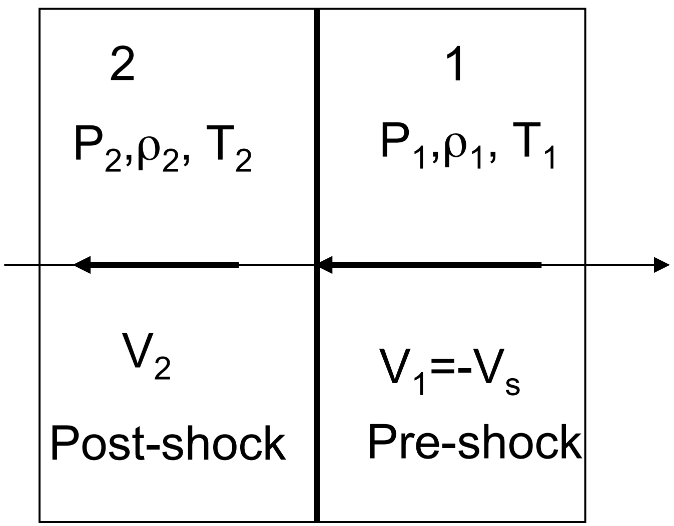
Let us assume that the expansion occurs in a uniform stationary polytropic medium with density $\rho_o$ and pressure $p_o$

how does the shock wave evolve in time?

First, since the surrounding medium is uniform, the expansion will have spherical symmetry

# Blast shock waves

We worked out the physics of the (strong) shock in the shock reference frame, where it is stationary




Shock rest frame

$$\frac{\rho_2}{\rho_1} \approx \frac{(\gamma+1)}{(\gamma-1)}$$

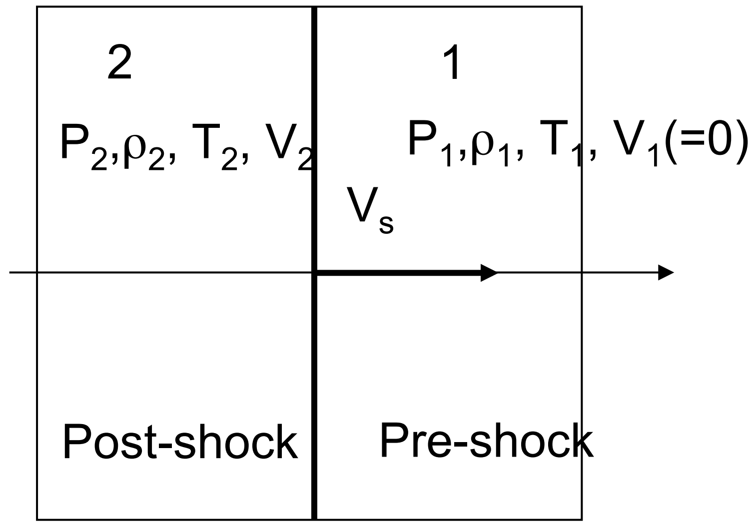
$$\frac{p_2}{p_1} \approx \frac{2\gamma M_s^2}{(\gamma+1)}$$

$$\frac{V_2}{V_1} \approx \frac{\gamma-1}{\gamma+1}$$

$$\frac{T_2}{T_1} = \frac{2\gamma(\gamma-1)}{(\gamma+1)^2} M_s^2$$

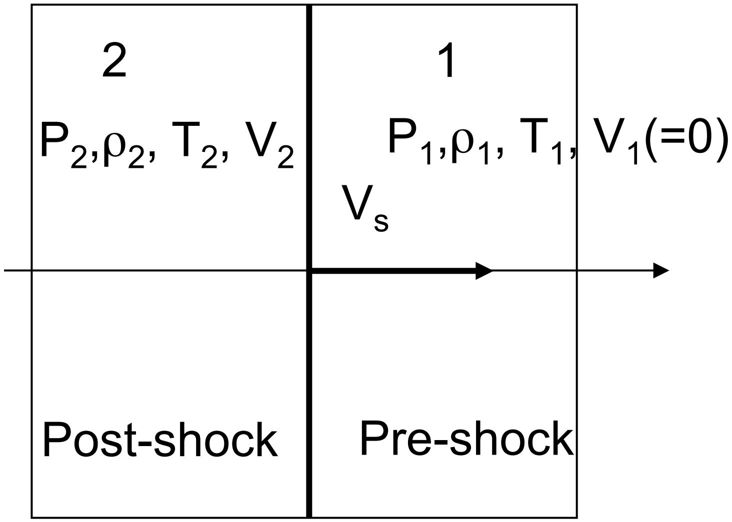
# Blast shock waves

Now we have an (supersonic) expansion in a uniform, stationary ($V_1$=0) polytropic medium with density $\rho_1 = \rho_o$ and pressure $p_1 = p_o$ and a supersonic shock wave propagating to the right (+x dir) with speed $V_s$ ahead the expanding gas, as in the case of the supersonic piston

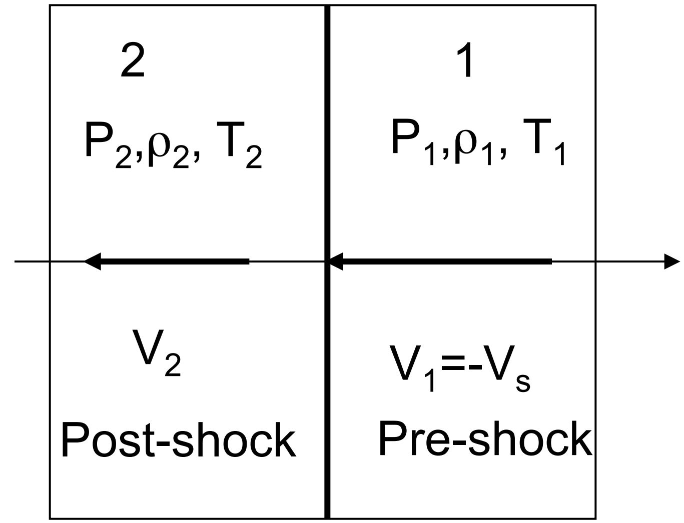



As in the supersonic piston example, the situation in a reference frame where the shock is traveling with speed $V_s$ can be obtained, for non-relativistic shocks, from a simple galileian transformation

# Blast shock waves




Observer frame




Shock rest frame

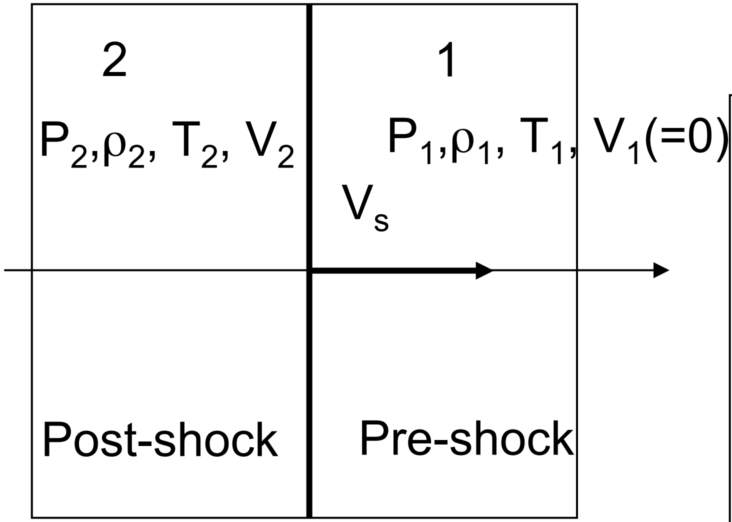
The RH conditions can be still applied, provided one interprets the speeds $V_1$ and $V_2$ as relative speeds with respect to the shock, ie apply a galileian velocity tranformation

$$\vec{V} \rightarrow \vec{V}_{rel} = \vec{V} - \vec{V}_s$$

For a shock propagating with velocity $\mathbf{V}_s$ into a medium at rest, $\mathbf{V}$=0, one has $\underline{\mathbf{V_1=-V_s}}$ and $\theta_s$=0 (as for the supersonic piston) ➔ in this case any shock is a normal shock, with $\mathbf{V}_t$=0, even when the shock surface itself is not a plane!

➔ spherical shocks are normal shocks

# Blast shock waves



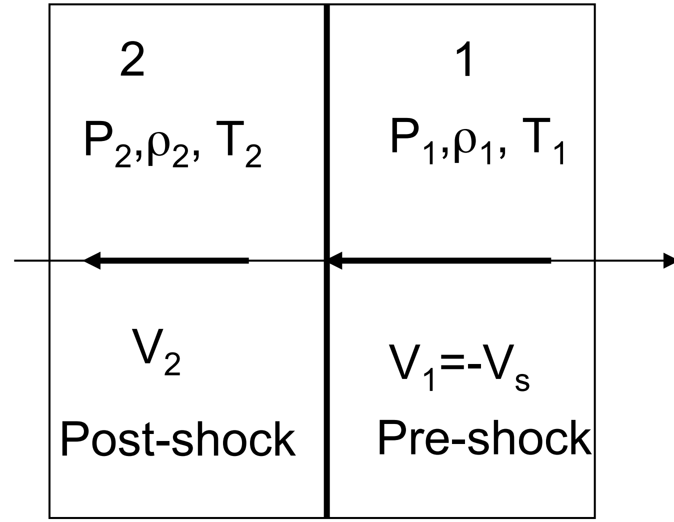

$$\vec{V}_{sh} = \vec{V}_{obs} - \vec{V}_s$$

$$\frac{\rho_2}{\rho_1} \approx \frac{(\gamma+1)}{(\gamma-1)}$$

$$\frac{p_2}{p_1} \approx \frac{2\gamma M_s^2}{(\gamma+1)}$$

$$\frac{T_2}{T_1} = \frac{2\gamma(\gamma-1)}{(\gamma+1)^2} M_s^2$$

$$\frac{V_2}{V_1} \approx \frac{\gamma-1}{\gamma+1}$$

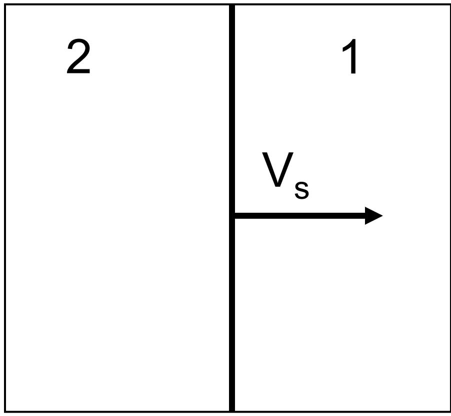



All the conditions may be directly applied, except the speed for which we need to make the substitution $V_2^{sh}=V_2^{obs} - V_s$ and $V_1=-V_s$

$$\frac{V_2}{V_1} \rightarrow \frac{V_2 - V_s}{-V_s} \approx \frac{\gamma-1}{\gamma+1} \quad \Rightarrow \quad V_2 \approx \frac{2V_s}{\gamma+1}$$

Is the post shock speed in the obs frame

# Blast shock waves




Let assume also that the explosion occurs in a uniform stationary polytropic medium with density $\rho_1 = \rho_o$ and pressure $p_1 = p_o$

The strong shock satisfies the relation $M_s^2 = (\frac{V_s}{c_s})^2 = \frac{\rho_o V_s^2}{\gamma p_o} \gg 1$

The RH relations then give

$$\frac{\rho_2}{\rho_1} \approx \frac{(\gamma+1)}{(\gamma-1)}$$

$$\boxed{\frac{p_2}{p_1} \approx \frac{2\gamma M_s^2}{(\gamma+1)}}$$

$$V_2 \approx \frac{2V_s}{\gamma+1}$$

$$\frac{T_2}{T_1} = \frac{2\gamma(\gamma-1)}{(\gamma+1)^2} M_s^2$$

We can have the pressure $p_2$ immediately behind the shock

$$p_2 \approx \frac{2\gamma M_s^2}{(\gamma+1)} p_o = \frac{2\rho_o V_s^2}{(\gamma+1)}$$

Inverting this relation, one can calculate the shock speed as a function of the post-shock pressure and the pre-shock density

$$V_s = (\frac{\gamma+1}{2})^{1/2} (\frac{P_2}{\rho_o})^{1/2}$$

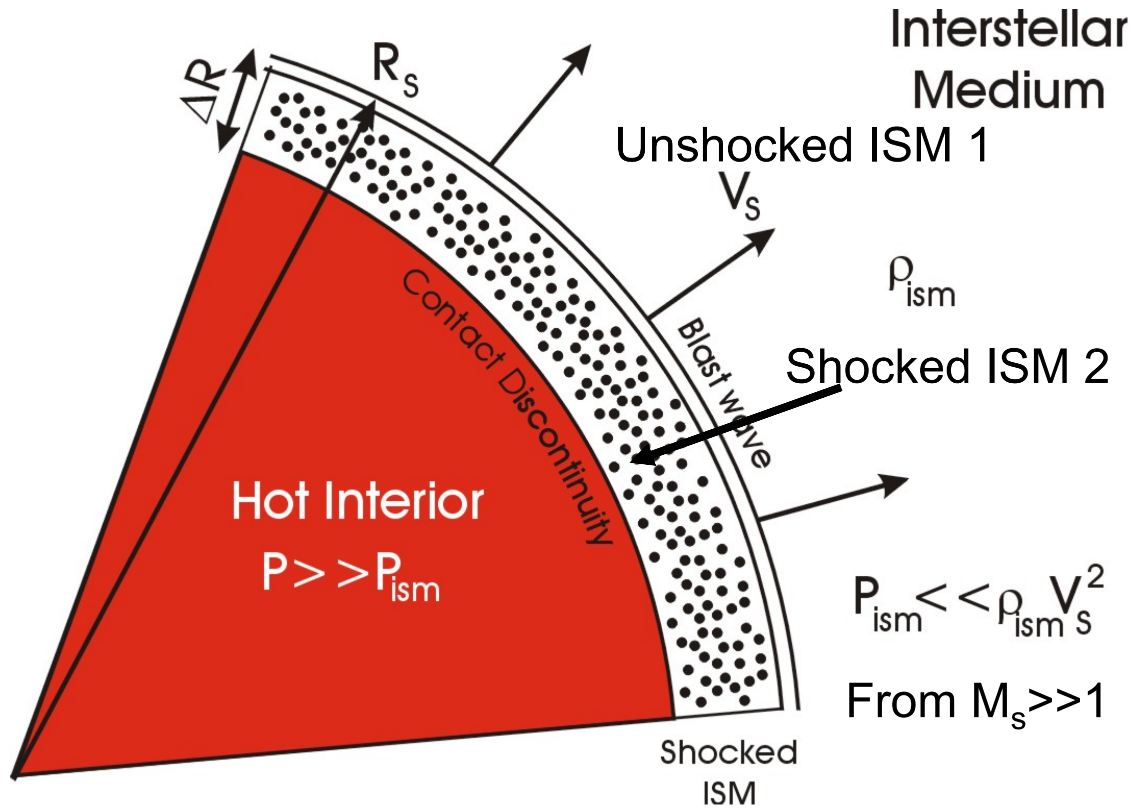
This result can be applied for the formation of high pressure bubbles in a stationary medium

As for istance SuperNova remnants (SNRs) and stellar wind bubbles in the interstellar medium

(...and to nuclear explosions too, unfortunately)

# Blast shock waves

Consider a spherical bubble containing a low density, very hot gas with internal pressure $p_i$ and density $\rho_i$ embedded in a cold, dense stationary medium with low pressure $p_o$ and a high density $\rho_o$ (...we can have low pressure with high density at low T because $p=nkT$)




Because of the high pressure difference, the bubble will start to expand rapidly

# Blast shock waves

If the difference between internal and external pressure is sufficiently large, the expansion speed will be supersonic with respect to sound speed of the surrounding medium

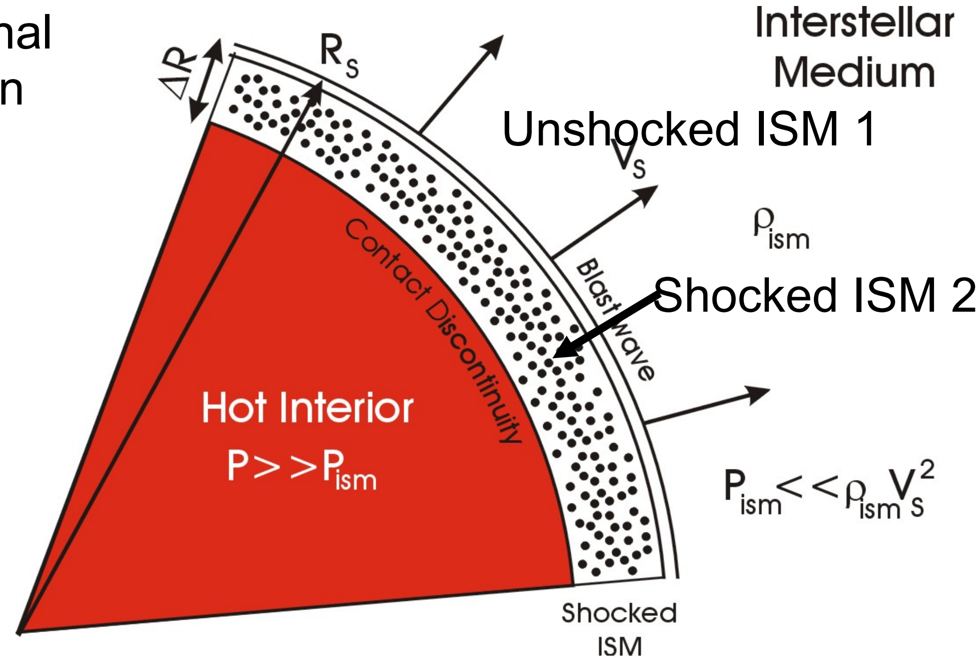
$$M_s^2 = (\frac{V_s}{c_s})^2 = \frac{\rho_o V_s^2}{\gamma p_o} \gg 1$$

$$V_s = (\frac{\gamma + 1}{2})^{1/2} (\frac{P_2}{\rho_o})^{1/2}$$

Substituting Vs in $M^2_s$ we get

$$\mathbf{M^2_s = [(\gamma + 1)/\gamma](p_2/p_0)}$$

Imply that $M_s > 1$ if $p_2/p_o > 2\gamma/(\gamma+1)$




**The key point is that at the interface between hot interior and the shocked material a contact discontinuity forms where there MUST be pressure equilibrium $p_2 = p_i$, because the relative speed between the two media is zero (like in the case of supersonic piston) so that M>>1 if $p_i >> p_0$**

For istance, the typical observed expansion speed of a supernova remnant is ~10000 km/s, while the sound speed in the ISM ranges 10-100 km/s

Because of supersonic speed, a shock will form at the outer edge of the bubble (which acts as a supersonic piston) .

This shock is usually called blast wave

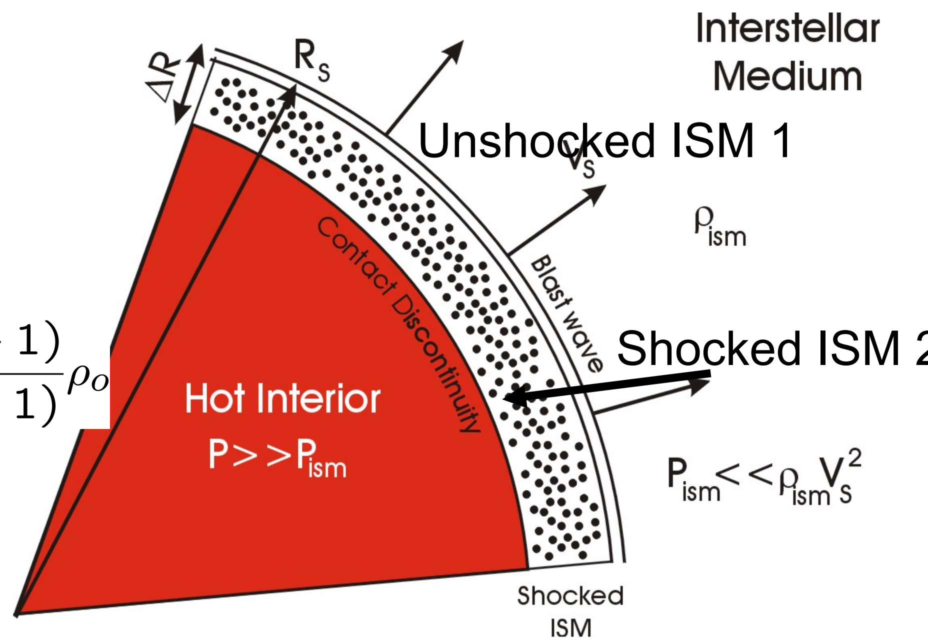
# Blast shock waves

The mass that has been swept up by the expanding bubble will collect in a dense shell at its outer rim

If $M_s$>>1, the typical density of the shocked material in the shell is

$$\rho_{sh} \approx \frac{(\gamma+1)}{(\gamma-1)}\rho_o$$

This allows us to calculate the thickness of the shell




Neglecting, for now, the mass of the bubble (because $\rho_i << \rho_o$), a bubble with radius R has swept up a mass from surrounding ISM

$$M_{sw} = \frac{4}{3}\pi\rho_o R^3$$

This mass is now in a shell with thickness $\Delta R$ with density $\rho_{sh}$ → if $\Delta R << R$

$$M_{sw} \approx 4\pi\rho_{sh} R^2 \Delta R$$

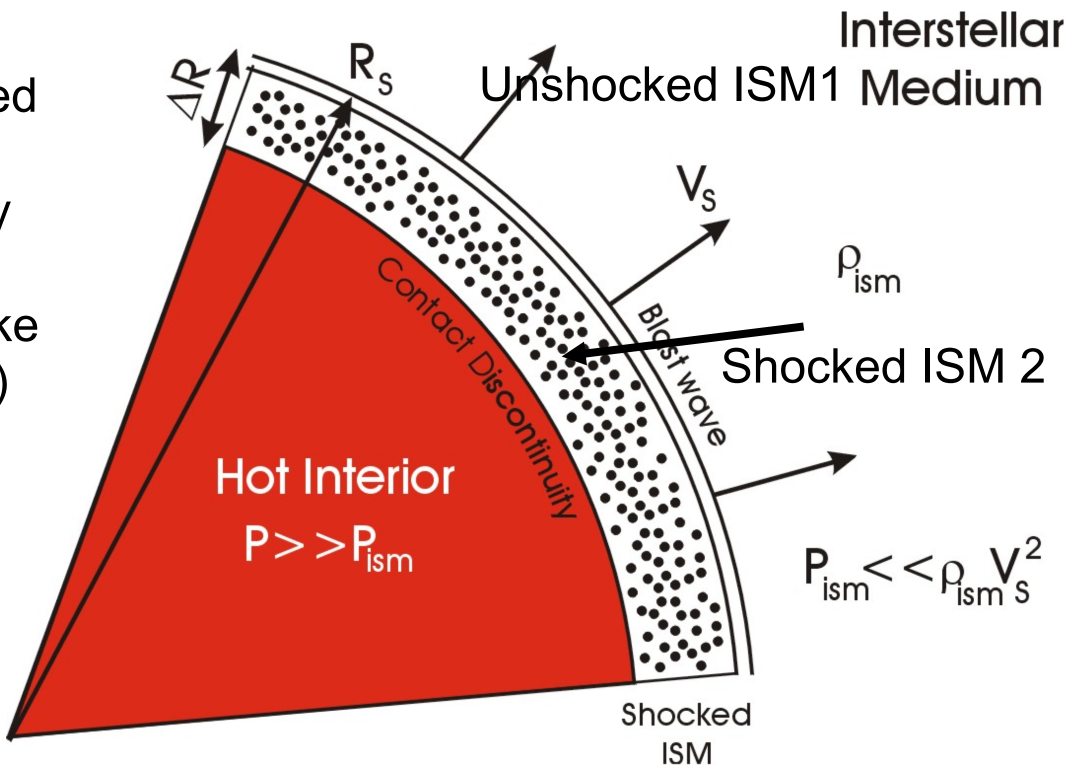
Combining we get $$\Delta R \approx \frac{(\gamma-1)}{3(\gamma+1)} R$$

For $\gamma$ =5/3, $\Delta R$=0.083 R

→ thin shell approximation is fine

# Blast shock waves

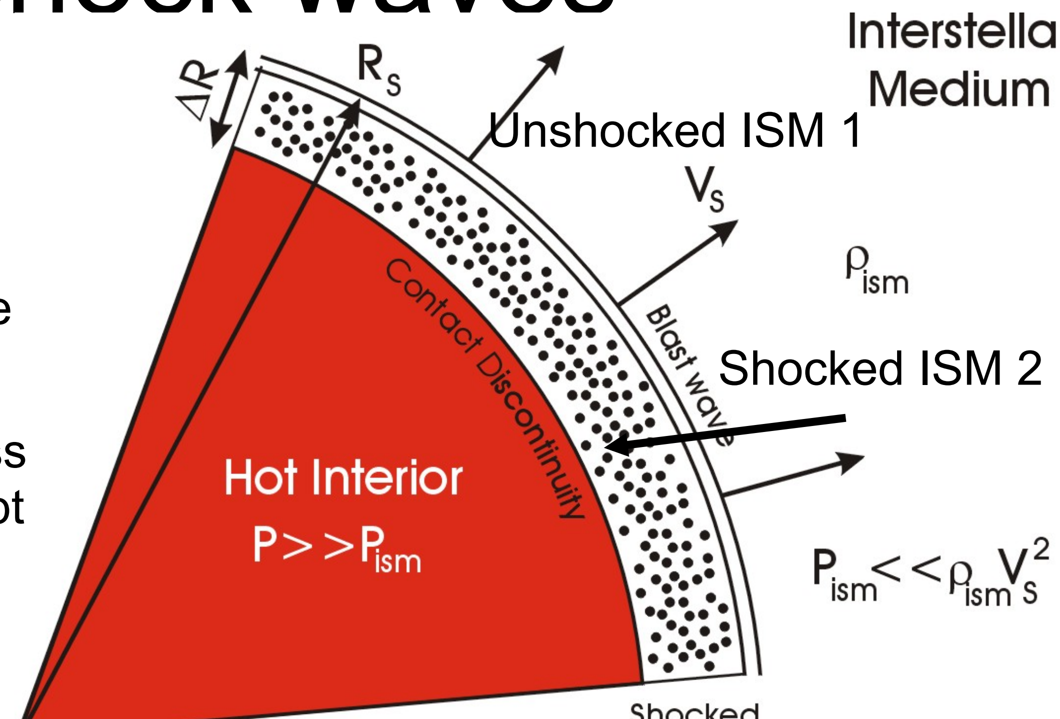
The swept up material is separated from the hot material inside the bubble by a contact discontinuity because the relative speed between the two media is zero (like in the case of supersonic piston)




The expansion speed is

$$\frac{dR}{dt} \approx V_s = (\frac{\gamma+1}{2})^{1/2}(\frac{p_2}{\rho_o})^{1/2}$$

# Blast shock waves




$$\frac{dR}{dt} \approx V_s = \left(\frac{\gamma+1}{2}\right)^{1/2}\left(\frac{P_2}{\rho_o}\right)^{1/2}$$

We can get the time evolution of the shock radius

Let us assume that most of the mass of the expanding bubble is the swept up mass and which resides in the shocked shell of thickness ΔR<<R

→ the mass contained in the bubble is then $M(t) \approx M_{sw} = \frac{4\pi}{3}\rho_o R^3(t)$

The total energy of the bubble consists of the kinetic energy of the expanding massive shell and the internal energy of the hot tenous gas inside the bubble

mistake: here $\rho_o$ is $\rho_i$

$$E(t) = \frac{1}{2}M(t)\left(\frac{dR(t)}{dt}\right)^2 + \left(\frac{4\pi}{3}\rho_o R^3\right)\left(\frac{p_i(t)}{\rho_o(\gamma-1)}\right)$$

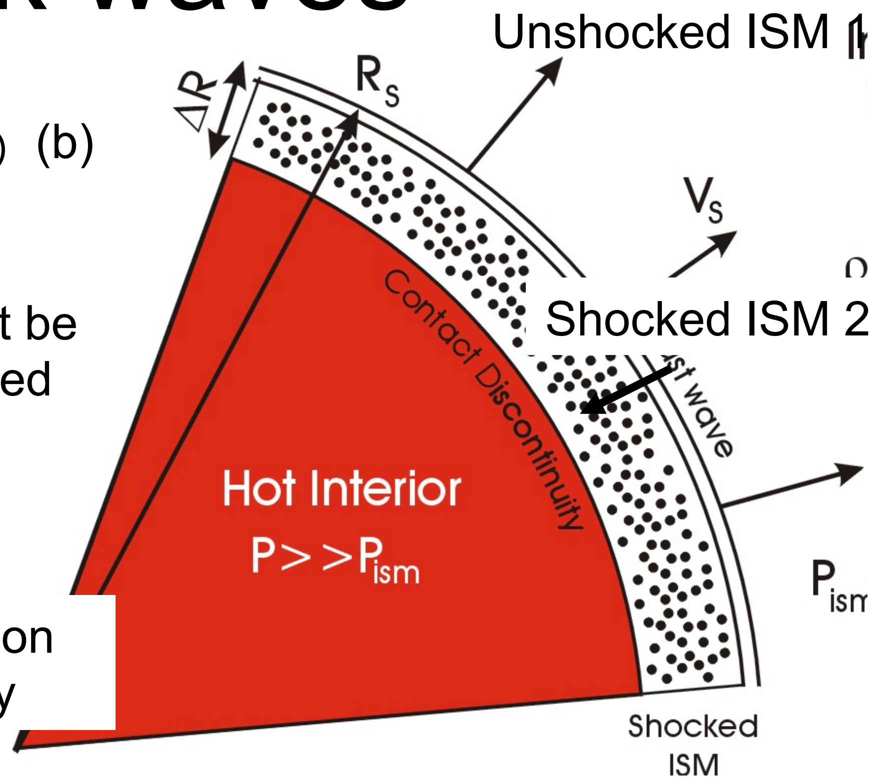
Assuming uniform pressure $p_i$ inside the bubble

Internal energy U = uV = (ρe)V, e = specific energy

# Blast shock waves

mistake: here $\rho_o$ is $\rho_i$

$$E(t) = \frac{1}{2}M(t)(\frac{dR(t)}{dt})^2 + (\frac{4\pi}{3}\rho_o R^3)(\frac{p_i(t)}{\rho_o(\gamma-1)}) \quad \text{(b)}$$




The interior pressure in the hot bubble must be roughly equal to the pressure in the shocked material in the shell

$$p_i \approx p_2 \approx \frac{2}{\gamma+1}\rho_o(\frac{dR}{dt})^2 \quad \text{(a)}$$

Which is simply the pressure-balance condition which must hold at the contact discontinuity

Inserting (a) into (b) we get $$E(t) = \frac{2\pi}{3}\rho_o R^3(\frac{dR}{dt})^2[1+\frac{4}{\gamma^2-1}]$$

→ the ratio of thermal energy and kinetic energy is a constant $E_{th}/E_{kin} = \frac{4}{\gamma^2-1}$

For $\gamma$=5/3, $E_{th}/E_{kin}$ =9/4 → $E_T$=$E_{th}$+$E_{kin}$= (1+9/4)$E_{kin}$=(13/4)$E_{kin}$

➔ $E_{kin}$ = (4/13)$E_T$ ~ 31% $E_T$...most part of the initial energy is thermal energy

# Blast shock waves

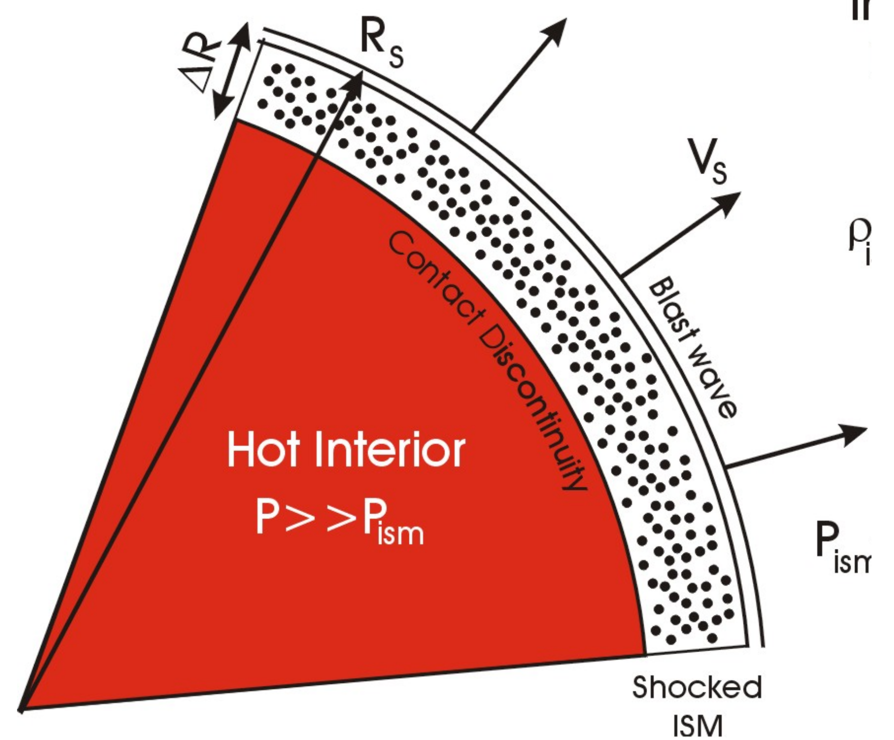
$$E(t) = \frac{2\pi}{3}\rho_o R^3 (\frac{dR}{dt})^2 [1 + \frac{4}{\gamma^2 - 1}]$$

→ the total energy can be then written as

$$E(t) = C_\gamma M(t) (\frac{dR}{dt})^2$$

With $$C_\gamma = \frac{\gamma^2 + 3}{2(\gamma^2 - 1)}$$

For an ideal monoatomic gas $C_\gamma$=1.625

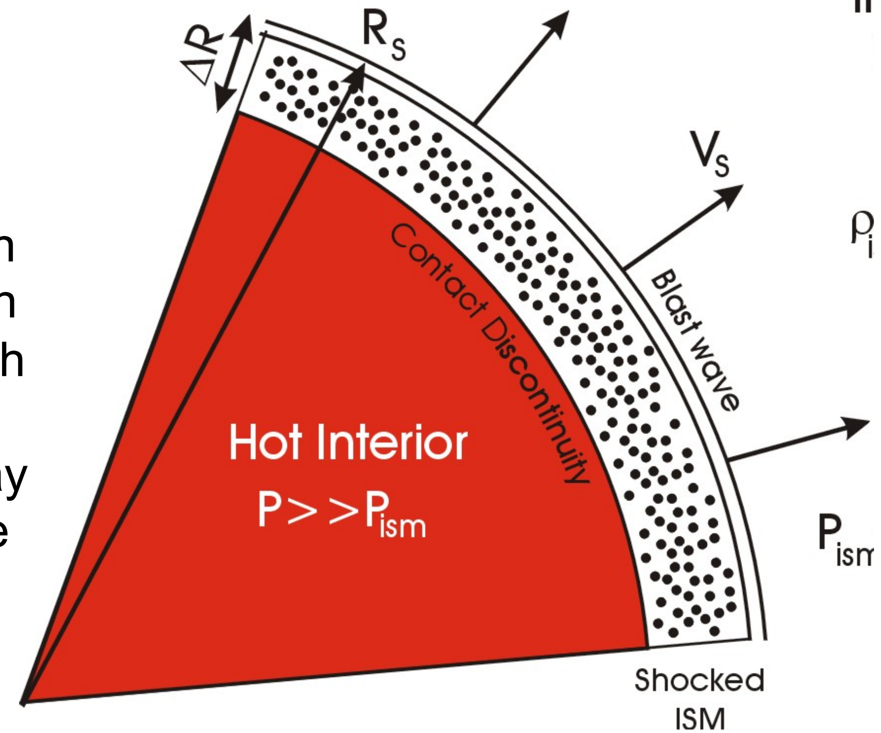



$C_\gamma$ is approximate because of the various approximations made in the derivation. However, more exact treatments come at the same result with a somewhat smaller value

# Blast shock waves

$$E(t) = C_\gamma M(t)(\frac{dR}{dt})^2$$

The energy may depend on time, because in general there may be some energy source in the bubble which fuels the hot bubble through emission of power L, as for istance a star driving a strong stellar wind, or the SNRs may radiate away part or all their energy at some point of its evolution




Two cases:
i) a point explosion where a fixed amount of energy is supplied impulsively at t=0 and where no energy losses occur afterwards (blast shock wave, SN explosions)
Ii) a constant energy supply at some luminosity L=dE/dt so that E(t)=L x t, which can serve as a crude model of the energy of a bubble blown into the ISM by a strong stellar wind

# Blast shock waves

Blast wave generated by a 4.8 kiloton explosion

Explosions (eg nuclear detonations) generate blast shock wave, very much like a supernova: the piston is the exploding materal

# Blast shock waves

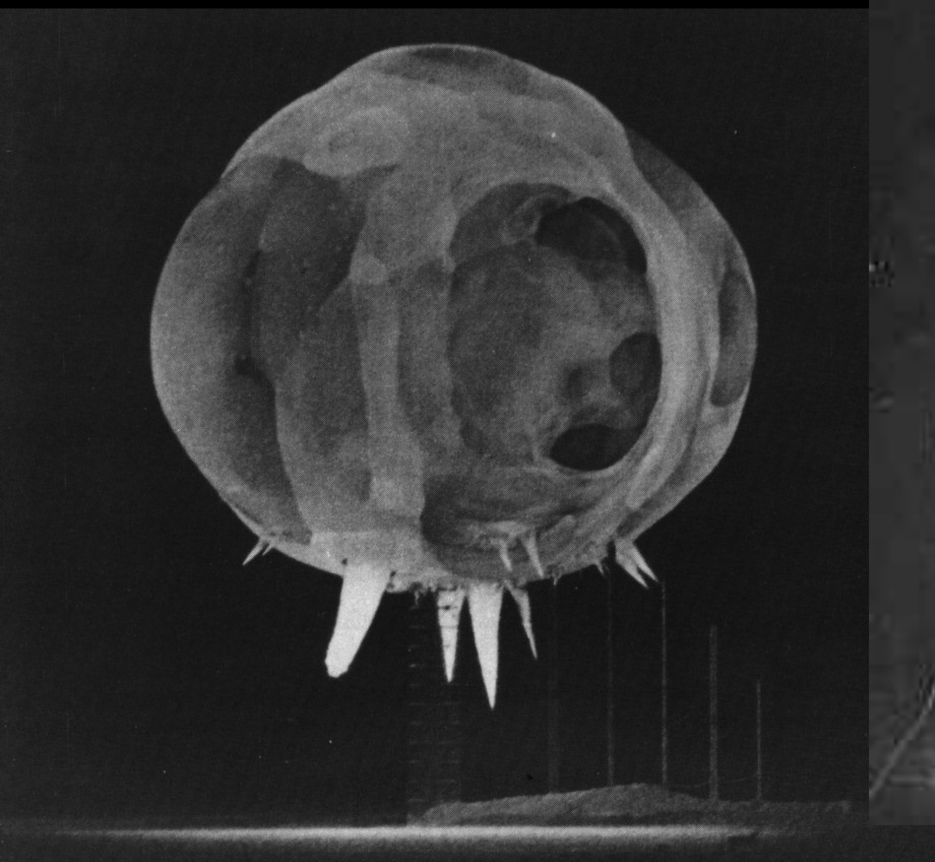
Most of the "experimental" knowledge comes from military nuclear tests during the cold war

For long time, and still now, most of the data on nuclear tests are classified top secret

Whatever the mechanism by which initial energy is liberated, the subsequent blast wave evolution does not depend on it!

# Blast shock waves

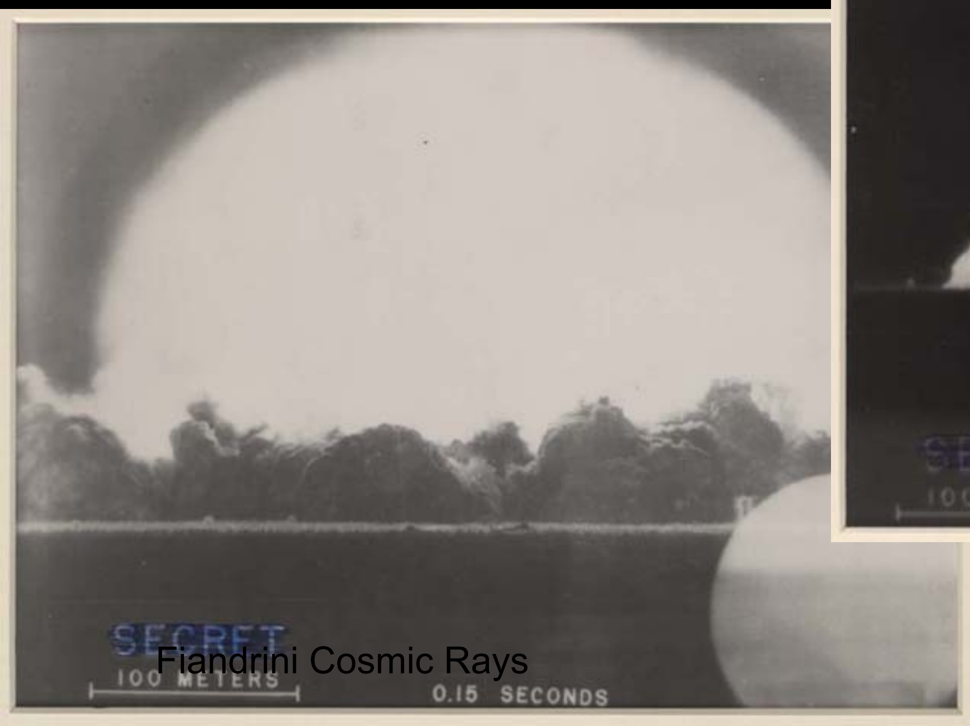
10 ms from detonation



1 ms from detonation

# Trinity test





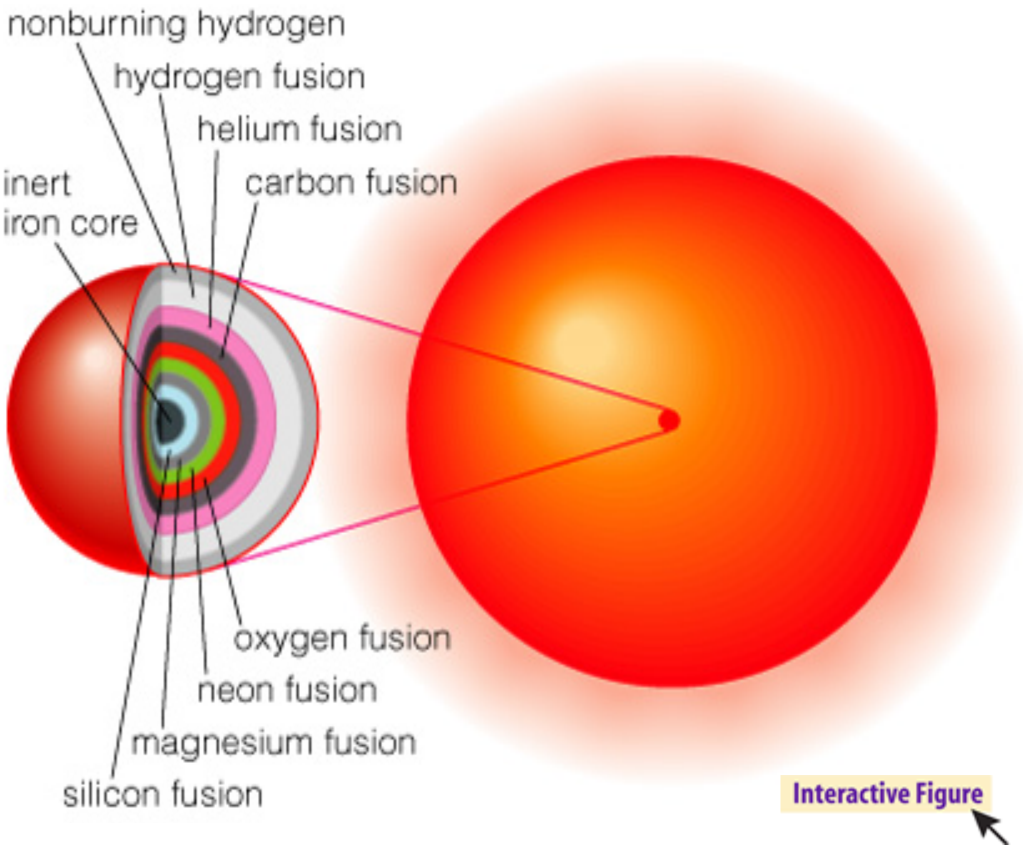
# Supernova Blast Waves

| | |
|---|---|
| Supernovae: | |
| Type Ia | Mechanism: explosive carbon burning in a mass-accreting white dwarf |
| Type Ib-Ic & Type II | Core collapse of massive star |

We may neglect the details of the explosive event: for us is a sudden, point-like release (a $\delta(x,t)$) of energy

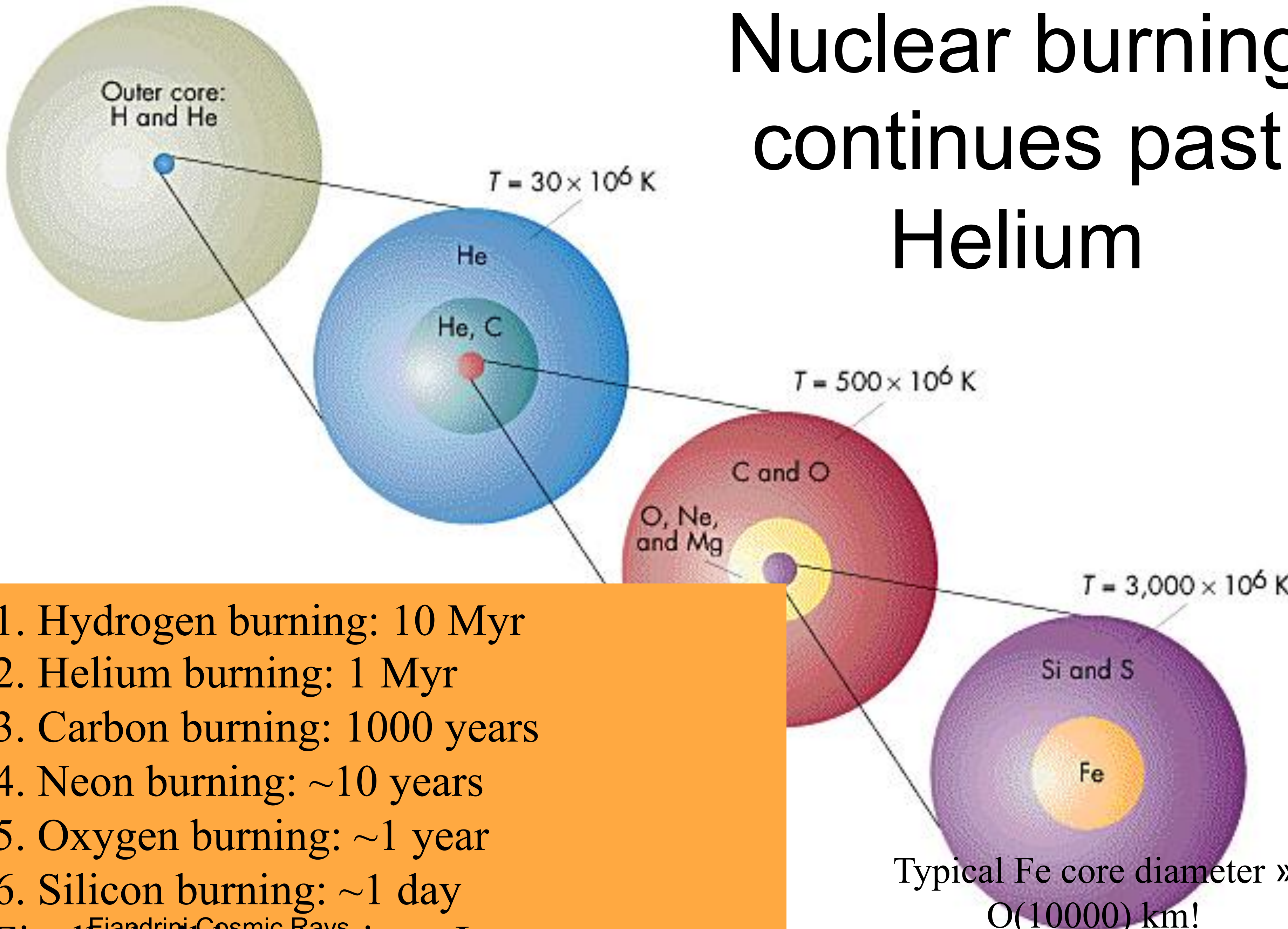
While the composition of the material of the SNR does depend on the mechanism (ie star composition) and on the surrounding ISM composition

# Multiple Shell Burning




- Advanced nuclear burning proceeds in a series of nested shells, like an onion

# Nuclear burning continues past Helium



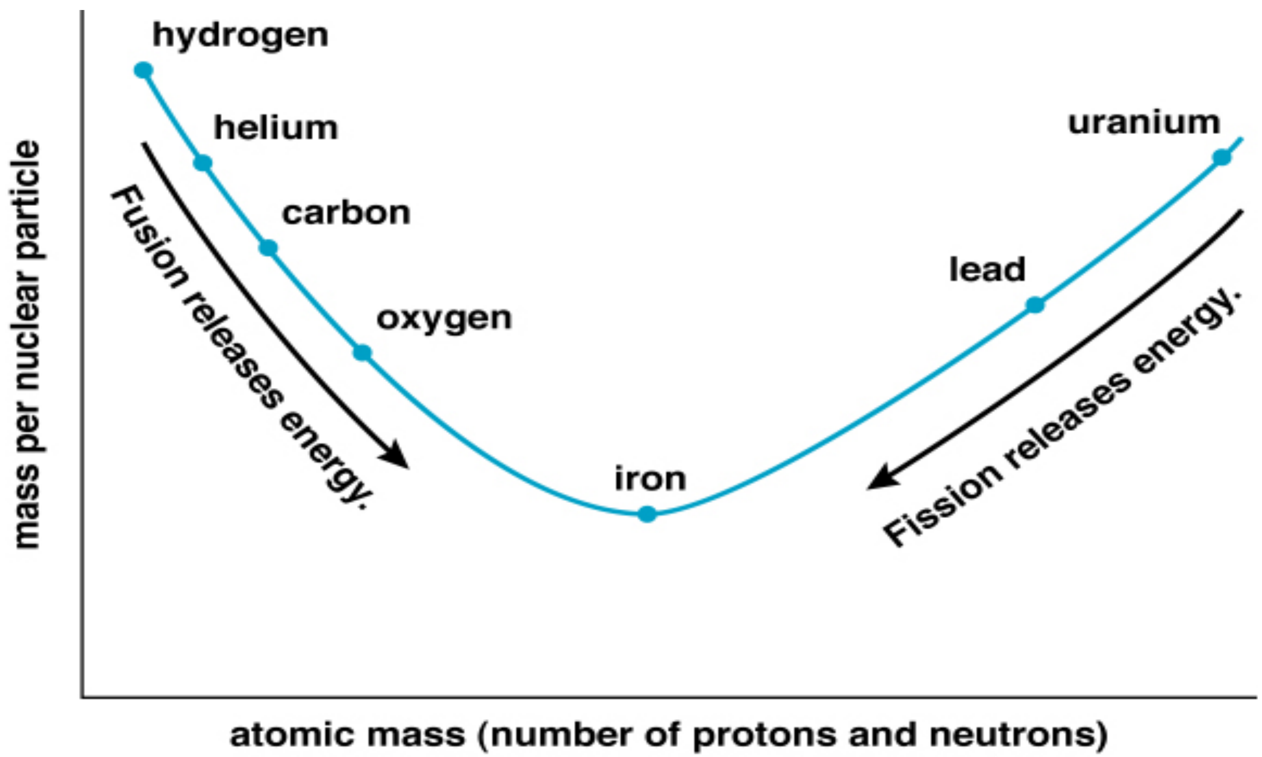

1. Hydrogen burning: 10 Myr
2. Helium burning: 1 Myr
3. Carbon burning: 1000 years
4. Neon burning: ~10 years
5. Oxygen burning: ~1 year
6. Silicon burning: ~1 day

Finally builds up an inert Iron core

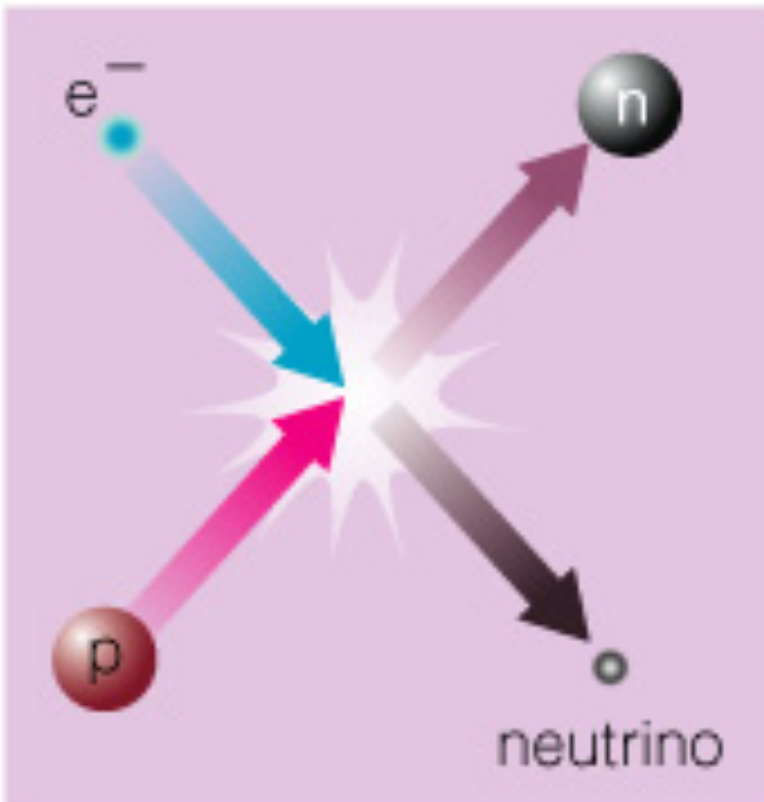
Typical Fe core diameter » O(10000) km!

- The supergiant has an inert Fe core which collapses & heats
  - Fe can not fuse
  - It has the lowest mass per nuclear particle of any element
  - It can not fuse into another element without *creating* mass

So the Fe core continues to collapse until it is stopped by electron degeneracy.
(like a White Dwarf)



# Supernova Explosion



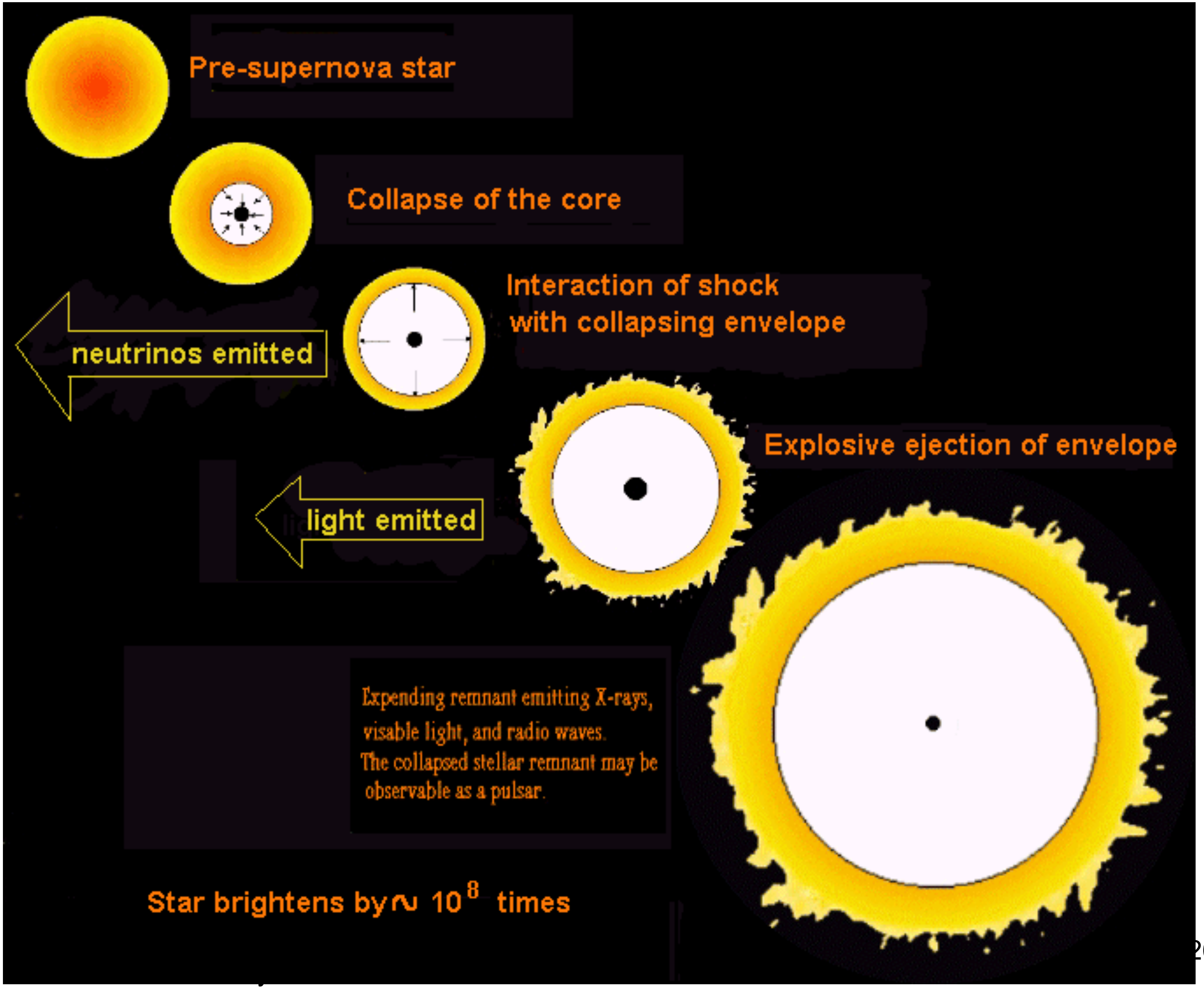

- Core degeneracy pressure goes away because electrons combine with protons, making neutrons and neutrinos

$P + e^- \rightarrow n + \nu_e$

- Neutrons collapse to the center, forming a **neutron star**

# Core collapse

- Iron core is degenerate
- Core grows until it is too heavy to support itself ($M_{nuc} > M_{chandrasekhar}$)
- Core collapses, density increases, normal iron nuclei are converted into neutrons with the emission of neutrinos
- Core collapse stops, neutron star is formed
- Rest of the star collapses on the core, but bounces off the new neutron star (also pushed outwards by the neutrinos)


Pre-supernova star
Collapse of the core
Interaction of shock
with collapsing envelope
neutrinos emitted
Explosive ejection of envelope
light emitted
Expending remnant emitting X-rays,
visable light, and radio waves.
The collapsed stellar remnant may be
observable as a pulsar.
Star brightens by ~ 10^8 times

# Shock Waves in Supernovae

- Discontinuity in velocity, density and pressure in a flow of matter.
- Unlike a sound wave, it causes  a permanent change in the medium
- Shock speed >> sound speed -  between 30,000 and 50,000 km/s.
- Shock wave may be *stalled* if energy goes into breaking-up nuclei into nucleons.
- This consumes a lot of energy, even though the pressure (nkT) increases because n is larger.

# Supernova Energetics

Same source for supernovae (Ib/Ic and II),

- Explosion powered by the collapse (death) of a massive core
- Energy source: Potential Energy from the collapse of the iron core down to a neutron star or black hole:

$$\text{Energy} = G\, M^2_{core}/r_{NS,BH} - G\, M^2_{core}/r_{\text{before collapse}}$$
$$> 10^{53} \text{ ergs}$$

$M_{core} \sim$ 1.4-3 solar masses

$R_{NS,BH} \sim 10$ km

$R_{core} \sim 10{,}000$ km

# Supernova energetics

## Core Collapse Supernova Energetics

Liberated gravitational binding energy of neutron star:

$$E_b \approx 3 \times 10^{53}\ \text{erg} \approx 17\%\ M_{SUN}c^2$$

This shows up as

| | |
|---|---|
| 99% | Neutrinos |
| 1% | Kinetic energy of explosion (1% of this into cosmic rays) |
| 0.01% | Photons (outshine host galaxy) |

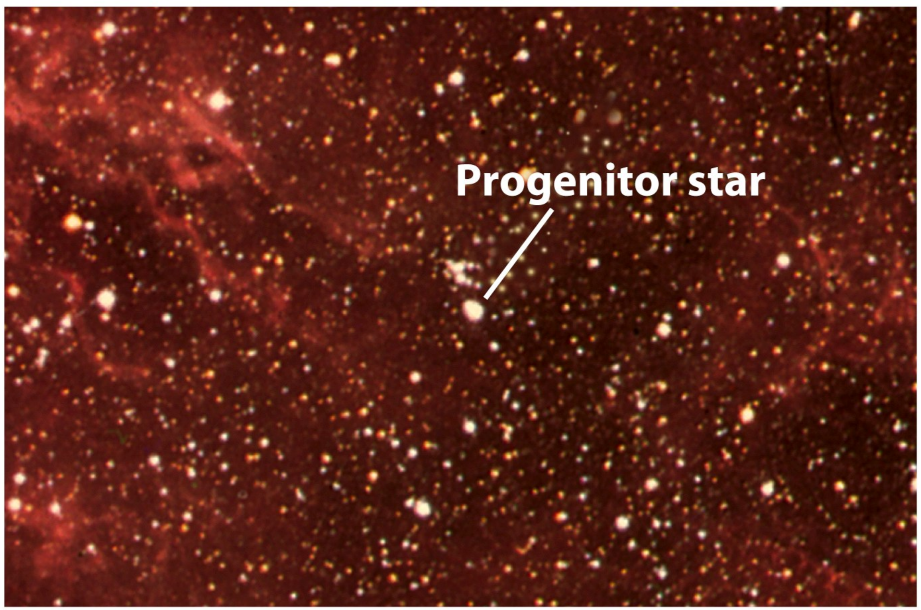
Neutrino luminosity

$$L_\nu \approx 3\times10^{53}\ \text{erg} / 3\ \text{sec} \approx 3\times10^{19}\ L_{SUN}$$

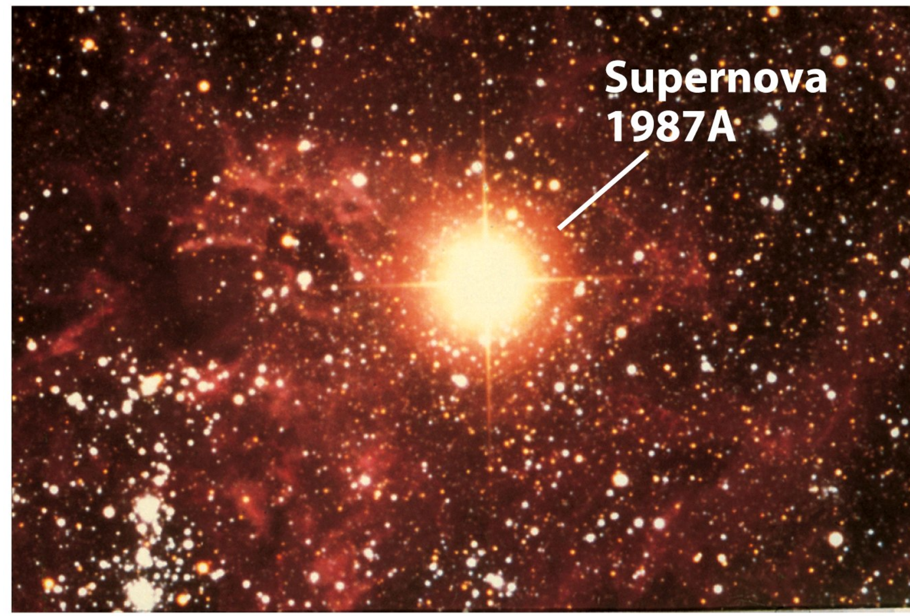
While it lasts, outshines the photon luminosity of the entire visible universe!

Georg Raffelt, Max-Planck-Institut für Physik (München)

# In 1987 a nearby supernova gave us a close-up look at the death of a massive star



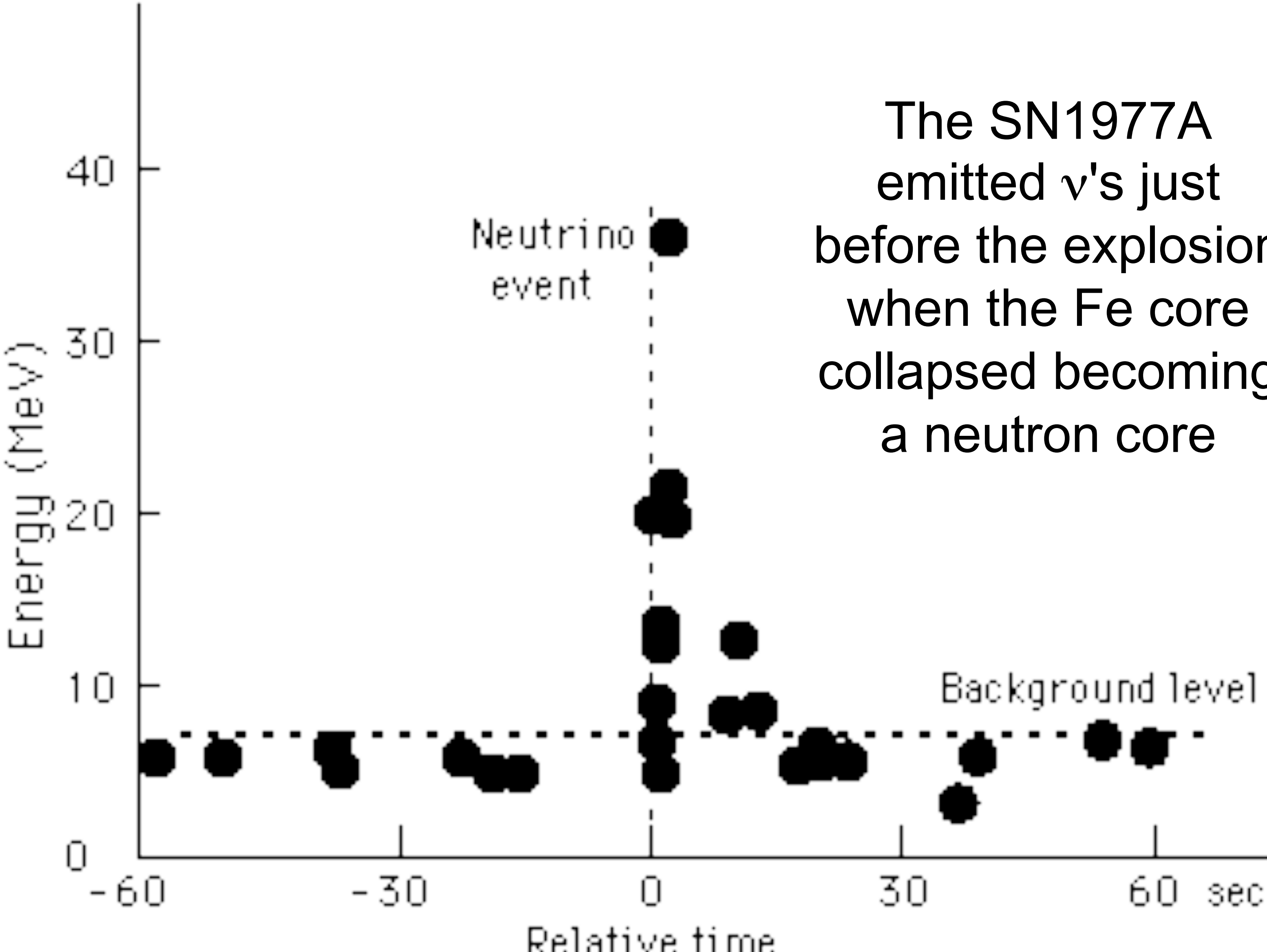

Before the star exploded




After the star exploded

At peak its luminosity was greater than host galaxy luminosity (Large magellanic Cloud)

# Neutrinos from SN1987A

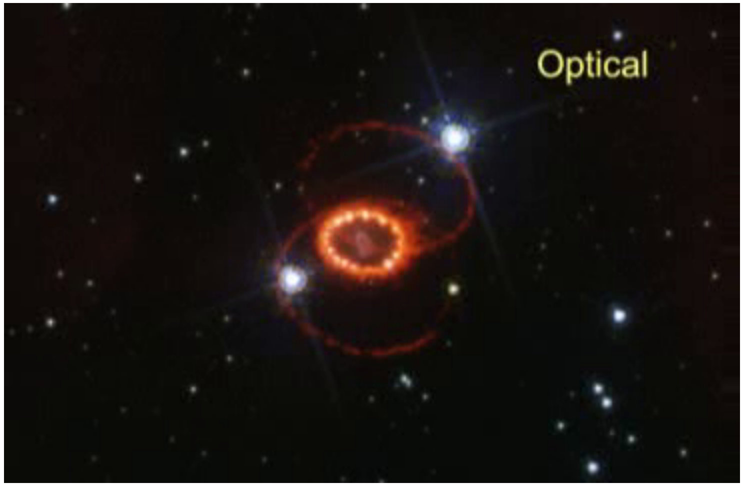



The SN1977A emitted ν's just before the explosion when the Fe core collapsed becoming a neutron core

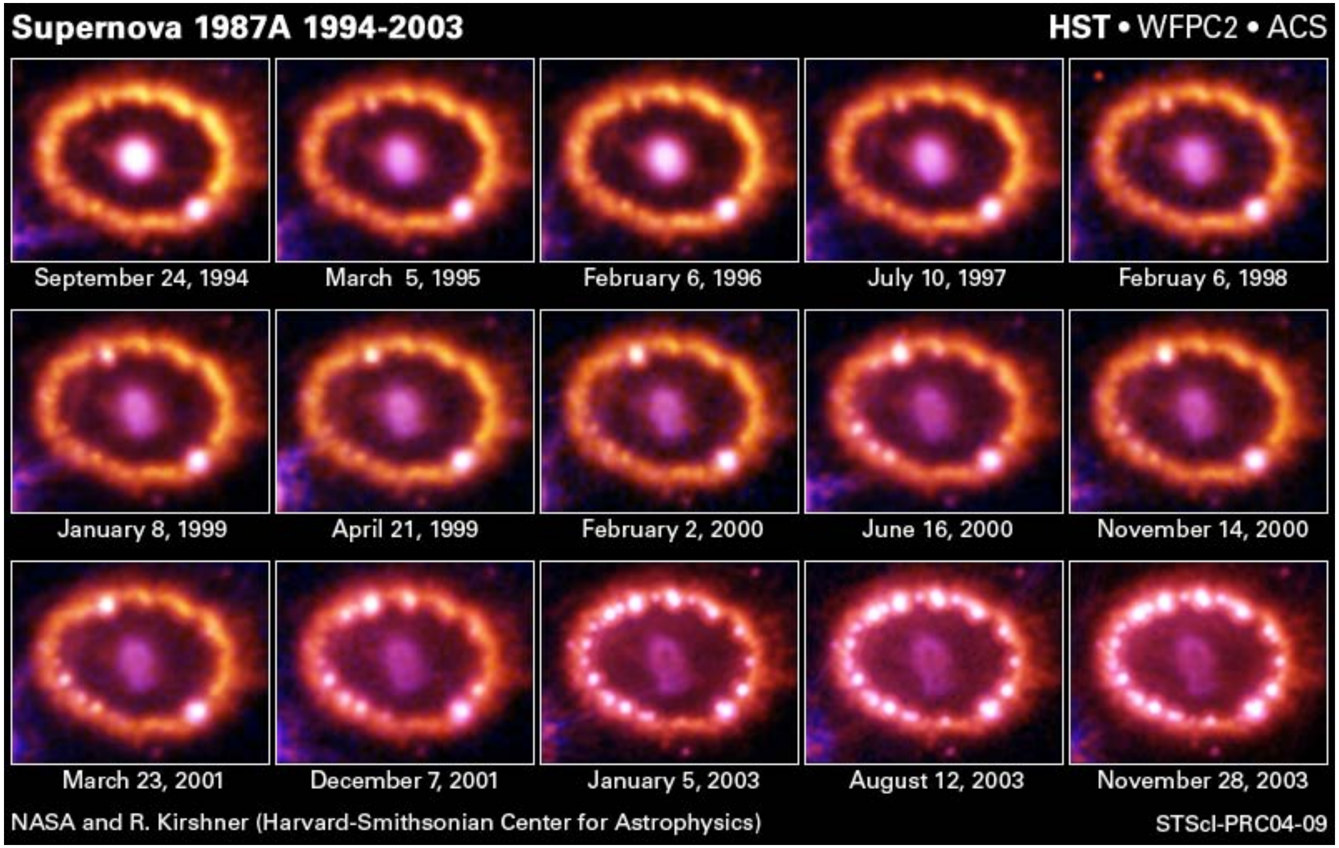
# SN1987A – Blue Supergiant Supernova?

The progenitor of SN1987A was a blue giant with a mass of about 18 $M_{sun}$.

- Probably, the high-mass progenitor of SN1987A lost most of its outer layer by a slow stellar wind long before the supernova explosion.
- Right before the supernova explosion, a fast wind pushes the envelop to make a cavity around the star. Making the outer layer of the star unusually thin and warm
- The outer gas cloud forms a ring.
- The shockwave from the supernova explosion was expected to hit the edge of the ring around 1999.
- Chandra X-ray images from 1999 to 2005 shows brightening of the ring.



# Shock hits inner ring

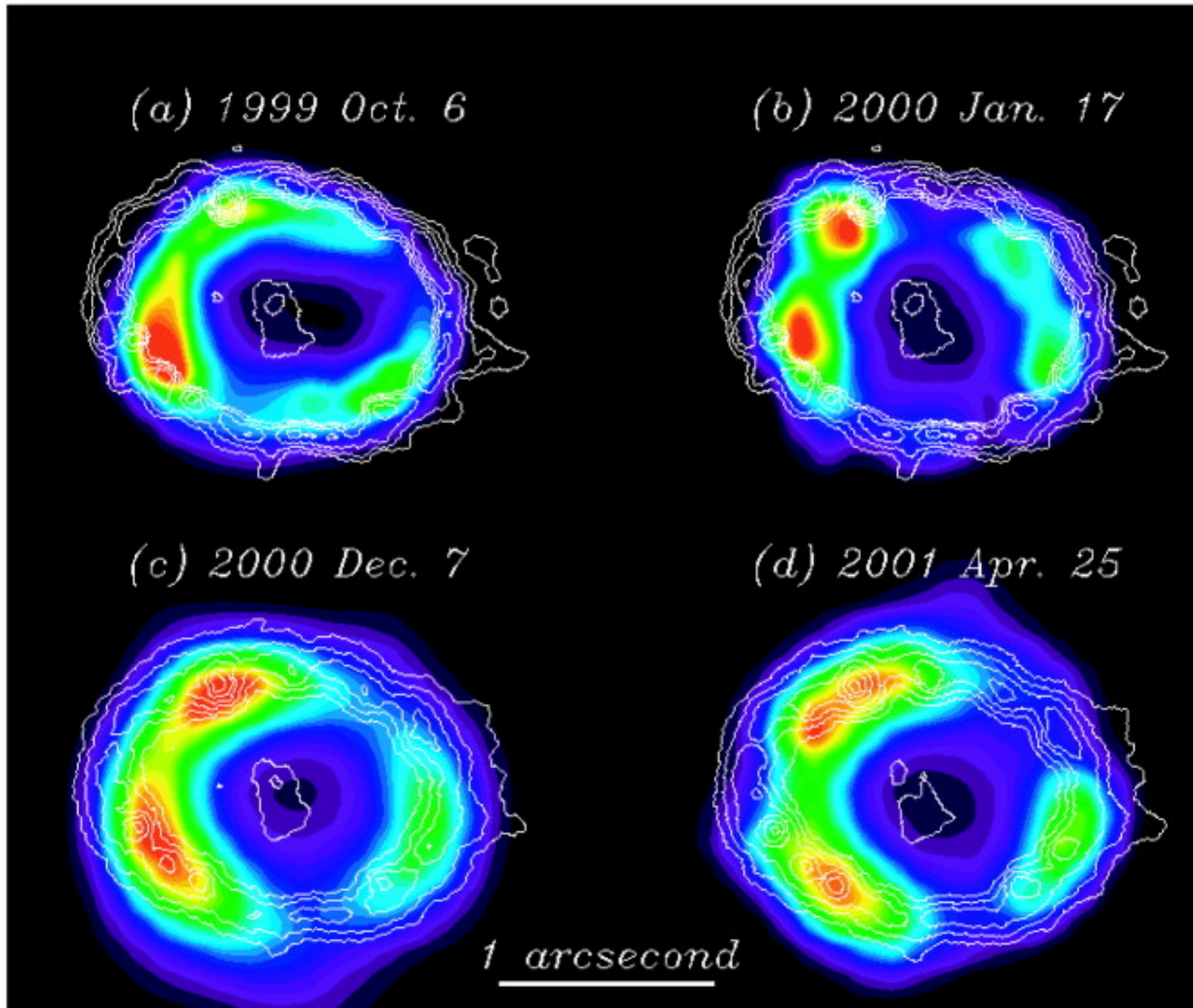



The shock has hit the inner ring at 20,000 km/s, lighting up knots of shocked, compressed and heated material in the ring which is 160 billion km wide.

# Chandra X-ray Images of SN 1987A

- X-ray intensities (0.5 – 8.0 keV) in colour with HST Hα images as contours

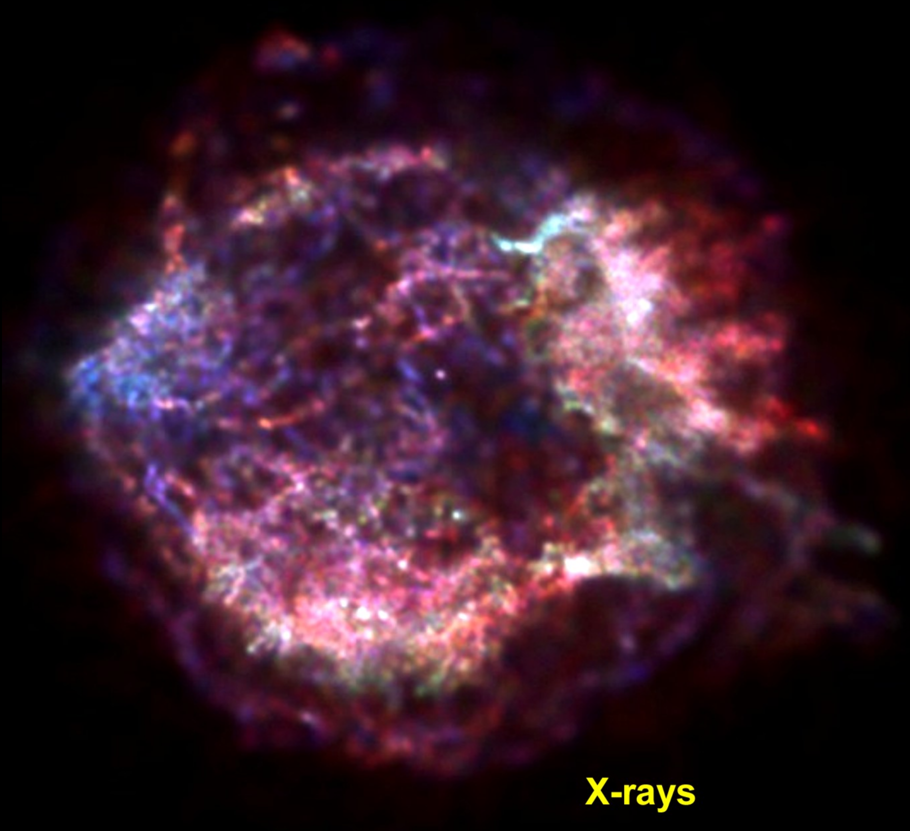
- Low energy X-rays are well correlated with optical knots in ring – dense gas ejected by progenitor?
- Higher energy X-rays well correlated with radio emission – fast shock hitting circumstellar H II region?
- No evidence yet for emission from central pulsar

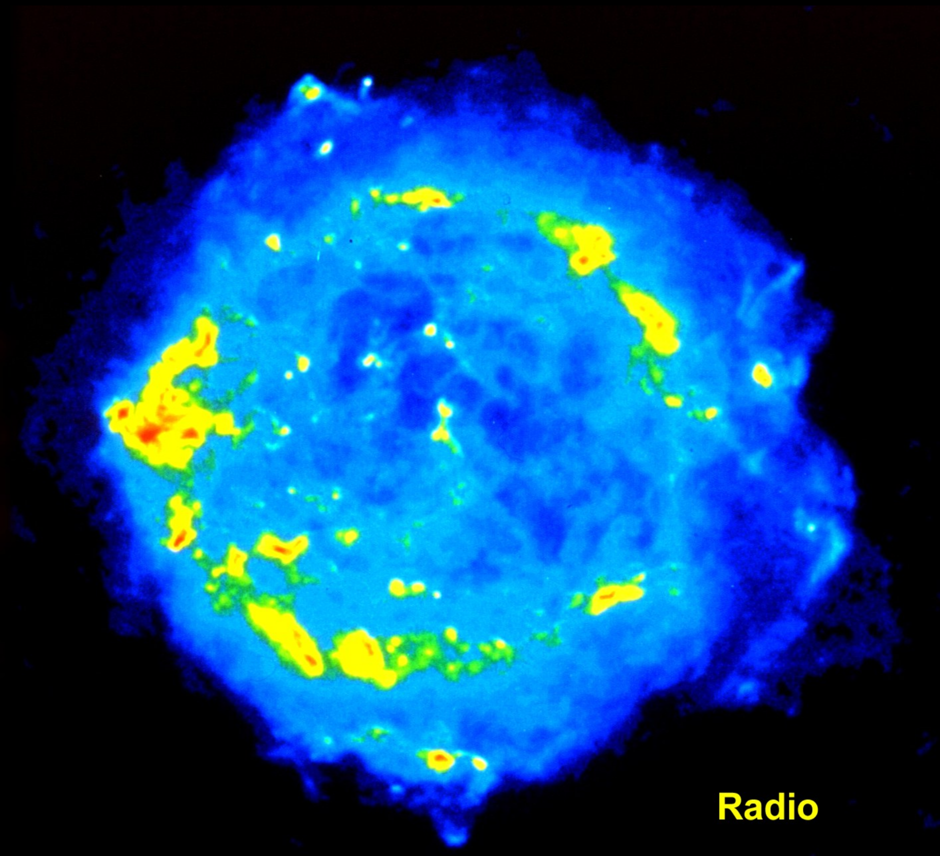


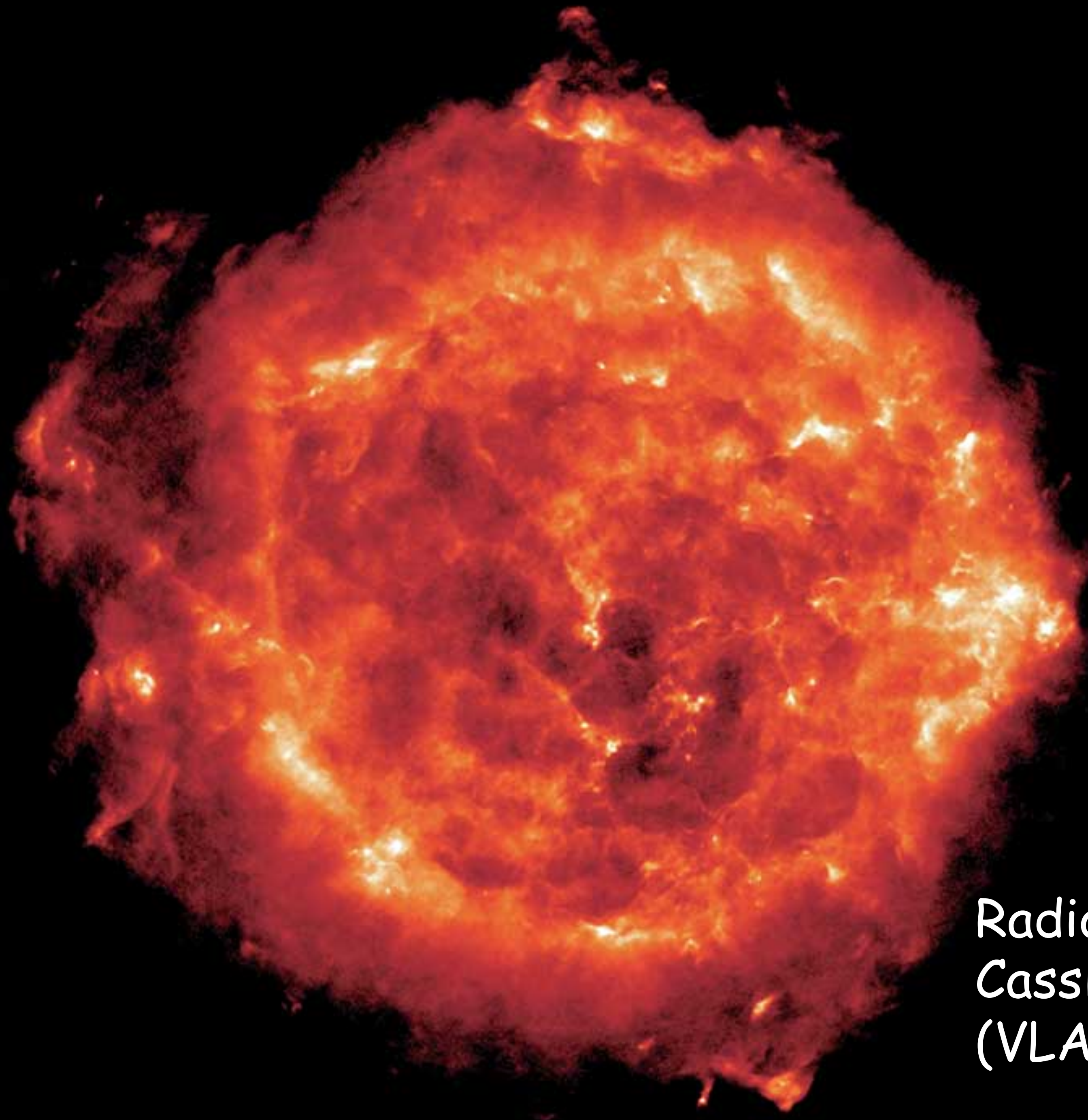
# SuperNova Remnants (SNRs)

A supernova remnant (SNR) consists essentially of the stellar ejecta of the SN explosion embedded in a hot expanding bubble, preceded by swept-up interstellar material and an outer blast wave (strong shock) propagating into the interstellar medium

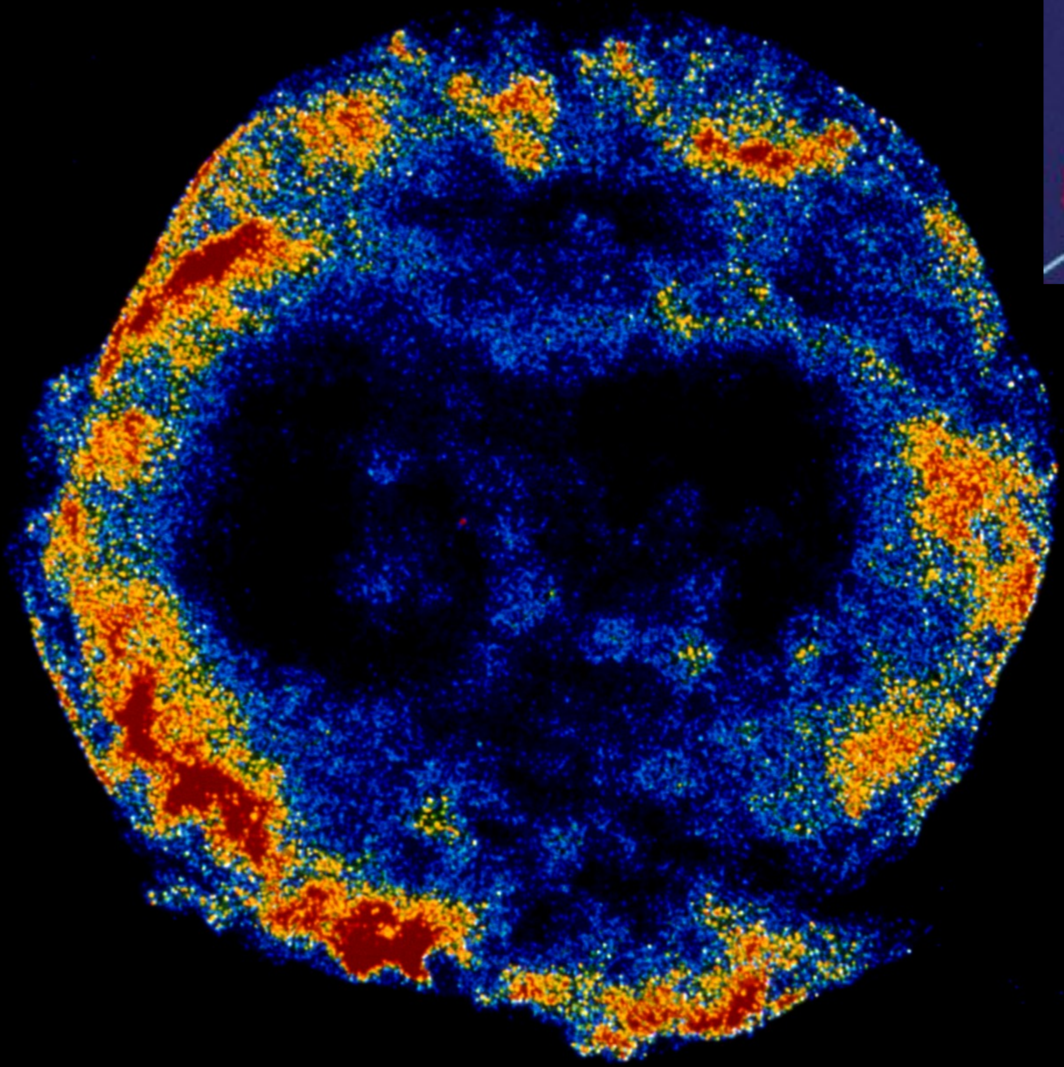
# Supernova Remnant Cassiopeia A

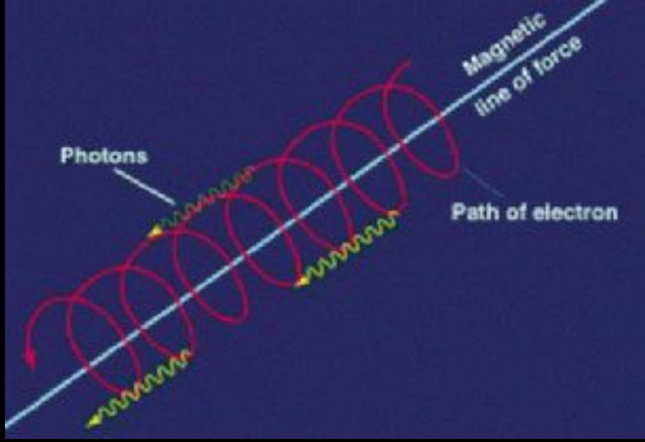




Radio map
Cassiopeia A
(VLA)






Remnant of Tycho's supernova of 1572 AD

# SNR evolution: energy budget

The mechanical energy of the ejecta is of the order $E_{snr}$~ 0.01x $E_{sn}$~ $10^{51}$ erg

This is the energy that fuels the explosive event and ultimately creates a SNR

The typical speed $V_s$ of this material can be estimated by energy conservation

Let us assume that all the mechanical energy is converted into kinetic energy of the remnant and that energy loss is negligible

$$E_{snr} = \frac{1}{2} M_{snr} V_s^2 \quad \Rightarrow \quad V_s = (\frac{2E_{snr}}{M_{snr}})^{1/2}$$

The mass is $\mathbf{M_{snr} = M_{ej} + M_{sw}}$, the sum of explosively ejected mass from the star at the time of explosion and the mass swept up added later as the remnant sweeps up more and more ISM material

# SNR evolution: free expansion phase

If the density of the ISM is constant, then $M_{snr} = M_{ej} + \frac{4\pi}{3}\rho_o R_s^3$

Since we know the typical energy involved and that the mass of the remnant must be several solar masses, we can estimate the typical expansion speed

$$V_s \simeq 3000\left(\frac{E_{snr}}{10^{51}\ erg}\right)^{1/2}\left(\frac{M_{snr}}{10 M_{sun}}\right)^{-1/2}\ km/s$$

Initially, the mass consists almost entirely of the ejecta mass $M_{snr} \sim M_{ej} \sim 2 - 10\ M_{sun}$

➔ the expansion speed is almost constant $V_s \approx \left(\frac{2E_{snr}}{M_{ej}}\right)^{1/2} = 10000\ km/s$

This is >> of sound speed (~10-100 km/s)→ shock must form

The bubble expands as if there is not surrounding medium, therefore this phase is called free expansion phase and lasts few hundreds of years until the swept up mass increases enough to start the deceleration

# Sedov-Taylor phase

As more and more ISM gas is swept up, the mass of remnant increases

After few hundreds of years, the mass is dominated by the swept up material so that

$$M_{snr} = M_{ej} + \frac{4\pi}{3}\rho_o R_s^3 \approx \frac{4\pi}{3}\rho_o R_s^3 \qquad V_s \approx (\frac{2E_{snr}}{M_{ej}})^{1/2} \approx (\frac{6E_{snr}}{4\pi\rho_o})^{1/2} R_s^{-3/2}$$

Speed decreases as $R^{-3/2}$, result of the increasing remnant mass ~ $R^3$

This is the so called Sedov-taylor phase or energy-conserving phase

The typical expansion speed remains supersonic for a considerable time, typically >100000 years, so that the shock at the outer boundary persists in this evolutionary phase

# Sedov-Taylor phase

The transition between the free expansion and the Sedov-taylor phases occurs gradually when the radius of the remnant reaches the deceleration radius, defined as the radius at which $M_{ej} = M_{sw}$

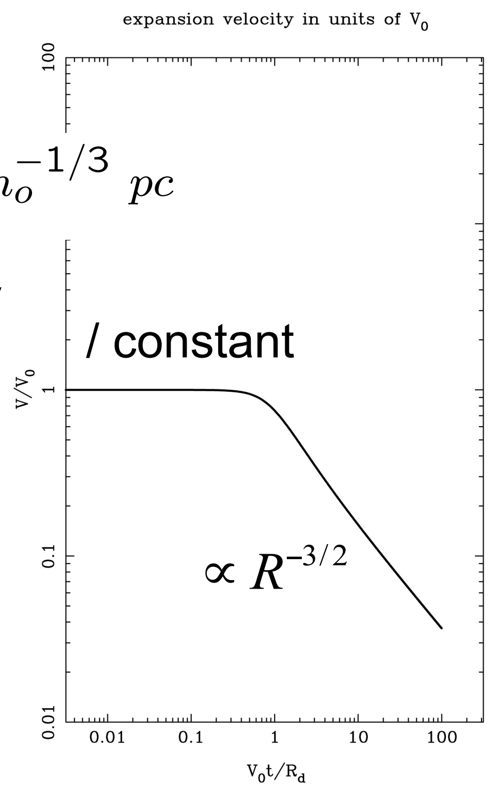
$$\frac{4\pi}{3}\rho_o R_d^3 = M_{ej}$$

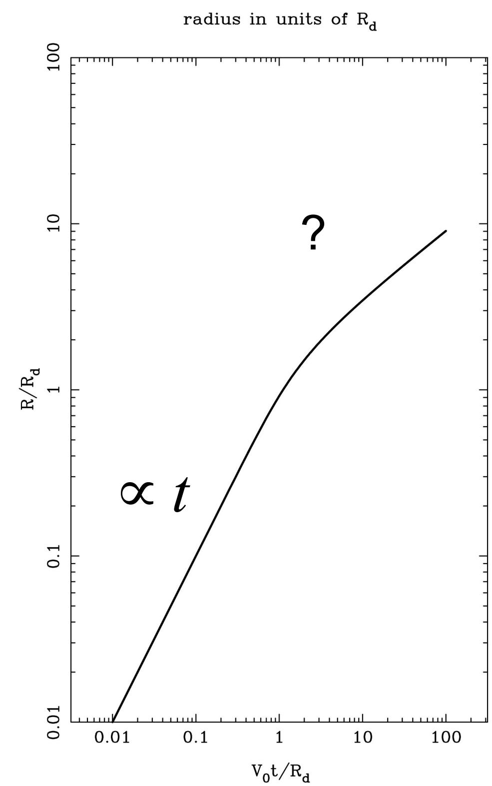
$$R_d = \left(\frac{3M_{ej}}{4\pi\rho_o}\right)^{1/3} \approx 2.2\left(\frac{M_{ej}}{M_{sun}}\right)^{1/3} \times n_o^{-1/3} \; pc$$

Here $n_o = \rho_o/m_p$ is the number density of ISM, which is typically ~1 $cm^{-3}$

Assuming $M_{ej} = 5M_{sun}$ we have $R_d \sim 4$ pc

A crude extimation of time spent during the free expansion phase is ~ $R_d/V_{free}$ ~ 400 years



# Sedov-Taylor phase

In a SN explosion, the mechanical energy $E_0$~ $E_{snr}$ the drives the expansion is supplied impulsively in a point explosion at t=0

<u>If no energy is lost, for istance through radiation losses, E remains constant for t>0</u>

$$E(t) = C_\gamma M(t)(\frac{dR}{dt})^2 = constant \qquad C_\gamma = \frac{\gamma^2 + 3}{2(\gamma^2 - 1)}$$

Once the remnant has expanded to a radius larger than $R_d$, the mass is

$$M(t) \approx M_{sw} = \frac{4\pi}{3}\rho_o R^3(t)$$

→ the energy equation can written as $R^{3/2}\frac{dR}{dt} = (\frac{3E_{snr}}{4\pi C_\gamma \rho_o})^{1/2} = constant$

# Sedov-Taylor phase

$$R^{3/2}\frac{dR}{dt} = (\frac{3E_{snr}}{4\pi C_\gamma \rho_o})^{1/2} = constant$$

This relationship between the speed and the radius of the bubble is the same one as derived above using a simple conservation law for the kinetic energy but in this derivation we also take into account of the thermal energy of the hot bubble material

The integration is straightforward

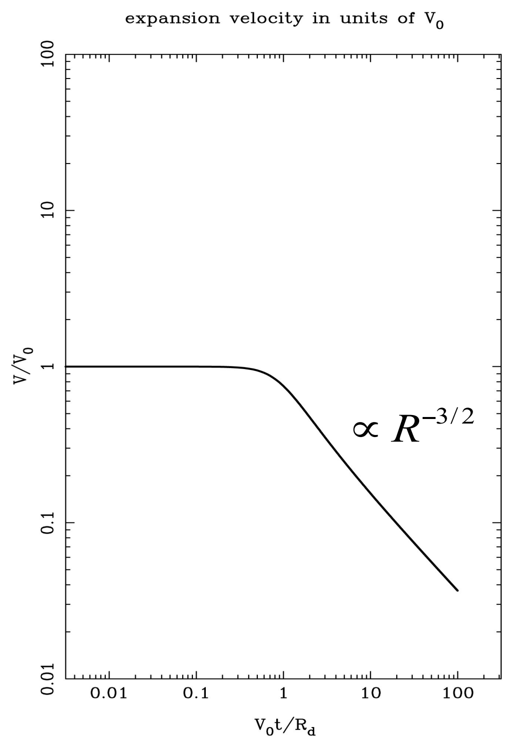
$$\frac{2}{5}R^{5/2} \approx (\frac{3E_{snr}}{4\pi C_\gamma \rho_o})^{1/2}t$$

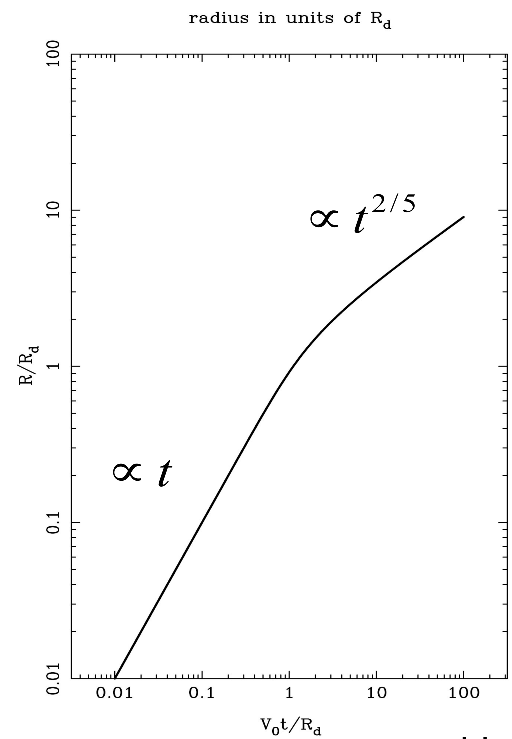
Assuming R(0) = 0

$$R = (\frac{75}{16}\frac{E_{snr}}{\pi C_\gamma \rho_o})^{1/5}t^{2/5}$$

$$V_s = (2/5)(\frac{75}{16}\frac{E_{snr}}{\pi C_\gamma \rho_o})^{1/5}t^{-3/5}$$

Fiandrini Cosmic Rays



# S-T phase: exact solution

The solution interior to the shock obeys the equations

$$\frac{\partial \rho}{\partial t} + \frac{1}{r^2}\frac{\partial}{\partial r}(r^2 \rho u) = 0,$$

$$\frac{\partial u}{\partial t} + u\frac{\partial u}{\partial r} = -\frac{1}{\rho}\frac{\partial P}{\partial r},$$

$$\frac{\partial}{\partial t}\left[\rho\left(\mathcal{E} + \frac{U^2}{2}\right)\right] + \frac{1}{r^2}\frac{\partial}{\partial r}\left[r^2 \rho u\left(\mathcal{E} + \frac{P}{\rho} + \frac{U^2}{2}\right)\right] = 0$$

with $\mathcal{E} = P/[\rho(\gamma - 1)]$, for an ideal gas.

Sedov recognised that the solution must be self-similar, *i.e.*, that at any time the pressure, density, velocity, *etc.* at all points interior to the shock can be expressed in terms of *single* similarity variable

$$\xi = r\left(\frac{\rho_0}{t^2 E}\right)^{1/5} = \xi_0 \frac{r}{r_s}.$$

Thus

$$\rho(r,t) = \rho_2 \alpha(\xi),$$

$$u(r,t) = u_2(t)\frac{r}{r_s}v(\xi),$$

$$P(r,t) = P_2(t)\left(\frac{r}{r_s}\right)^2 p(\xi),$$

where $\alpha$, $v$ & $P$ are dimensionless *time independent* functions of the similarity variable $\xi$. All three functions are deliberately scaled so that they reach the value unity at $\xi = \xi_0$.

The value of the constant $\xi_0$ is determined by the condition that total energy is conserved, which can be written

$$\int_0^{r_s(t)} \rho\left(\mathcal{E} + \frac{U^2}{2}\right) 4\pi r^2 dr = E$$

or

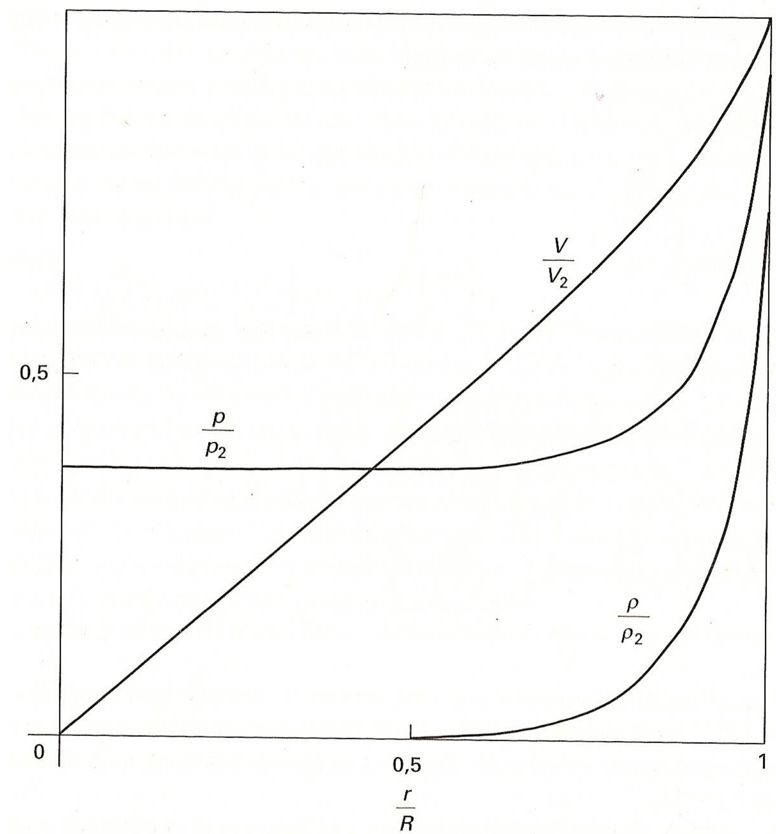
$$\frac{32\pi}{25(\gamma^2 - 1)}\int_0^{\xi_0} \left(p(\xi) + \alpha(\xi)v^2(\xi)\right)\xi^4 d\xi = 1.$$

# Sedov-Taylor: $\rho$, p, V distribution in the bubble

Integration has to be made numerically

The solution is self-similar: at any stage of the expansion the functional form is the same

NB: the quantities $P_2$, $V_2$ and $\rho_2$ are the values immediately post-shock as given by the RH conditions




Our assumption that all the swept up mass is in a thin layer of thickness $\delta R/R \sim 0.1$ behind the shock is well justified
Also the pressure is concentrated in the layer between the contact discontinuity and the blast wave

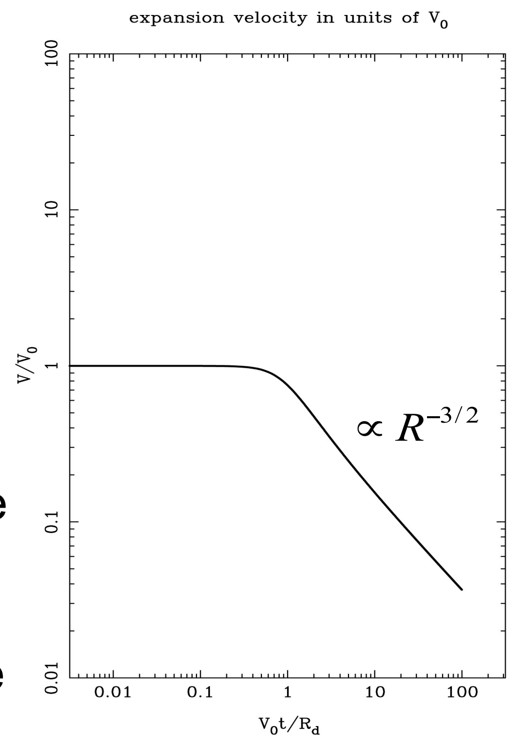
The pressure in the bubble decays as the bubble expands

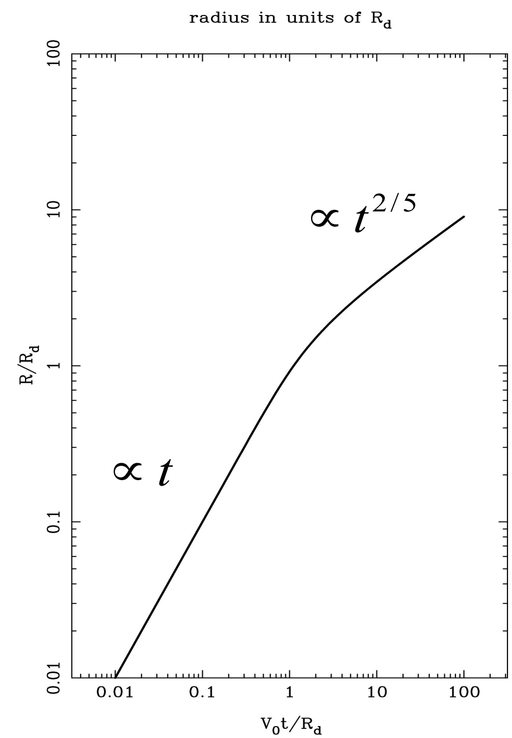
From $p_i \approx p_2 = \frac{2\rho_O V_s^2}{(\gamma+1)}$ One finds $p_i \sim t^{-6/5}$ or $p_i \sim R_s^{-3}$

This decay is simply an expansion loss as the internal pressure is converted into kinetic energy of the expanding shell

This energy conserving Sedov-taylor solution applies for R> $R_d$ and until the radiation losses become important

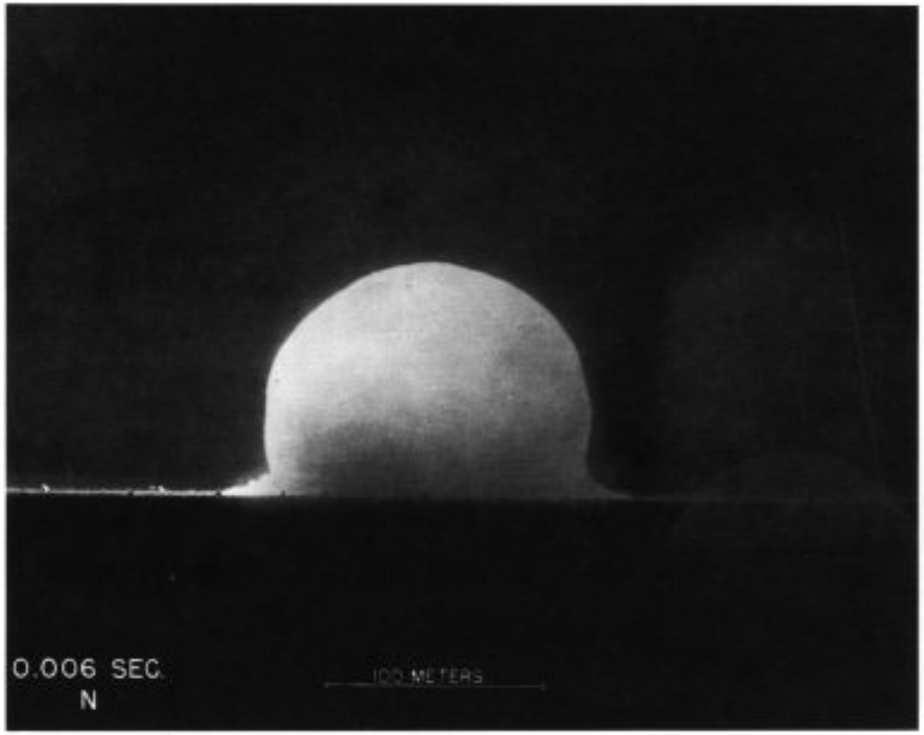
Radiative cooling makes the pressure inside the hot bubble decay faster and consequently the remnant looses energy and the expansion slows down more rapidly than in Sedov-taylor phase

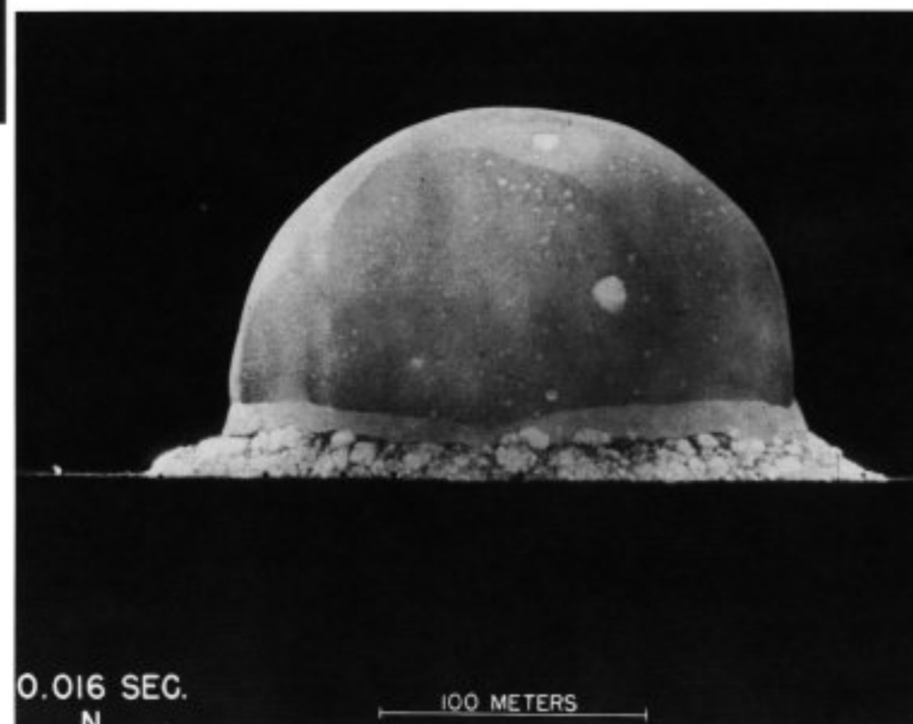



The cooling dominated stages of the evolution set in after about 10000 years

# Sedov & Taylor



# Sedov & Taylor

Crab Nebula: SNR exploded in 1054 ad
Distance 2000+-500 pc
Diameter ~1.7 pc

# Temperature

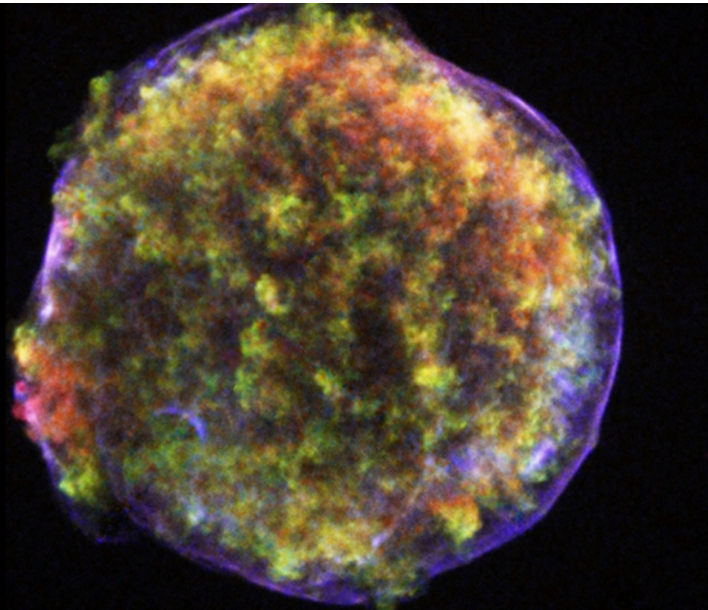
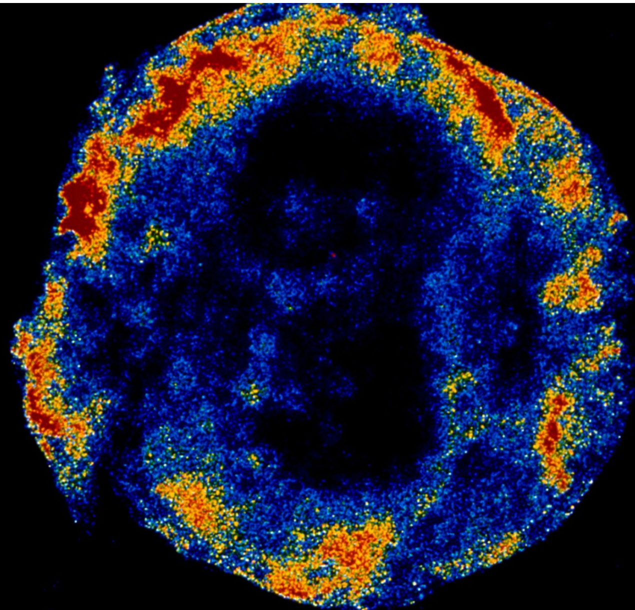
$$\frac{T_2}{T_1} = \frac{2\gamma(\gamma-1)}{(\gamma+1)^2} M_s^2 = \frac{2\gamma(\gamma-1)}{(\gamma+1)^2} \frac{V_s^2 \rho_o}{\gamma p_o} \qquad T_2 = \frac{2(\gamma-1)}{(\gamma+1)^2} \frac{V_s^2 \rho_o T_1}{p_o} = \frac{2(\gamma-1)}{(\gamma+1)^2} V_s^2 \frac{\mu}{R_{gas}}$$

For typical values, $T_2$ is several million K (>$10^7$ K) after $10^4$ years and the SNR is a bright X ray source

After $10^5$ years the post-shock T decreases to ~ $2\times10^5$ K, when radiative cooling becomes important and the energy conservation assumption breaks down

This occurs because at sufficiently low T, nuclei and electrons recombine so that radiation is no longer trapped in the shell and radiates away, cooling the shell

*Two pictures of the remnant of Tycho's supernova (AD 1572), a picture in X-Rays (left) , made with the CHANDRA satellite, and a radio picture made with the* Very Large Array *radio synthesis telescope (right). The X-ray picture shows the hot ($T \sim 10^8$ K) gas in the remnants interior in yellow. This is mostly* line emission *from exited nuclei. The blue radiation at the outer rim of the remnant is* synchrotron continuum emission, *caused by relativistic electrons moving in a weak magnetic field. The radio emission is also synchrotron radiation. It is believed that these relativistic electrons are accelerated at the outer shock.*
*This is a 'classical' remnant with a nearly perfect spherical shape. It is believed to be entering the Sedov-Taylor phase. Note the sharp outer edge of the remnant, which is believed to coincide with the position of the outer blast wave.*

# Tycho's Remnant (SN 1572AD)

X-Rays (CHANDRA Observatory) Radio (21cm)

# Snowplow phase

When the SNR becomes sufficiently old, radiative cooling becomes important and the total energy is no longer conserved

In the energy conserving Sedov-Taylor phase, pressure forces accelerate the swept up ISM converting thermal energy (which came from original explosion) into kinetic energy of the shell of swept up mass

Since radiative cooling depends on the particle density ($\sim n^2$), ie the higher is particle density, the higher is the radiation, and since the density in the shocked shell is much higher than in the bubble interior, most of the cooling occurs in the shocked ISM layer

# Snowplow phase

In the snowplow approximation, it is assumed that all the energy in the shocked shell is radiated away, but that the hot interior does not cool because there radiative processes are much less effective due to the very low density (ie the time scale of radiative cooling in the hot interior is much longer than in the shocked shell)

In such a case the shell must collapse until it becomes very thin because the pressure in the shell is decreased together with temperature but the pressure equilibrium at the contact discontinuity must still hold and the pressure in bubble interior is not changed much due to lack of radiative cooling, so the shell is compressed until a new pressure equilibrium is reached

The hot interior can be therefore considered as adiabatic → $p_i \sim \rho_i^{\gamma}$

Since the mass in the interior is conserved, one has $\rho_i = \frac{M_{ej}}{(4\pi/3)R_s^3}$

Combining the two relations one gets $p_i = (\frac{M_{ej}}{(4\pi/3)})^{\gamma} R_s^{-3\gamma}$ ~ $r_s^{-5}$ for $\gamma$=5/3

To be compared with the scaling law in the Sedov-Taylor phase ~ $R_s^{-3}$

# Snowplow phase

$$p_i = (\frac{M_{ej}}{(4\pi/3)})^{\gamma} R_s^{-3\gamma}$$

The motion of the collapsed shell, containing most of the mass, is driven by the pressure $p_i$ of the hot interior (neglecting the ISM pressure, which very low compared to $p_i$)

The motion equation of the expanding blast shock is then

$$\frac{d}{dt}(M(R_s)\frac{dR_s}{dt}) = 4\pi R_s^2 p_i(R_s) \quad = 4\pi(\frac{3M_{ej}}{4\pi})^{\gamma} R_s^{2-3\gamma} \equiv AR_s^{2-3\gamma}$$

Taking into account that $M_s(R_s) = \frac{4\pi}{3}\rho_o R_s^3$

$$\frac{4\pi}{3}\rho_o \frac{d}{dt}(R_s^3 \frac{dR_s}{dt}) = 4\pi(\frac{3M_{ej}}{4\pi})^{\gamma} R_s^{2-3\gamma} \Rightarrow \frac{d}{dt}(R_s^3 \frac{dR_s}{dt}) = \frac{3}{\rho_o}(\frac{3M_{ej}}{4\pi})^{\gamma} R_s^{2-3\gamma}$$

# Snowplow phase

$$\frac{d}{dt}(R_s^3\frac{dR_s}{dt}) = \frac{3}{\rho_o}(\frac{3M_{ej}}{4\pi})^\gamma R_s^{2-3\gamma}$$

We look for power law solutions **$R_s$(t) = B$t^\alpha$**, with B some constant and $\alpha$ to be determined

The index $\alpha$ is determined by substituting the trial solution into the equation

$$\alpha(4\alpha-1)B^3t^{(4\alpha-2)} = AB^\beta t^{(2-3\gamma)\alpha}$$

$$A = \frac{3}{\rho_o}(\frac{3M_{ej}}{4\pi})^\gamma$$

For the trial function to be solution, the two exponents must be equal

$$4\alpha - 2 = (2-3\gamma)\alpha$$

$$\alpha = \frac{2}{3\gamma+2}$$

Which for $\gamma$=5/3 yields $\alpha$ = 2/7=0.286

The actual value of the index of this pressure driven phase obtained by numerical, more accurate, simulations yields a value closer to 3/10=0.3

We have assumed uniform pressure in the bubble, no mixing at the contact discontinuity between bubble and shocked ISM, and neglected radiative losses

# Momentum driven phase

$$p_i = (\frac{M_{ej}}{(4\pi/3)})^\gamma R_s^{-3\gamma}$$

$$\frac{d}{dt}(M(R_s)\frac{dR_s}{dt}) = 4\pi R_s^2 p_i(R_s)$$

As the SNR evolves, the pressure inside the bubble decreases until becomes negligible since the remnant radiates away all its internal energy

In this limit, the total momentum is conserved since the total force acting the bubble is zero (hence the name momentum-conserving phase)

$$\frac{d}{dt}(M(R_s)\frac{dR_s}{dt}) = 0$$

$$M_s(R_s) = \frac{4\pi}{3}\rho_o R_s^3$$

This implies $M(R_s)\frac{dR_s}{dt} = constant$ $\Rightarrow$ $\frac{dR_s}{dt} = (\frac{3}{4\pi\rho_o})R_s^{-3}$

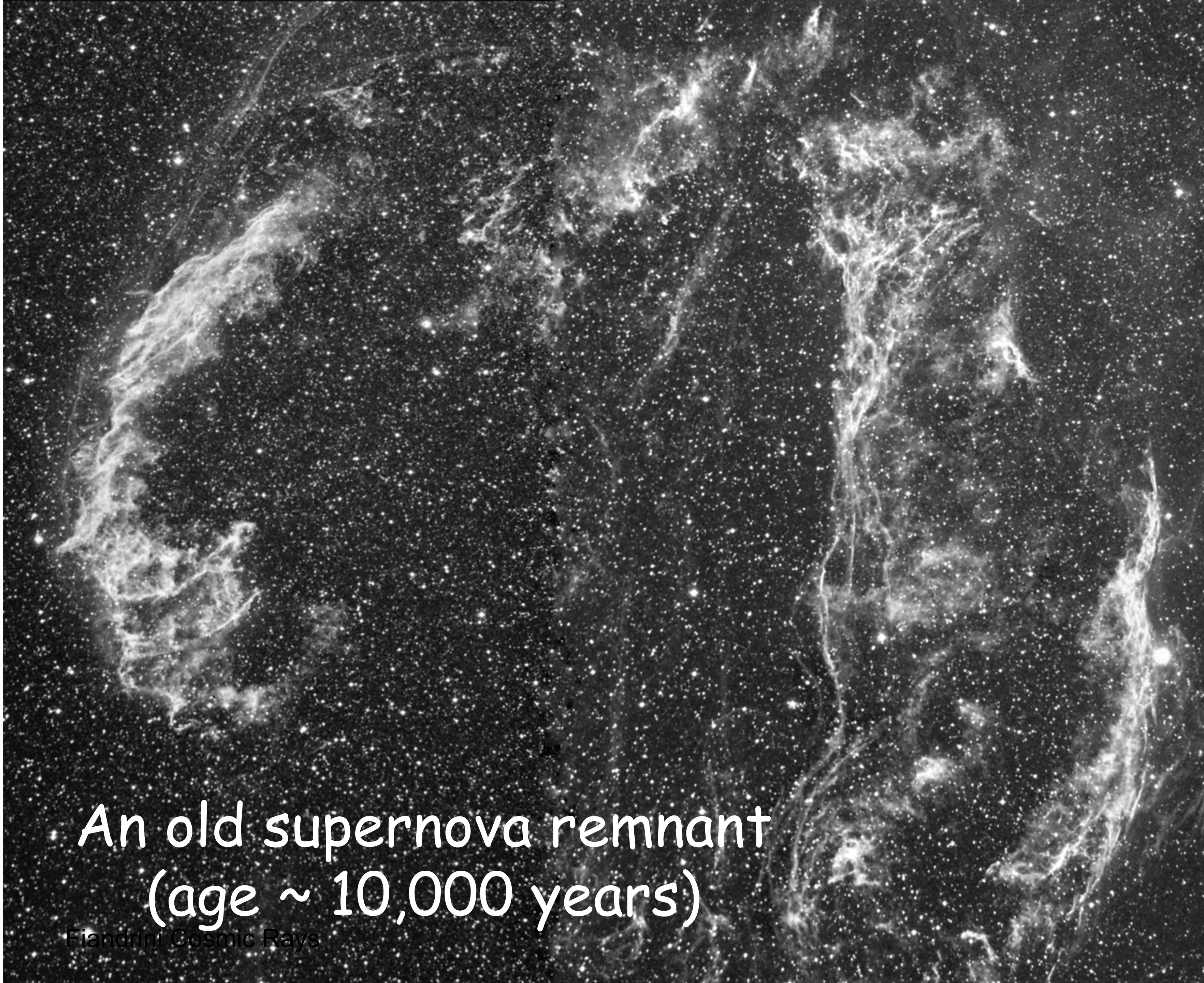
$$R_s^3 dR_s = (\frac{3}{4\pi\rho_o})dt \Rightarrow (1/4)R_s^4 dR_s = (\frac{3}{4\pi\rho_o})t \Rightarrow R_s = [4(\frac{3}{4\pi\rho_o})]^{1/4} t^{1/4}$$

# Coalescence phase

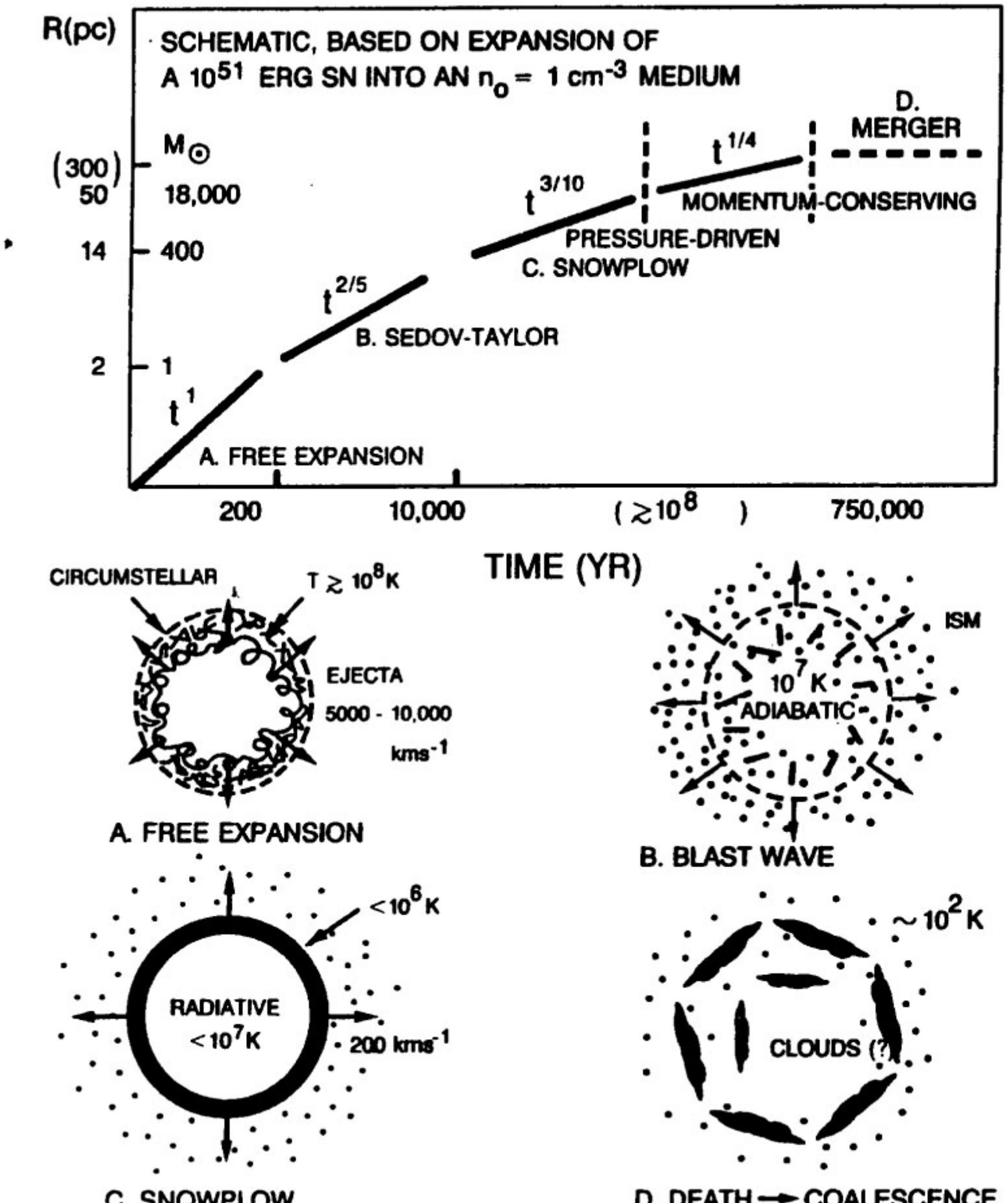
The momentum conserving phase lasts until the shock speed remains beyond sound speed in the ISM

The shock speed decreases as $R_s^{-3} \sim t^{-3/4}$, so at some point it approaches the sound speed and the shock itself disappear.
The material starts to straggle into the ISM (initially at sound speed) becoming part of it

Figure 7.15: *The old supernova remnant S147, which is in the process of dissolving into the general interstellar medium*. Photo credit: Robert Gendler


An old supernova remnant
(age ~ 10,000 years)

# STANDARD SNR EVOLUTION




Main properties:

Different expansion stages:

- Free expansion stage (t<1000 yr) $R \propto t$

- Sedov-Taylor stage (1000 yr< t< 10,000 yr) $R \propto t^{2/5}$

- Pressure-driven snowplow (10,000 yr< t<250,000 yr) $R \propto t^{3/10}$

- Momentum-conserving (250,000<t<750,000 yr) $R \propto t^{1/4}$

# MagnetoHydroDynamics (MHD)

We have seen the equations of hydrodynamics and some astrophysical applications of them, but we have neglected any electromagnetic phenomenon relative to the fluids

However, in astrophysics environments fluid temperatures are usually very high and the most of the atoms are completely ionized

This is particularly true for H and He, which make the most part cosmic matter, since their ionization potentials are quite low

# Plasmas

Plasmas are gases in which the constituent particles are electrically charged

The first consequence is that the electric fields are not important, because the abundances of free electric charges ensure that any E field would be short-circuited by their motion ie the charges move in such a way to cancel the (external) field

But if the fluid is in a B field, its motion wrt B builds up E fields too by induction and this generates currents

In turn, currents are influenced by the B fields and generate new B fields, that influence again the fluid motion

# Plasmas

The plasmas are made of particles of opposite sign

In a volume containing many charged particles, we expect that volume to be close to charge-neutral, since any charge imbalance would produce strong electrostatic forces to restore charge neutrality

Charge imbalances may exist only over a short distance (called Debye length) or for a short period of time (the inverse of the so called plasma frequency)

What makes the macroscopic behavior of plasmas so different from that of neutral gases is the fact that an electromotive force applied to a plasma can drive large currents→ volumes of plasmas can substain large currents in spite of being nearly charge neutral

Since many charged particles in a plasma can interact simultaneously through long range electromag interactions, there can be many phenomena in plasma which are caused by collective interactions

# Plasmas

In neutral gases, the interactions between partilces are mediate by collisions

In a completely ionized plasma, instead, the interactions are mediated by long range elm forces

In a partially ionized gas, both processes play a role

if we put a test charge q in the plasma, the presence of many charges of opposite sign around q screens the elm fields felt by "far" charges

The characteristic length over which the screening occurs is a measure of the assumption of charge neutrality of the overall gas: the Debye length

# Plasmas: Debye length

To understand how good is the assumption of charge neutrality is, let us consider the charge separation produced by introducing a charge q inside a plasma

If $n_e$ and $n_i$ are the number densities of e- and ions, the charge density at a point x is

$$\nabla^2\Phi = -4\pi(n_i - n_e)e$$

If the plasma in thermodynamical equilibrium and n is the density of e- and ions far away from the charge q, then we expect, according to Maxwell-Boltzman distribution

$$n_i = n \times exp(-\frac{e\Phi}{kT}) \qquad n_e = n \times exp(\frac{e\Phi}{kT})$$

Substituting in the potential eqn $\nabla^2\Phi = -4\pi ne(exp(-\frac{e\Phi}{kT}) - exp(\frac{e\Phi}{kT}))e$

Usually $e\Phi \ll kT$ →we may expand in Taylor series the exp, neglecting higher terms

$\nabla^2\Phi = -\frac{\Phi}{\lambda_D^2}$ with $\lambda_D = (\frac{kT}{8\pi ne^2})^{1/2}$ Called Debye length

# Plasmas: Debye length

$$\nabla^2\Phi = -\frac{\Phi}{\lambda_D^2} \qquad \lambda_D = (\frac{kT}{8\pi n e^2})^{1/2}$$

The solution is easy to find assuming a spherical symmetry $\nabla^2 = \frac{1}{r^2}\frac{\partial}{\partial r}(r^2\frac{\partial}{\partial r})$

We get $\Phi(r) = \frac{q}{r}e^{-r/\lambda_D}$

It thus appears that the effect of the charge is screened beyond a distance $\lambda_D$

So a plasma can be considered charge neutral when distances larger than the Debye lehgth are considered

Altough the E field of a charge in principle extends to ∞, the influence on a charged particle in a plasma is effectively felt to a distance $\lambda_D$, ie within a volume $\lambda_D^3$, called the Debye volume

Hence the nbr of particles interacting with q is $n\lambda_D^3$: this is a measure of the number of particles which can interact simultaneously

# Plasma parameter

The plasma parameter is defined as $g = 1/n\lambda_D^3 = \frac{(8\pi)^{3/2} e^3 n^{1/2}}{(kT)^{3/2}}$

When g is smaller, there is more collective interaction in the plasma (note g is smaller for smaller n)

Therefore the nbr of particles interacting collectively is more for a low density plasma, the Debye length being less effective so that Debye volume is much larger

The average distance between the particles of a plasma is of the order of $n^{-1/3}$ → the average potential energy between a pair of nearby particles is of the order of $e^2 n^{1/3}$

Hence the ratio of potential and average kinetic energy (~ kT) is

$$\frac{< PE >}{< KE >} \sim \frac{e^2 n^{1/3}}{kT} \propto n^{1/3}$$

Another interpretation of g is, therefore, that it is a measure of the potential energy of interactions compared to kinetic energy: when g is small (as for low n), the interaction amongst particles is weak, but a large nbr of particles interact simultaneously. On the other hand, a larger g implies few particles interacting collectively, but interacting strongly

The limit of small g is referred to as the ***plasma limit***

$$g = \frac{(8\pi)^{3/2} e^3 n^{1/2}}{(kT)^{3/2}}$$

$$\lambda_D = (\frac{kT}{8\pi n e^2})^{1/2}$$

# Types of plasmas

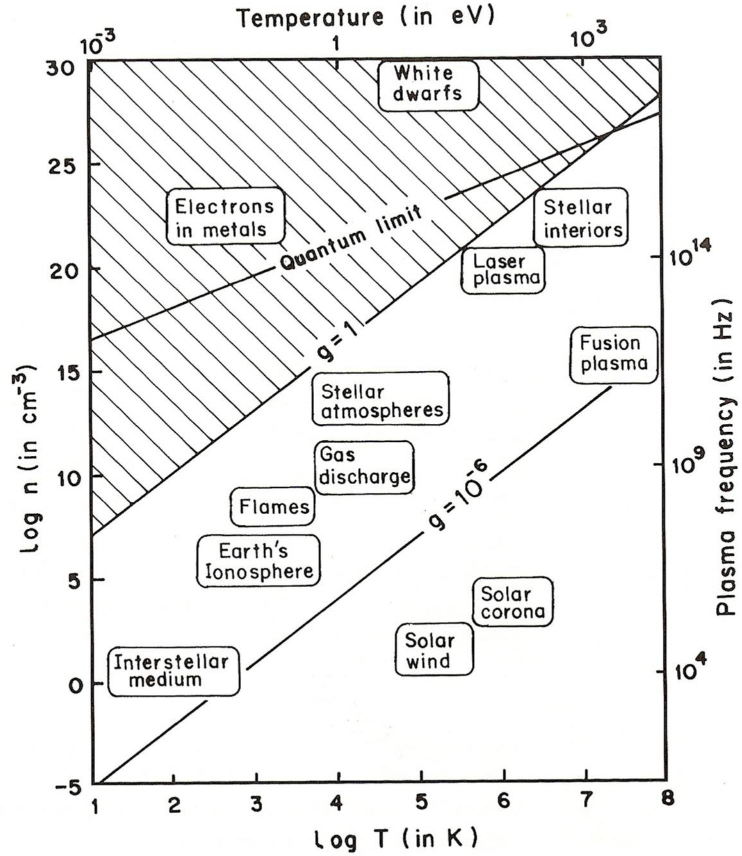
The characteristics of a plasma are determined by n and T

The condition g<1 is a requirement for a gas to be defined a plasma

In the region above g=1, the gas do not behave like a plasma

Above quantum limit, we have to use proper quantum mechanics, as for electrons in a metal and white dwarfs interiors (as for neutron stars)

In between QL and g=1 we find solids, liquids and crystals which can not be considered plasmas




The ratio $$\frac{< PE >}{< KE >} \sim \frac{e^2 n^{1/3}}{kT} \propto n^{1/3}$$

Implies that for small g plasmas kinetic energy exceeds potential energy, so that it can be treated as a perfect gas

This holds even for material at the centres of stars!